

## Physics of silver halides and their applications

D N BOSE and P A GOVINDACHARYULU\*

Materials Science Centre, Indian Institute of Technology, Kharagpur 721 302, India

\* Electrical Communication Engineering Department, Indian Institute of Science, Bangalore 560 012, India

MS received 6 May 1980 ; revised 19 July 1980

**Abstract.** Silver halides are notable for exhibiting both ionic and electronic conductivity which is responsible for their widespread use in photography. Due to differences in ionicity and bonding, their properties vary from the highly ionic silver fluoride to the more covalent silver iodide. These fundamental properties such as defect formation energies, carrier mobilities and band-structures, particularly of silver iodide, have only recently been determined. One interesting finding is the presence of high surface fields, caused by differences between interstitial and vacancy formation energies, which results in separation of photo-generated electron hole pairs. This throws new light on the photographic process and on the suitability of silver bromide as a photographic material.  $\alpha$ -AgI is also the first fast ion conductor to be discovered. Structural instability, low defect formation energy and optimum cation size are responsible for this phenomenon but existing models fail to explain all the experimental results. Recent microwave conductivity and neutron scattering experiments which provide fresh insight are discussed.

**Keywords.** Silver halides ; ionic conductivity ; defect energy ; carrier mobility ; surface field ; photography ; fast ion conduction.

### 1. Introduction

Although silver halides have long played an important role in the development of our understanding of crystalline solids, some of their fundamental properties are still in the process of being unravelled. Often overshadowed by the more famous group of alkali halides, the silver halides have nevertheless also been model solids for experimenting and testing theories, specially of the defect solid state (Brown 1967). The silver salts are notable for their unique combination of ionic as well as electronic conductivity which has given them their prominence in photography (Neblette 1964). Another intriguing phenomenon first discovered in silver iodide, fast ion conduction (Mahan and Roth 1976) is of great theoretical interest and is finding applications in solid state batteries. This paper will attempt to relate some of the fundamental ionic and electronic properties of these halides, as found recently, with these applications.

## 2. Structure

There are interesting similarities as well as differences in the family of silver halides starting with their crystal structure. While silver fluoride, chloride and bromide all possess the rock-salt structure (figure 1a) with ions having a 6-fold coordination, silver iodide in its stable  $\beta$ -phase has the hexagonal wurzite structure with 4-fold coordination (figure 1b). Earlier controversy has been resolved with the proof

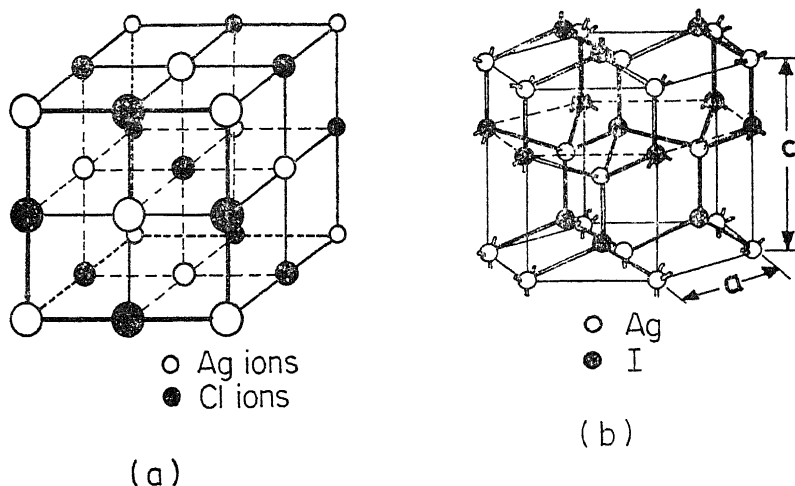


Figure 1. (a) Crystal structure of AgCl. (b) Crystal structure of AgI.

that the  $\gamma$ -phase of AgI which has the zinc-blende structure closely related to the wurzite structure and with the same coordination number, is meta-stable. At a pressure of 4 K bars  $\beta$ -AgI transforms into the rock-salt structure. At a temperature of 147° C it undergoes a phase transformation into  $\alpha$ -AgI, the first fast ion conductor to be discovered, in which the iodine atoms occupy a b.c.c. structure while the two silver ions, being disordered and statistically distributed among 12 stable sites, are responsible for the high ionic conductivity.

