

Structure of poly(propyl ether imine) dendrimer from fully atomistic molecular dynamics simulation and by small angle x-ray scattering

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We study the structure of carboxylic acid terminated neutral poly(propyl ether imine) (PETIM) dendrimer from generations 1–6 (G1–G6) in a good solvent (water) by fully atomistic molecular dynamics (MD) simulations. We determine as a function of generation the structural properties such as radius of gyration, shape tensor, asphericity, fractal dimension, monomer density distribution, and end-group distribution functions. The sizes obtained from the MD simulations have been validated by small angle x-ray scattering experiment on dendrimer of generations 2–4 (G2–G4). A good agreement between the experimental and theoretical value of radius of gyration has been observed. We find a linear increase in radius of gyration with the generation. In contrast, R_g scales as $\sim N^x$ with the number of monomers. We find two distinct exponents depending on the generations, $x=0.47$ for G1–G3 and $x=0.28$ for G3–G6, which reveal their nonspace filling nature. In comparison with the amine terminated poly(amidoamine) (PAMAM) dendrimer, we find that R_g of G th generation PETIM dendrimer is nearly equal to that of $(G+1)$ th generation of PAMAM dendrimer as observed by Maiti *et al.* [*Macromolecules* **38**, 979 (2005)]. We find substantial back folding of the outer subgenerations into the interior of the dendrimer. Due to their highly flexible nature of the repeating branch units, the shape of the PETIM dendrimer deviates significantly from the spherical shape and the molecules become more and more spherical as the generation increases. The interior of the dendrimer is quite open with internal cavities available for accommodating guest molecules, suggesting the use of PETIM dendrimer for guest-host applications. We also give a quantitative measure of the number of water molecules present inside the dendrimer. © 2006 American Institute of Physics. [DOI: [10.1063/1.2194538](https://doi.org/10.1063/1.2194538)]

I. INTRODUCTION

Dendrimer,¹ regularly branched polymers, of different initiator cores, branches, and the peripheral groups have been synthesized^{2–8} in the recent decades. Numerous kinds of experiment have been carried out to prove the potential application of this type of new material in biochemical, medical, technical, and industrial fields.^{9–12} In last two decades lots of efforts have been made to design and synthesize biocompatible dendrimer with different monomers having a variety of functionalities. The atomistic characterization of dendrimer structure has lagged this rapid progress in synthesis and design.¹³ The problem is that these molecules possess an

enormous number of energetically permissible conformations, and in solution there is rapid interchange between them. Thus diffraction techniques yield little structure information. Also many generations involve the same monomers, making it difficult to extract precise information about the local structure. Thus the first precise experimental data about the gross size came from size exclusion chromatography (SEC), which is now being complemented with small angle neutron scattering (SANS) and small angle x-ray scattering (SAXS) to determine gross size and some structural details of dendrimers.^{14–16} In recent years, computational and theoretical techniques^{17–35} proved to be very effective in elucidating the structural properties of the dendrimer in both good and bad solvents.^{18,19} Many of these theoretical results agree very well with the available experimental data^{36,37,15,16} on such systems.

Recently Jayaraman and co-workers³⁸ have synthesized

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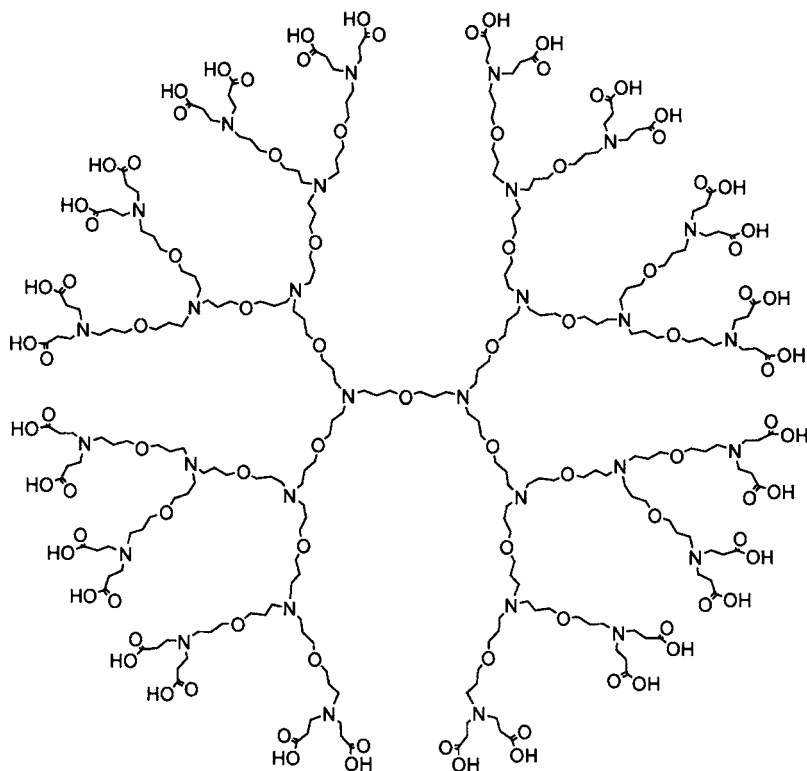


FIG. 1. Two-dimensional schematic representation of fourth generation (32 acid groups) PETIM dendrimer.

poly(propyl ether imine) (PETIM) dendrimers and reported their cytotoxic properties. Evaluation of the extent of cytotoxicity indicates that the toxicity levels of these dendrimer are very mild and this points to the possibility of using PETIM dendrimer in various applications. However, so far no structural information is available for these dendrimers, which is essential for their application ranging from drug delivery to molecular encapsulations. Here we report a comprehensive understanding of the structure of PETIM dendrimer using fully atomistic simulations and have used SAXS to validate some of our simulation results.

Following is the outline of the paper. In Sec. II A we describe the structure and the atomistic feature of the PETIM dendrimer, in Sec. II B the experimental details of the SAXS, and in Sec. II C the details of building of the atomistic model and simulation methods. In Sec. III we have summarized the results obtained from the simulation as well as SAXS studies. Finally, the summary of the main results and the conclusions are given in Sec. IV.

