

ON THE ANALYSIS OF THE BAND SPECTRUM OF CADMIUM.

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THE band spectrum of cadmium has been investigated by a number of observers including Mohler and Moore,¹ Winans,² Hamada,³ Kuhn and Arrhenius⁴ and Mrozowski.⁵ Winans has also discussed the origin of the bands and indicated the molecular states involved; he has also drawn some conclusions regarding the heat of dissociation of the cadmium molecule in the ground state, but Kuhn has shown that there is no basis for these conclusions. Kuhn and Arrhenius have tried to determine the heat of dissociation by studying the absorption of cadmium vapour at different temperatures. All these authors assume that the potential curve for the ground state possesses a very small depression with few vibrational levels and that the structure observed in several of the bands is almost entirely due to the vibrational levels of the upper states involved in their emission. A similar assumption has been made in the case of mercury and zinc. In a recent paper⁶ the author has succeeded in arriving at a vibrational analysis of five of the band systems of mercury and in thereby explaining a number of points which appeared obscure. This work, however, revealed a peculiarity in the behaviour of the ΔF^v 's of the ground state in that instead of decreasing steadily, they first went on increasing and then began to decrease, $\Delta^2 F$ during the decrease being the same as during the increase. This peculiarity made it difficult to draw definite conclusions regarding the heat of dissociation of the ground state. Since the similarity to be expected between the band spectra of zinc, cadmium and mercury has been experimentally proved to exist, it was thought that further light would be thrown on the question if the bands of cadmium were analysed and their structure compared with the structure of

¹ *Journ. Opt. Soc. America*, 1927, **15**, 74.

² *Phil. Mag.*, 1929, **7**, 556.

³ *Phil. Mag.*, 1931, **12**, 50.

⁴ *Zs. f. Phys.*, 1933, **82**, 716.

⁵ *Zs. f. Phys.*, 1934, **91**, 600.

⁶ *Proc. Ind. Acad. Sci.*, 1934, **1**, 166.

the corresponding bands of mercury. Accordingly the bands measured by Hamada were investigated and a vibrational analysis of these bands was made. It was found that by comparing their structure with that of the previously analysed bands of mercury, the bands corresponding to similar molecular transitions could be picked out without ambiguity. It was also found that the result so obtained was in entire agreement with the expectations of Hamada based on experimental grounds. For the sake of comparison the analysis of the bands attributed by Hamada to transitions from the $^1S_0 + ^3P_1$ state to the ground state is given in Table I while Table II gives the analysis of the corresponding bands of mercury observed by

TABLE I.

$^1S_0 + ^1S_0$											
$^1S_0 + ^3P_1$	26062	854	25208	486	24722						
	626		652		659						
	26688	828	25860	479	25381	362	25019	205	24814?	92	24722*
	597		595		586		596		567		569
	27285	830	26455	488	25967	352	25615	234	25381	90	25291
		463		488				479		482	
		26918	463	26455*				25860	87	25773	

* Used twice.

Hamada and attributed by him to a similar transition. Table III similarly contains the analysis of the cadmium bands attributed to transitions from the $^1S_0 + ^3P_0$ state and Table IV gives the results obtained for the similar Volkringer bands of mercury. Of the band systems observed by Mohler and Moore, Walter and Barratt⁷ attribute group II to CdO, groups III, IV and V to cadmium chloride and groups VI and VII to thallium chloride. Mrozowski,⁸ however, has established that group V belongs to Cd₂. On examining the bands of group II it was found that they fitted into the present scheme of vibrational levels (Table V) and their structure showed them to be analogous to the Steubing bands of mercury (given in Table VI). This similarity in their structure lends support to the opinion of Winans that the group II is due to transitions from the ground state to the $^1S_0 + ^1P_1$ state by absorption, and not due to CdO as Walter and Barratt think. The other groups given by Mohler and Moore do not seem to fit into the present scheme, except group V in which the difference $\approx 205 \text{ cm.}^{-1}$ characteristic of the ground state occurs twice in a system containing only five bands. This group may therefore belong to Cd₂ as has been shown by Mrozowski. From the

⁷ *Proc. Roy. Soc. A.*, 1929, **122**, 201.

