On the pairing interaction in high T_c superconductors

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Abstract. It is argued, from the dependence of the critical current density and of T_c on applied magnetic field, that the pairing interaction in high T_c superconductors must be strongly field-dependent.

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We have earlier shown (Ravi Kumar and Chaddah 1988a, b) that the existing magnetization data on high T_c superconductors is explained by a critical current density J_c that decays exponentially with applied field. Since data on single crystal YBa₂Cu₃O₇ are also consistent with $J_c \approx J_{c0} \exp(-H/H_0)$, we have considered this behaviour to be intrinsic to high T_c superconductors. Farrell et al (1987) have observed a zero field $J_c \simeq 2 \times 10^7$ A/cm² at 4·2 K, and this is close to the depairing limit of $\approx 10^8$ A/cm². Their sample also shows an exponential decay of J_c with field. One possible reason for the observed decay of J_c with H could be that the depairing current itself decays exponentially with field. Within the BCS model this would correspond to the isotropic and k-independent pairing interaction V varying with field as 1/(H+K) where K is a constant. Such a variation is not possible for phonon-mediated interactions, but could be visualized in a pairing interaction mediated by local spin configurations as modelled, for example, by Emery (1987).

Such a pairing interaction would result in a sharp decay of the superconducting gap with increasing field, but no field dependent measurements of the gap have so far been reported. The dependence of T_c on the applied field is, however, available from the H_{c2} data. It will now be argued that the H_{c2} data also indicates a field dependent pairing interaction.

It has been reported in literature (Orlando et al 1987) that for $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ and $\text{RBa}_2\text{Cu}_3\text{O}_7$ compounds the resistive transition broadens rapidly with increasing applied field. A second feature is not emphasized but is apparent from the published data and is more important. The curvature of T_c vs H (or that of H_{c2} vs T) is concave near T_c . Such a curvature is contrary to the BCS result (or even WHHM theory) which yields a convex curvature for a field-independent V. Since this feature is not emphasized in literature we will discuss it in some detail.

Since the first non-trivial solution to the BCS gap equation will occur at the temperature where resistivity is zero (and not at the midpoint of the resistive transition), the $H_{c2}(T)$ relevant to any discussion strictly corresponds to the temperature at which $\rho = 0$. A close look at the reported resistive transitions in La_{1.85}Sr_{0.15}CuO₄ (Kwok *et al*

1987) and in YBa₂Cu₃O₇ (Maple 1987) shows that the H_{c2} vs T_c ($\rho = 0$) data is concave. Both these data are on polycrystalline samples, and tentative explanations can be offered on why T_c ($\rho = 0$) may not be intrinsic to the material. $H_{c2}(T)$ data on single crystal YBa₂Cu₃O₇ are now available (Moodera et al (1987)) with the crystal showing little twinning. The concave curvature of H_{c2} vs T is apparent and is seen in both the basal plane and along c-axis. This feature is irrespective of whether T_c is defined as where one has 10% or 50% or 90% of the normal state resistivity, and the curvature is a better approximation to an exponential decay as the definition of H_{c2} is shifted towards the zero resistivity condition. An exponential decay would again imply a strong field dependence in V.

Our argument that the concave curvature implies a field-dependent V and thus a spin-dependent pairing interaction is strengthened by the fact that such a concave curvature was seen by Sarkissian and Tholence (1986) in Y_9Co_7 where strong spin correlations are known to precede superconductivity.

The most important experimental data now required is a measurement of the superconducting gap as a function of applied field. This can yield the H-dependence of the pairing interaction.

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