II. SYSTEM AND METHODS

A. Samples

The synthesis and cytotoxicity studies of carboxylic acid terminated PETIM dendrimer (up to generation 3) reported³⁸ recently prove its excellent biocompatibility. PETIM dendrimer starts growing (Fig. 1) three dimensionally from the oxygen as the core and branches out at each tertiary nitrogen, which is separated by eight-bond spacer for each generation of the dendrimer. The spacer containing all the sp^3 hybridized atoms of the PETIM dendrimer is flexible enough to have numerous numbers of molecular conformations. As the spacer length of the PETIM dendrimer (eight bonds) is larger than poly(amidoamine) (PAMAM) (seven bonds) and poly-

(propylene imine) (PPI) (four bonds) dendrimers, the size of PETIM is expected to be larger than that of PAMAM dendrimer as well as of PPI dendrimer for the same number of terminal groups and generations. Due to its larger spacer length, the compactness and space filling nature of the molecule are expected to be different compared to those of PAMAM and PPI dendrimers where the branching occurs at shorter distances. The fully atomistic description of the PETIM dendrimer has been described in Table I.

B. Small angle x-ray scattering

1. Experimental section

The samples were loaded in thin quartz capillary and were subjected to $Cu K\alpha$ (0.154 nm) radiation from a sealed tube x-ray generator (Philips, PW3830) with a line focus beam. A small angle resolution of 0.01 \AA^{-1} was achieved using a small angle Kratky camera (Hecus M. Braun, SWAXS) with line collimation. A slit size of $200 \mu\text{m}$ and a beamwidth of 10 mm were used in all of our experiments. The instrument resolution [full width at half maximum (FWHM)] was found to be 0.08 nm^{-1} . A one-dimensional

TABLE I. Atomistic description of PETIM dendrimer.

Generations (G)	No. of terminal group (N_t)	Molecular weight (g/mol)	Total No. of atoms (N)	Total No. of nitrogen atoms (N_n)
1	4	420.4547	61	2
2	8	1 169.4149	181	6
3	16	2 667.3269	421	14
4	32	5 663.1509	901	30
5	64	11 654.7988	1861	62
6	128	23 638.0938	3781	126

(1D) position sensitive detector (MBraun PS D50M) was used to measure the scattered intensity. An adjustable sample holder with temperature regulation unit (Anton PAAR, KHR) was used to maintain the temperature of the samples at 28 °C. The sample chamber was evacuated to ~ 0.01 mbar to prevent stray scattering. Typical exposure times were about 8 h.

2. Theoretical background

SAXS is a well-established technique for obtaining information about the size, shape, and structural correlations of macromolecules in solution. The scattered intensity from a collection of macromolecules is given by the expression

$$I(q) = AP(q)S(q), \quad (1)$$

where q is the wave vector transfer and is given by $q = (4\pi \sin \theta)/\lambda$, θ is the scattering angle, λ is the wavelength of the radiation used, A is the scattering amplitude which depends on the number density and electron density contrast, $P(q)$ is the particle form factor and has information regarding the size and shape of the macromolecule, and $S(q)$ is the structure factor having information regarding interparticle correlations. In a dilute collection of macromolecules $S(q) \sim 1$ and Eq. (1) reduces to

$$I(q) = AP(q). \quad (2)$$

For a collection of dilute spheres with radius of gyration R_g , the scattered intensity $I(q)$ in the low q limit $qR_g \ll 1$ is given by the Guinier relation

$$I(q) = I(0) \exp \left[-\frac{q^2 R_g^2}{3} \right]. \quad (3)$$

A log-linear plot of $I(q)$ versus q^2 in the low q limit will result in a straight line and its slope yields the radius of gyration R_g .

C. Simulation

The initial three-dimensional (3D) molecular models of various generations PETIM dendrimer were built as follows: First the core fragment, the branching repeat unit, and terminal repeat unit were built and energy minimized using 3D sketure of the CERIU2 software.³⁹ The partial charges on the atoms on each fragment were derived using the charge equilibration (QEq) method.⁴⁰ Then the fragments were assembled properly again using 3D sketure of the CERIU2 software to make a given generation and was optimized for 5000 steps using conjugate gradient minimization. Dreiding force field⁴¹ was used to describe the interatomic interactions. The charge equilibrated minimized structures of all generations have been directed through the annealing algorithm using the open force field (OFF) module of CERIU2 software as follows.

The initial minimized structure was heated at a rate of 100 K/4 ps from 300 to 2000 K, followed by quenching to 1000 K at the same rate, followed by four such cycles between 1000 and 2000 K, and finally cooling to 300 K. The annealed structures were solvated with TIP3P water model using the xleap module of AMBER 7.⁴² The box dimensions

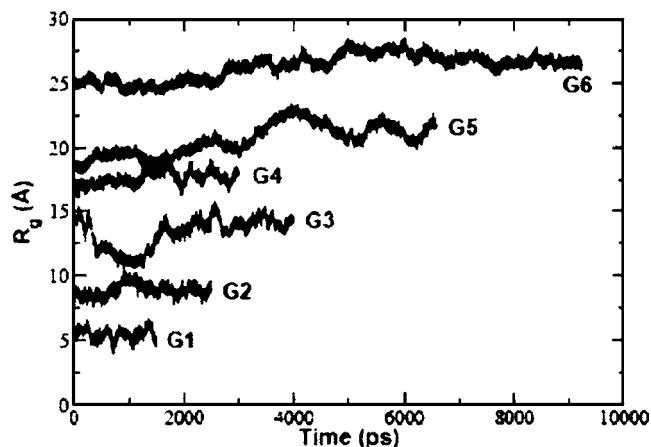


FIG. 2. Variation of the radius of gyration with time (ps) for generations 1–6 (upwards).

were chosen in order to ensure a 10 Å solvation shell around the dendrimer structure. This procedure resulted in solvated structures, containing between 3301 atoms for G1 and 67669 atoms for G6. Molecular dynamics (MD) simulation was performed using the AMBER 7 (Ref. 42) software suite, using the Dreiding force field.⁴¹ The solvated structures were subjected to 1000 steps of steepest descent minimization of potential energy, followed by another 2000 steps of conjugate gradient minimization. During this minimization the dendrimer structure was kept fixed in their starting conformations using a harmonic constraint with a force constant of 500 kcal/mol/Å². This allowed the reorganization of the water molecules to eliminate bad contacts with the dendrimer structure. The minimized structure was then subjected to 45 ps of MD, with 2 fs time step. During the dynamics, the system was gradually heated from 0 to 300 K with harmonic constraints on the solute using the SHAKE method. This was followed by 200 ps constant-volume–constant-temperature (*NVT*) dynamics with a temperature-coupling constant of 0.5–1.0 ps on the solute. Finally, 2–10 ns (depending on the generations) *NPT* unrestrained production dynamics was carried out with a time constant for heat bath coupling of 1 ps. The electrostatics interactions were evaluated with the particle mesh Ewald⁴³ (PME) method, using a real space cutoff of 9 Å. This simulation times proved to be long enough to get equilibrium properties as seen from the time evolution of the radius of gyration of the dendrimer shown in Fig. 2.