⁸ *Zs. f. Phys.*, 1930, **62**, 314.

TABLE II.
Hamada's Bands of Mercury.

$^1S_0 + ^1S_0$ $v'' =$	8	9	10	11	12	13	14	15
$^1S_0 + ^3P_1$ $v' = 0$								
1	30891							
2	(29961)							
3	496?							
4	30457							
	525							
5	30982							
6	31810	1067	30743	735	30008			
			337		353			
7			31080	719	30361			
					336			
8				30697	483	30214		
				335		339		
9			31760	728	31032	479	30553	288
							288	288
10						30840	287	30553* 96
						240		240
11						31080*	288	30792 95
								190
12								30982* 91
								149
13								31131* 99
								99
14								31131 51
								51
15								31182 51
								31131*

* Used twice.

TABLE III.

$^1S_0 + ^1S_0$				
$^1S_0 + ^3P_0$				22401
				122
		23020	497	22523
		128		117
	23507	359	23148	508
	122		124	123
	23629	357	23272	509
	118		114	120
	23747	361	23386	503
				22883

present analysis we find that the peculiarity observed in the $\Delta F^{v'}$'s of the ground state of mercury is present also in that of cadmium. Since the bands corresponding to the 'wing' bands of mercury have not

TABLE IV.
Volkringer's Bands.

$^1S_0 + ^1S_0$ $v''=$	3	4	5	6		
$^1S_0 + ^3P_0$ $v'=0$	26680 230					
1	26910 219	746	26164 228	829 233	25335	
2	27129	737	26392 196	824	25568 171	
3			26588 200 ?	849	25739 147	
4			26788 ?	902 ?	25886 142	1042 24844
5					26028 136	
6			27027 102	863	26164* 131	1058 25106
7			27129*	834	26295	

* Used twice.

TABLE V.

$^1S_0 + ^1S_0$					
$^1S_0 + ^1P_1$					35948 182 36130 144 36274 145 36419 120 36539
			37012 96	218	36673 121 36794 375 36903 364 37012*
			37108 97	205	37012* 109 37205 193
	37302 98	97	37205 97		
	37400 98	98	37302* 98		
	37498 99	98	37400* 98		
	37597 85	99	37498* 85		
	37682				

* Used twice.

been measured in cadmium, the vibration quantum numbers could not be fixed. As in the case of mercury it is difficult to determine the heat of dissociation of the ground state. Since the data are not even as extensive as in the case of mercury it is difficult even to determine the relative position of the $v''=0$ level with respect to the $^1S_0 + ^1S_0$ electronic level as has been possible to do in the case of mercury. If the assumption

TABLE VI.
Steubing's Bands.

${}^1S_0 + {}^1S_0$ $v''=$	1	2	3	4	5				
${}^1S_0 + {}^1P_1$			47373 188	801?	46572 214				
		48035	479	47556	770?	46786			
		162		177		213			
		48197	464	47733	734	46999	857	46142	
		161		148		182		227?	
	48650	292	48358	477	47881	700	47181	812	46369
	138		146		154				
	48788	284	48504	469	48035*				
	126		146						
	48914	264	48650						

* Used twice.

is made that the similarity in the structure of the bands of mercury and cadmium extends to all the details, we may estimate the difference between $v''=0$ and the ${}^1S_0 + {}^1S_0$ electronic level as 0.18 volt while it is 0.25 volt in mercury. Because of the difficulty involved in the measurement of these diffuse and weak bands great accuracy is not found. At the present stage, therefore, the attempt to draw potential curves for the various states has been abandoned. Five of the bands recorded by Hamada could not be fitted into the present analysis so that their vibration quantum numbers relative to the rest could not be fixed. The numerical analysis shows, however, that the peculiar behaviour of the ground state in having ΔF 's which increase at first and then decrease, and in having a large anharmonic term is common to both mercury and cadmium. When this behaviour is explained, the structure of the band spectra of these metals will be better understood.