III. RESULTS AND DISCUSSION

To characterize the structure and properties of dendrimers as a function of generation we have chosen the following quantities: radius of gyration, shape tensor, asphericity, monomer density distribution, molecular surface area, end-group distribution, solvent accessible surface, molecular volume, and spatial arrangement of branch points. We have also done SAXS experiment to compare our simulation results with those obtained from experiment. We have also studied the penetration of water inside dendrimer by computing the number of water inside the dendrimer as a function of the distance from the center of mass of the dendrimer.

TABLE II. Radius of gyration (R_g) obtained from three different ways as the function of generation of PETIM dendrimer. $R_g(a)$ =radius of gyration calculated from Eq. (4) $R_g(s)$ =radius of gyration obtained from the theoretical Guinier plot. $R_g(\text{SAXS})$ =radius of gyration obtained from the experimental Guinier plot. R_T =radius of gyration calculated considering the terminal group only. The quantities have been averaged over 500 ps after the equilibration.

No. of terminal groups	$R_g(a)$ (Å)	$R_T(\text{PETIM})$ (Å)	$R_g(s)$ (Å)	$R_g(\text{SAXS})$ (Å)
4	5.65±0.27	6.59	5.19±0.12	
8	8.90±0.37	10.64	8.19±0.36	7.86±0.1
16	14.14±0.39	16.88	13.18±0.24	10.74±0.16
32	17.73±0.41	20.58	16.10±0.53	14.88±0.21
64	21.12±0.60	23.73	19.85±0.33	
128	26.62±0.16	27.80	24.91±0.35	

A. Size and shape

To obtain a quantitative estimate of the average size of the dendrimer, the mean square radius of gyration $\langle R_g^2 \rangle$ defined as

$$\langle R_g^2 \rangle = (1/M) \left\langle \left[\sum_{i=1}^N m_i |r_i - R|^2 \right] \right\rangle \quad (4)$$

has been computed [$R_g(a)$ given in Table II] over the trajectory. Here R is the center of mass of the dendrimer, m_i and r_i are the mass and position vector of the i th atom, and M and N are the total mass and total number of atoms of the dendrimer. Table II shows the radius of gyration (denoted as $R_g(a)$) obtained from our simulation as a function of generation. To compare our simulation results we have also shown the size obtained from our SAXS studies, details of which are given the next section.

The radius of gyration of the dendrimers in solution has been determined from the small angle x-ray scattering (SAXS). Figure 3 shows the Guinier plots ($I(q)$ versus q^2) for three different generations of the dendrimers after subtracting the background due to the solvent and the capillary. The black lines are the fits to the data using Eq. (3). The concentration of the third generation [Fig. 3(b)] and fourth generation [Fig. 3(c)] dendrimers was 0.5 wt % and that for the second generation [Fig. 3(a)] was 1.5 wt % to have a good signal to noise ratio. In spite of this threefold increase in concentration over the third and fourth generation dendrimers, the scattering from the second generation was low and hence the data shown for the second generation were smoothed using a five-point adjacent-averaging method. The values of R_g obtained from Guinier plot for generations G2–G4 have been tabulated in Table II. Closely related to the results of SAXS experiments is the spherically averaged Fourier transform of the single particle density $I(q)$ given by Eq. (5)²²

$$I(q) = \frac{1}{4\pi N^2} \int_0^{2\pi} d\phi \int_0^\pi \sin \theta d\theta \left| \sum_{i=1}^N \exp[i\mathbf{q} \cdot \mathbf{r}_i] \right|^2, \quad (5)$$

where

$$\mathbf{q} = q \sin \theta \cos \phi \hat{x} + q \sin \theta \sin \phi \hat{y} + q \cos \theta \hat{z}$$

is the scattering vector and r_i is the position vector of the scattering center.

To make comparison with the SAXS data, we have calculated $I(q)$ using orientation averaging at intervals of 9° in both $\theta(0 < \theta < 180^\circ)$ and $\phi(0 < \phi < 360^\circ)$. Figure 4 shows the plot of $\ln(I(q))$ versus q^2 as obtained from our simulation data. Linear regression fit to the theoretical Guinier plots obtained in this way ($\ln(I(q))$ versus q^2) gives the radius of gyration R_g for each generation (denoted by $R_g(s)$ in Table II) and has been shown in Fig. 5 as a function of generation. Table II shows variation of R_g as a function of generation

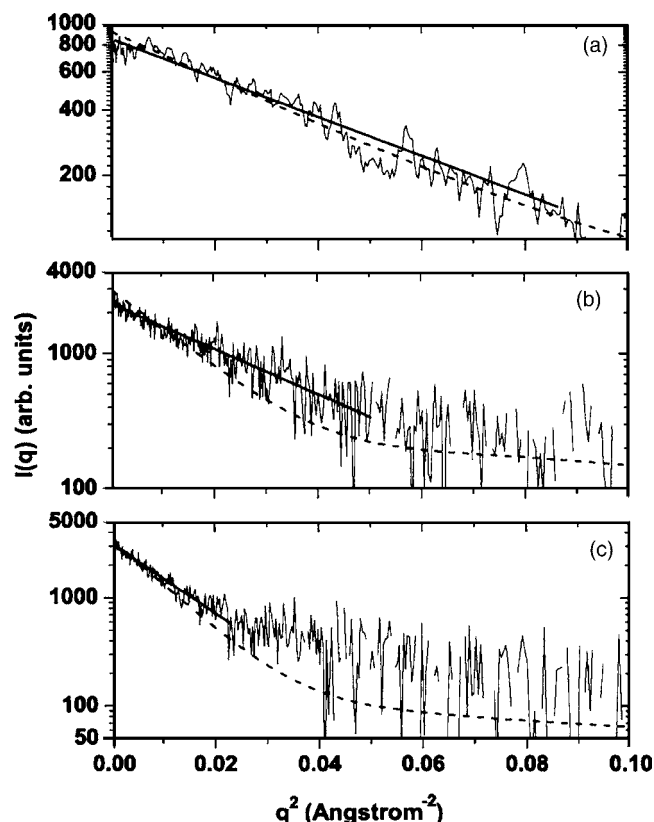


FIG. 3. Log-linear plots of the $I(q)$ vs q^2 for (a) PETIM generation 2, (b) PETIM generation 3, and (c) PETIM generation 4. The thin black lines are from the experiment and the dashed lines from MD simulation. The thick black lines are Guinier fits to the experimental data.

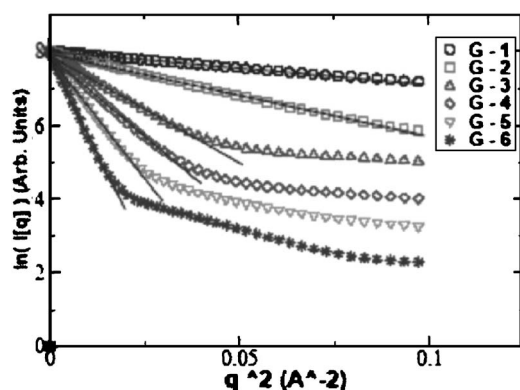


FIG. 4. Theoretical guinier plot ($\ln(I(q))$ vs q^2) for different generations of PETIM dendrimer. The black lines are the linear least squares fitting for the lower q value.

obtained from three different ways and we find a good agreement between our calculated R_g with the values obtained from the SAXS data.

Figure 5 shows that the radius of gyration increases linearly, having slightly different slopes with the generation number for all the three cases. This linear dependence of the radius of gyration on the generation has been observed for PPI dendrimer.¹⁷ We also see that our calculated values for R_g are consistently higher than that of experimental one. This discrepancy could be attributed to the fact that the shape of the dendrimer (discussed below) is far from spherical, an assumption that was made to extract R_g when using Eq. (3). The discrepancy between the experimental and simulation values may also be due to the solution conditions. The experiments were conducted in methanol solution, whereas the simulations were performed in water. The consistent higher R_g values for the aqueous solutions indicate that the propylene spacers in dendrimers are stretched under this solution condition. In methanol solutions, the propylene spacers experience relatively less hydrophilic environments and thus less affected by the solution condition. The lesser R_g values of dendrimers in methanol solutions reflect the extent of shrinking of the dendrimers under this solution condition. Swelling in aqueous solutions and shrinking in methanol solutions may contribute in addition to the asphericity for unequal R_g values in these two solvents.

The variation of root mean square radius of gyration with number of monomers shown in Fig. 6 follows the scaling relation $\langle R_g \rangle \sim N^\alpha$, with $\alpha=0.47$ for G1–G3 and $\alpha=0.28$ for G3–G6. It is clear that a single law $\langle R_g \rangle \sim N^\alpha$ does not

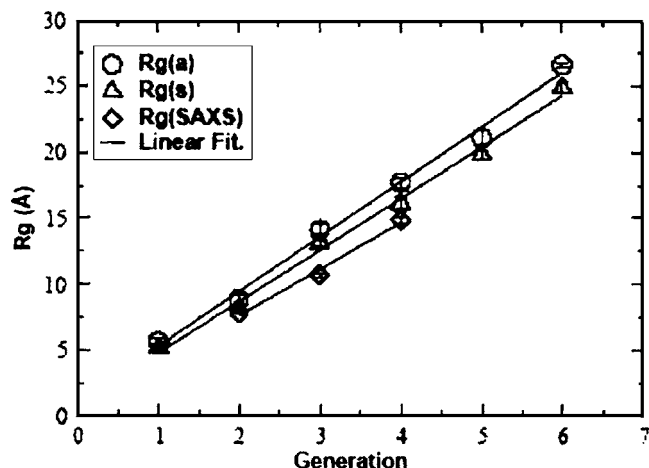


FIG. 5. Plot of the root mean square radius of gyration ($\langle R_g \rangle$) vs generation. All the averaging has been taken over 500 ps after the equilibration $R_g(a)$ = radius of gyration calculated from Eq. (4). $R_g(s)$ = radius of gyration obtained from the theoretical Guinier plot (Fig. 4). $R_g(\text{SAXS})$ = radius of gyration obtained from the experimental Guinier plot (Fig. 3).

describe the R_g dependence in the whole range of N studied here and far from the universal power law of the type $\langle R_g \rangle \sim N^{0.33}$ previously obtained for PAMAM and PPI dendrimers both theoretically and experimentally.^{18,19,44} Such nonuniversal scaling law behavior was recently found to be true for flexible dendrimers.⁴⁵ As the generation number increases, the dendrimer structure becomes more and more compact and space filling and the exponent is approaching the limiting value of 0.33. The above scaling exponent can be used to calculate the fractal dimension ($d_f=1/\alpha$) of the dendrimer from the relation

$$N \propto R_g^{1/\alpha}. \quad (6)$$

We find the fractal dimension d_f to be 2.1 for G1–G3 and 3.4 for G3–G6, respectively. PAMAM dendrimer shows this exponent as 3.0,¹⁹ which is equal to the dimensionality of the space indicating their space filling and compact structures. For lower generation PETIM dendrimer (G1–G3), the fractal dimension (~ 2.1) is far from the dimensionality of the space, which indicates that these molecules are nonspace filling and open in nature. For higher generation of the PETIM dendrimer, fractal dimension goes near to the dimensionality of the space; still the spatial arrangement of the branches within the molecule remains nonspace filling. This nonspace filling and open nature of the PETIM dendrimer have been further confirmed from the calculation of the single particle

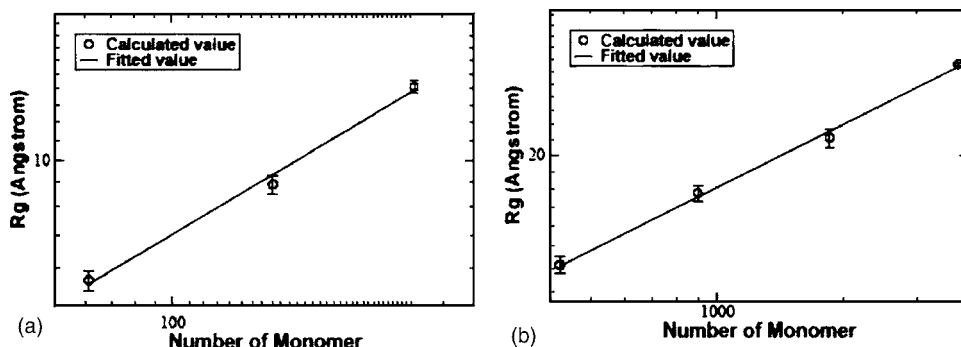


FIG. 6. Radius of gyration ($\langle R_g \rangle$) as the function of the number of monomers (atom) (in log-log scale), (a) G1–G3 and (b) G3–G6.

TABLE III. A comparison between radius of gyration of PETIM with that of amine functionalized PAMAM dendrimer as the function of generation. Our simulation has been done on the neutral dendrimer molecule. So we have taken the R_g of PAMAM observed by Maiti *et al.* Ref. 19 at high pH for the comparison. We also include the data observed by Han *et al.* Ref. 35 at slightly basic pH.

No. of terminal groups	R_g (PETIM) (Å)	R_g (PAMAM) (Å)	
		Maiti <i>et al.</i> ^a	Han <i>et al.</i> ^b
4	5.65		
8	8.90		6.0
16	14.14		8.0
32	17.73		10.0
64	21.13	16.78	13.0
128	26.63	20.67	17.0
256		26.76	21.3
512			27.0

^aReference 19.

^bReference 35.

form factor (see Kratky plots shown in Fig. S1 in the supplementary materials). Due to larger and flexible nature of the spacer, the molecules get larger span of space to orient themselves resulting in their nonspace filling nature.

Finally to make a comparison with the available data on other types of dendrimer, we find that the radius of gyration of PETIM dendrimer is larger than that of PAMAM dendrimer for a given generation. To a good approximation the radius of gyration of the G th generation of PETIM dendrimer is approximately greater or equal to that of the $(G+1)$ th generation (observed by Maiti *et al.*)¹⁹ and $(G+2)$ th generation of the PAMAM dendrimer (observed by Han *et al.*) (Table III).³⁵ Again this is a consequence of larger and more flexible nature of the spacer in PETIM dendrimer compared to other dendrimers such as PAMAM and PPI.

B. Shape

The shape of PAMAM dendrimer has been studied extensively using transmission electron microscopy⁴⁶ (TEM) and with tapping mode atomic force microscopy (AFM).⁴⁷ However, so far there are no data on the shape of the PETIM dendrimer. The flexibility and larger length of the repeat unit determine the shape of these molecules. To visualize the variation of the shape as a function of generation, we have shown a snapshot of the final configuration for each generations G1–G6 in Fig. 7. To provide a more quantitative criteria, the aspect ratio, which is the ratio of the two principle

moments of inertia, has been calculated and averaged over the dynamics trajectory. I_x , I_y , and I_z are three principle moments of inertia of the molecule, which have been found by diagonalizing the gyration metrics G ,⁴⁸

$$G_{mn} = \frac{1}{M} \left[\sum_i^N m_i (r_{mi} - R_m)(r_{ni} - R_n) \right], \quad m, n = x, y, z, \quad (7)$$

where r and R are the coordinates of atoms and the center of mass of the dendrimer, respectively, m_i is the mass of i th atom and M is the total mass of the dendrimer. The average values of the three principal moments of inertia are tabulated in Table IV, while Fig. 8(a) shows the average ratios for different generations of dendrimer. We see that I_z/I_x and I_z/I_y vary from 20.163 to 2.372 and from 5.54 to 1.58, respectively, from generations 1–6 (Table IV).

As the generation increases the value of the aspect ratio goes towards 1.0 (Table IV) indicating that the shape of the dendrimer approaches to the spherical one. More quantitative information about the shape of the dendrimer can be obtained by calculating the *asphericity*, introduced by Rudinck and Gaspari as⁴⁹

$$\delta = 1 - 3 \left(\frac{\langle I_2 \rangle}{\langle I_1^2 \rangle} \right), \quad (8)$$

where I_1 , I_2 , and I_3 are defined as

$$I_1 = I_x + I_y + I_z,$$

$$I_2 = I_x I_y + I_y I_z + I_x I_z,$$

$$I_3 = I_x I_y I_z.$$

The relative shape anisotropy of the simulated dendrimer shown in Fig. 8(b) shows that with the increase in generation, the dendrimer becomes a more and more compact spherical structure. The asphericity decreases from 0.4 for G1 to 0.05 for G6.

C. Radial monomer density profiles

The average radial monomer density $\rho(r)$ can be defined by counting the number $N(R)$ of atoms whose centers of mass are located within the spherical shell of radius r and thickness Δr . Hence, the integration over r yields the total number of atoms as

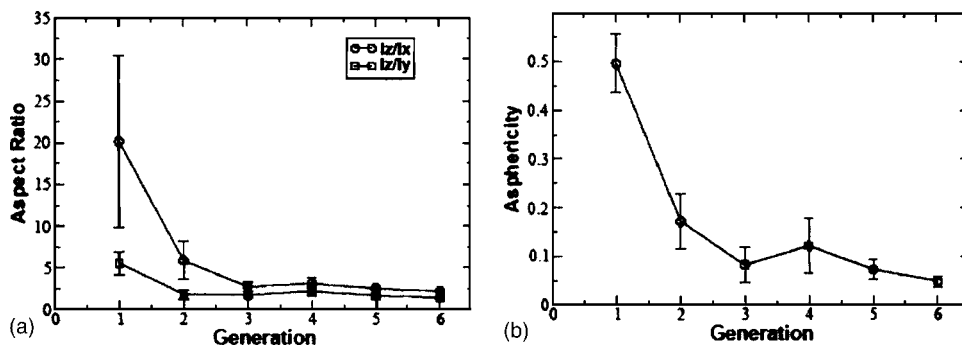


FIG. 7. (a) Aspect ratio as the function of generation, I_z/I_x and I_z/I_y . (b) Asphericity has been plotted as the function of generation.

TABLE IV. Information about the geometric properties (three principle moments of inertia, aspect ratio, and asphericity) of PETIM dendrimer in comparison. The quantities have been averaged over 500 ps after the equilibration.

Generation	$\langle I_z \rangle$	$\langle I_y \rangle$	$\langle I_x \rangle$	$\langle I_z/I_y \rangle$	$\langle I_z/I_x \rangle$	δ (asphericity)
1(4)	25.63	4.83	1.61	5.54±1.34	20.16±10.4	0.49±0.06
2(8)	45.30	225.67	8.43	1.85±0.53	5.86±2.3	0.17±0.06
3(16)	101.44	60.25	38.25	1.72±0.36	2.70±0.63	0.08±0.04
4(32)	173.99	83.13	57.56	2.19±0.6	3.10±0.75	0.12±0.06
5(64)	223.82	134.42	88.42	1.67±0.26	2.53±0.26	0.07±0.02
6(128)	327.20	232.95	148.87	1.41±0.17	2.20±0.15	0.04±0.01

$$N(R) = 4\pi \int_0^R r^2 \rho(r) dr.$$

In Fig. 9 we show the radial monomer density for various generation PETIM dendrimers in water. In each case the plot shows the contributions to a particular generation from each of its component generations. We take the origin as the center of mass. We see a very high-density region around the origin and a tailing zone in which the monomer density is gradually decreasing with the radial distance. This indicates that the core region is very dense compared to the middle of the dendrimer, which is fairly hollow supporting the dense core picture from earlier theoretical and computational studies.^{13,22,31} However, there is no constant density region in monomer density distribution in the middle of the dendrimer as has been observed for PAMAM dendrimers.¹⁸ The radial monomer density distribution for each subgeneration shows how the inner subgenerations are distributed throughout the interior of the dendrimer and indicates that there is significant back folding of the outer subgeneration. The extent of back folding increases with the increase in generations. It is

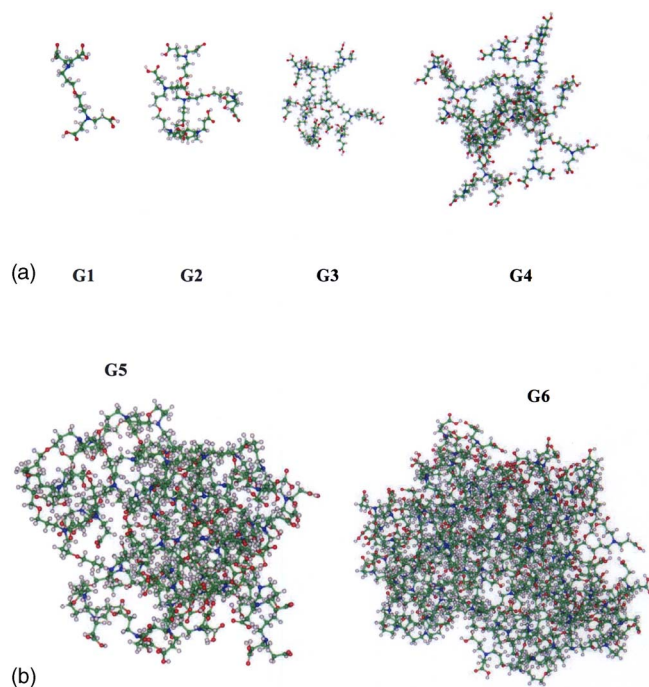


FIG. 8. Instantaneous snapshots of G1–G6 PETIM dendrimers after long MD simulations at $T=300$ K. All figures are to the same scale.

clear that the monomer density is higher at the core region for all the generation compared to the exterior of the molecule. On the basis of this observation it can be concluded that the subgeneration for a particular generation has folded back towards the core and makes the core region compact compared to the region far from the core. With increase in the generation number the monomer density increases and extends radically outward from the core. This dense core picture is in agreement with the results obtained by Boris and Rubinstein²¹ for the dendrimer containing the flexible repeat unit.

D. Water penetration

In a good solvent such as water PETIM has lots of internal voids and cavities, which can act as a binding site for small molecules for drug delivery and skin care products, and these cavities can accommodate a large number of water molecules as well. Due to the favorable interaction of water with the various functional groups of the dendrimer, significant number of water penetrates inside the dendrimer and it helps swelling the dendrimer. A quantitative estimate of the solvent penetration is given by counting the number of water bound by the dendrimer outer surface. Due to the nonuniformity as well as asphericity of the dendrimer surface, special care must be taken to identify the bound water, as simple spherical cutoff will overestimate the numbers of water within the dendrimer. To have an accurate estimate of the number of bound water we have used the following criteria:^{19,50} We first calculated the molecular surface area (MSA) for each of the dendrimer atom using a large probe radius (6 Å). With this probe radius the generated surface of the dendrimer becomes almost spherical and smooth. Those atoms with nonzero MSA represent the surface atoms of the dendrimer. Using these surface atoms we identify all the surface water that are within 4 Å of the surface atoms. Next we identify all the water close to the inner atoms (with zero MSA) excluding all the previously defined surface water. The number of bound water calculated this way is listed in Table V. This significant penetration of solvent molecules inside the dendrimer structure is in agreement with the recent SANS studies on poly(benzyl ether)⁵¹ and polycarbosilane dendrimers.⁵² In these experiments the number of solvent molecules inside the dendrimer was calculated from the change in neutron scattering density. We find that the number of bound water for the G th generation of the PETIM is larger than the number of water in same generation PAMAM den-

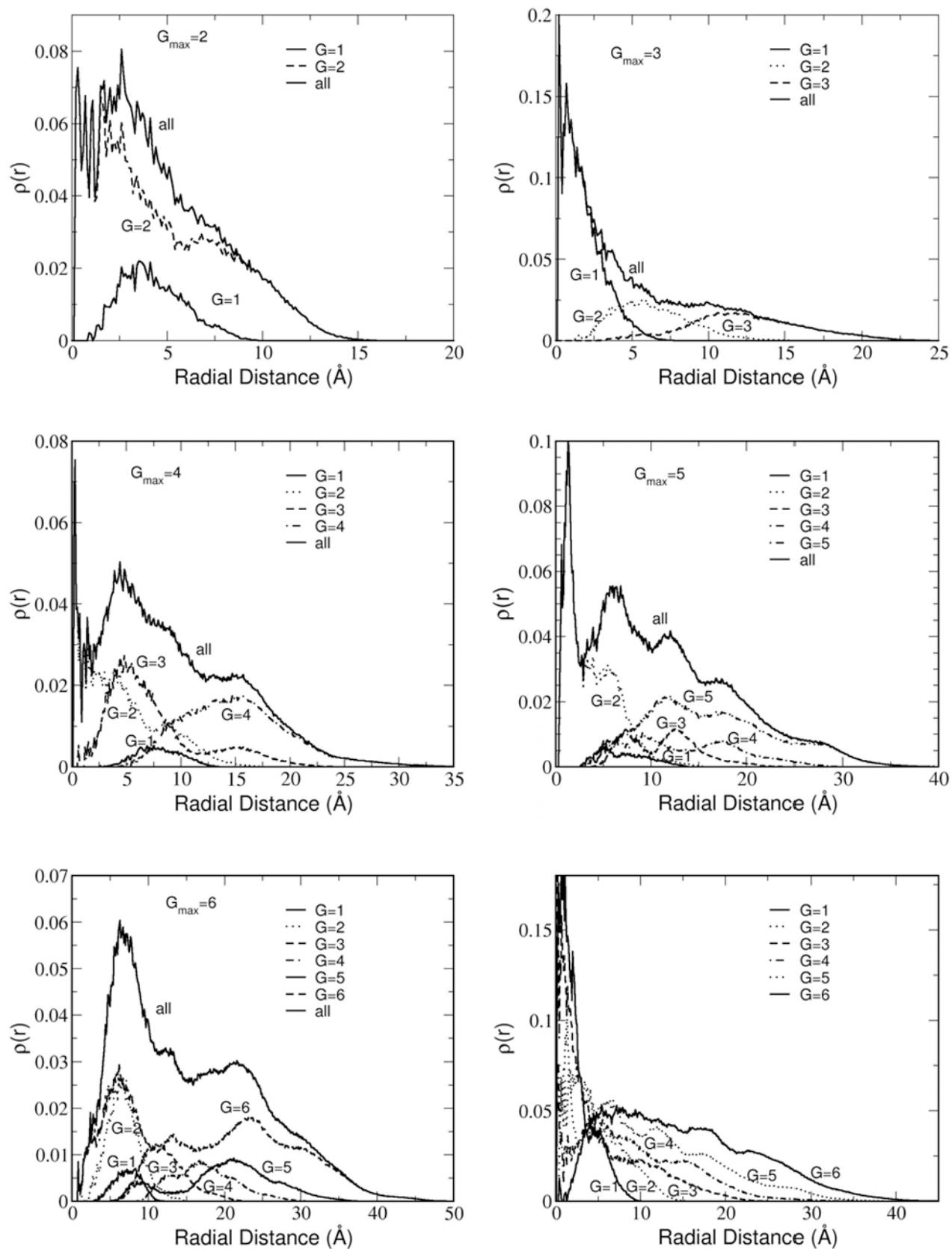


FIG. 9. Radial monomer densities (arbitrary units) for different generations of ethylene diamine (EDA) cored PAMAM dendrimers for all generations. Each figure is for a specific generation. The numbers shown were averaged from snapshots every 0.5 ps. The origin is at the center of mass. The last plot compares the total density profiles for all generations of PAMAM dendrimer from G1 to G6.

TABLE V. Average number of water inside the dendrimer for different generation PETIM dendrimers. For criteria used in identifying the bound water, see the text. For comparison we have also shown the number water inside PAMAM dendrimer with same number of terminal group.

No. of terminal groups	Bound water	
	PETIM	PAMAM
32	161	
64	491	138
128	1275	378
256		890

dendrimer. The number difference of bound water between G th generation of PETIM and $(G+1)$ th generation of PAMAM increases as we go to higher generation. For example, the number difference between PETIM-32 and PAMAM-64 is about 14%, where as for PETIM-128 and PAMAM-256 it is 30%. So for various applications lower generation PETIM dendrimer can be used in place of higher generation PAMAM to avoid some extra toxicity towards the living cell.

IV. CONCLUSION

SAXS has been used to study the size of the carboxylic acid terminated neutral poly(propyl ether imine) (PETIM) dendrimer for generations 2–4 (G_2 – G_4) in water. The size obtained from the SAXS measurement is in good agreement with the atomistic molecular dynamics simulation in explicit water. The dependence of the radius of gyration as a function of the number of monomers N does not obey the scaling law $R_g \sim N^{1/3}$ as has been observed for PAMAM and PPI dendrimers. Instead for higher generation PETIM dendrimer studied here (up to G_6), we find the scaling form to be $R_g \sim N^{0.28}$. The monomer density distribution shows the dense core nature of the dendrimer, which is expected for dendrimer having flexible repeat unit. Significant back folding is observed for all the generations studied. This, along with our previous simulation studies on PAMAM dendrimer, demonstrates that back folding is a universal phenomenon for dendrimer architecture. With the very flexible repeat unit the shape of the dendrimer is far from spherical one. We find significant penetration of solvent molecules in the interior of dendrimer molecules for all generations.

The following supporting information are available: (1) the xyz coordinates for a snapshot of the trajectory for each dendrimer from generation 1 to generation 6 in ASCII format, (2) the monomer density profiles with respect to the center of core of the dendrimer, (3) branch point distribution, and (4) Solvent accessible surface area and volume.

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- ¹G. R. Newkome, C. N. Moorfield, and F. Vögtle, *Dendrimers and Dendrons: Concepts, Synthesis, Applications* (Wiley-VCH, Weinheim, Germany, 2001).
- ²M. Ballauff, *Top. Curr. Chem.* **112**, 177 (2001).
- ³A. Adronov and J. M. J. Fréchet, *Chem. Commun. (Cambridge)* **2000**, 1701.
- ⁴M. A. Hearshaw and J. R. Moss, *Chem. Commun. (Cambridge)* **1999**, 1.
- ⁵K. Inoue, *Prog. Polym. Sci.* **25**, 453 (2000).
- ⁶A. W. Bosman, H. M. Janssen, and E. W. Meijer, *Chem. Rev. (Washington, D.C.)* **99**, 1665 (1999).
- ⁷S. Hecht and J. M. J. Fréchet, *Angew. Chem., Int. Ed.* **40**, 74 (2001).
- ⁸J. Nithyanandhan and N. Jayaraman, *J. Org. Chem.* **67**, 6282 (2002).
- ⁹A. Bielinska, J. F. Kukowska-Latallo, J. Johnson, D. A. Tomalia, and J. R. Baker, Jr., *Nucleic Acids Res.* **24**, 2176 (1996).
- ¹⁰D. S. Shah, T. Sakthivel, I. Toth, A. T. Florence, and A. F. Wilderspin, *Int. J. Pharm.* **208**, 41 (2000).
- ¹¹J. A. Hughes, I. A. Aronsohn, A. V. Avrutskaya, and R. L. Juliano, *Pharm. Res.* **13**, 404 (1996).
- ¹²H. Yoo, P. Sazani, and R. L. Juliano, *Pharm. Res.* **16**, 1799 (1999).
- ¹³D. A. Tomalia, A. M. Naylor, and W. A. Goddard III, *Angew. Chem., Int. Ed. Engl.* **29**, 138 (1990).
- ¹⁴S. Rosenfeldt, N. Dingenouts, M. Ballauff, P. Lindner, C. N. Likos, N. Werner, and F. Vögtle, *Macromol. Chem. Phys.* **203**, 1995 (2002); M. Ballauff and C. N. Likos, *Angew. Chem., Int. Ed.* **43**, 2998 (2004).
- ¹⁵A. Topp, B. J. Bauer, D. A. Tomalia, and E. J. Amis, *Macromolecules* **32**, 7232 (1999).
- ¹⁶T. J. Prosa, B. J. Bauer, and E. J. Amis, *Macromolecules* **34**, 4897 (2001).
- ¹⁷R. Scherrenberg, B. Coussens, P. van Vliet, G. Edouard, J. Brackman, and E. de Brabander, *Macromolecules* **31**, 456 (1998).
- ¹⁸P. K. Maiti, T. Cagin, G. Wang, and W. A. Goddard III, *Macromolecules* **37**, 6236 (2004).
- ¹⁹P. K. Maiti, T. Cagin, S. T. Lin, and W. A. Goddard III, *Macromolecules* **38**, 979 (2005).
- ²⁰E. Canetta and G. Maino, *Nucl. Instrum. Methods Phys. Res. B* **213**, 71 (2004).
- ²¹D. Boris and M. Rubinstein, *Macromolecules* **29**, 7251 (1996).
- ²²M. L. Mansfield and I. L. Klushin, *Macromolecules* **26**, 4262 (1993).
- ²³A. V. Lyulin, G. R. Davies, and D. B. Adolf, *Macromolecules* **33**, 6899 (2000).
- ²⁴A. V. Lyulin, G. R. Davies, and D. B. Adolf, *Macromolecules* **33**, 3294 (2000).
- ²⁵T. Cagin, G. Wang, R. Martin, W. A. Goddard III, and N. Breen, *Nanotechnology* **11**, 77 (2000).
- ²⁶K. J. Naidoo, S. J. Hughes, and J. R. Moss, *Macromolecules* **32**, 331 (1999).
- ²⁷P. G. de Gennes and H. Hervet, *J. Phys. (France)* **44**, L351 (1983).
- ²⁸S. Tretiak and S. Mukamel, *Chem. Rev. (Washington, D.C.)* **102**, 3171 (2002).
- ²⁹T. Tada, D. Nozaki, M. Kondo, and K. Yoshizawa, *J. Phys. Chem. B* **107**, 14204 (2003).
- ³⁰C. J. Hawker and J. M. J. Fréchet, *J. Am. Chem. Soc.* **112**, 7638 (1990).
- ³¹R. L. Lescanec and M. Muthukumar, *Macromolecules* **23**, 2280 (1990).
- ³²M. Murat and G. S. Grest, *Macromolecules* **29**, 1278 (1996).
- ³³A. M. Naylor, W. A. Goddard III, G. E. Keiffer, and D. A. Tomalia, *J. Am. Chem. Soc.* **111**, 2339 (1989).
- ³⁴C. B. Gorman and J. C. Smith, *Polymer* **41**, 675 (2000).
- ³⁵M. Han, P. Chen, and X. Yang, *Polymer* **46**, 3481 (2005).
- ³⁶S. Rosenfeldt, N. Dingenouts, M. Ballauff, P. Lindner, C. N. Likos, N. Werner, and F. Vögtle, *Macromol. Chem. Phys.* **203**, 1995 (2002).
- ³⁷N. C. Beck Tan, L. Balogh, S. F. Trevino, D. A. Tomalia, and J. S. Lin, *Polymer* **40**, 2537 (1999).
- ³⁸T. R. Krishna and N. Jayaraman, *J. Org. Chem.* **68**, 9694 (2003); T. R. Krishna, M. Belwal, U. S. Tatu, and N. Jayaraman, *Tetrahedron* **61**, 4281 (2005).
- ³⁹CERIUS2 Modeling Environment, Release 4.0., Accelrys Inc., San Diego, CA, 1999.
- ⁴⁰A. K. Rappe and W. A. Goddard III, *J. Phys. Chem.* **95**, 3358 (1991).
- ⁴¹S. L. Mayo, B. D. Olafson, and W. A. Goddard III, *J. Phys. Chem.* **94**, 8897 (1990).
- ⁴²D. A. Case, D. A. Pearlman, J. W. Caldwell *et al.*, AMBER 7, University of California, San Francisco, 2002.
- ⁴³T. Darden, D. York, and L. Pedersen, *J. Chem. Phys.* **98**, 10089 (1993).

- ⁴⁴G. Evmenenko, B. J. Bauer, R. Kleppinger, B. Forier, W. Dehaen, E. J. Amis, N. Mischenko, and H. Reynaers, *Macromol. Chem. Phys.* **202**, 891 (2001).
- ⁴⁵I. O. Gotze and C. N. Likos, *Macromolecules* **36**, 8189 (2003).
- ⁴⁶C. L. Jackson, H. D. Chanzy, F. P. Booy, B. J. Drake, D. A. Tomalia, B. J. Bauer, and E. J. Amis, *Macromolecules* **31**, 6259 (1998).
- ⁴⁷J. Li, L. T. Piehler, D. Qin, J. R. Baker, D. A. Tomalia, and D. J. Meier, *Langmuir* **16**, 5613 (2000).
- ⁴⁸D. C. Rapaport, *The Art of Molecular Dynamics Simulation* (Cambridge University Press, Cambridge, 1995).
- ⁴⁹G. Rudnick and G. Gaspari, *J. Phys. A* **4**, L191 (1986).
- ⁵⁰S. T. Lin, P. K. Maiti, and W. A. Goddard III, *J. Phys. Chem. B* **109**, 8663 (2005).
- ⁵¹G. Evmenenko, B. J. Bauer, R. Kleppinger, B. Forier, W. Dehaen, E. J. Amis, N. Mischenko, and H. Reynaers, *Macromol. Chem. Phys.* **202**, 891 (2001).
- ⁵²A. I. Kuklin, A. N. Ozerin, A. K. Islamov *et al.*, *J. Appl. Crystallogr.* **36**, 679 (2003).