The difference in structure between the silver halides has been explained very successfully by the theory of ionicity and bonding (Phillips 1973). This predicts a 6-fold coordination for ionicity greater than 0.785 on the Phillips scale and a 4-fold coordination for lower ionicity and greater covalence. It has been shown that the ionicity of AgCl, AgBr and AgI are 0.856, 0.850 and 0.770 respectively, thus explaining the difference in crystal structures and providing a good test of the theory (table 1).

## 3. Ionic conductivity

The classical experiments of Tubandt in 1913 revealed that silver iodide exhibited cationic conductivity. Since then it has been established that just as alkali halides exhibit Schottky-type disorder, the silver halides have predominantly Frenkel defects, comprising of interstitial silver ions and silver ion vacancies. It is the optimum size of the silver ion and its electronic structure which is responsible for its high

Table 1. Structure and defect properties of silver halides.

Com- pound	Ionicity	Structure	$h_F$ (eV)	$W_i$ (eV)	$W_v$ (eV)	$e\phi$ (eV)	$h_m^+$ (eV)	$h_m^-$ (eV)
AgCl	0.856	Rock-salt	1.44	0.97	0.47	-0.25	0.145	0.350
AgBr	0.850	Rock-salt	1.06	0.21	0.85	0.32	0.149	0.384
AgI	0.770	Wurzite	0.64	0.20	0.44	0.12	0.41, 0.38 (   c) (   a)	0.50 0.52 (   c) (   a)

mobility. Measurements of ionic conductivity over a wide range of temperatures and with different dopants have been used for determining the defect formation and migration energies (Barr and Lidiard 1970). As is well-known, at relatively high temperatures, the conductivity is given by

$$\sigma T = A_1 \exp(-E_1/kT). \quad (1)$$

When conduction is due to one type of charge carrier,

$$E_1 = (h_F + h_m)/2,$$

where  $h_F$  is the formation energy of the Frenkel defect and  $h_m$  is the migration energy of the interstitial and  $A_1$  is a constant dependent on the jump distance and vibrational frequency.

At lower temperatures where conductivity is determined by frozen-in defects already present,

$$\sigma T = A_2 \exp(-E_2/kT), \quad (2)$$

where  $E_2 = h_m$  is the migration energy of the mobile defect already present. At still lower temperatures, an activation energy  $E = h_m + \frac{1}{2}h_a$ , where  $h_a$  is the association energy of defects, may occur but will not be discussed further here.

Such measurements had been used to find  $h_F$  and  $h_m$  for AgCl and AgBr as given in table 1 and found in many texts in solid-state physics (Brown 1967). Here  $h_m^+$  and  $h_m^-$  are the activation energies of migration of positive and negative ions respectively. These are determined, for example, by introducing divalent cations to replace monovalent silver which thus increases the concentration of negatively charged vacancies. Calculations have also been carried out to determine  $h_F$  and  $h_m$  assuming various microscopic models for ionic motion. A comparison between theoretical and experimental values of defect formation and migration energies for AgCl and also experimental values for AgBr (Friauf 1962) are given in table 2. It is seen that for  $h_F$  there is reasonable agreement between theory and experiment in the case of AgCl. It is also evident that the collinear interstitial jump mechanism has by far the lowest activation energy for both AgCl and AgBr while direct interstitial jump requires prohibitively high energies. It has thus been possible to establish by comparison with the results of diffusion

Table 2. Comparison between theoretical and experimental defect energies.

Formation and migration energies (eV)	AgCl (Expt.)	AgCl (Theory)	AgBr (Expt.)
Frenkel defect formation ( $h_F$ )	1.69	1.76	1.27
Vacancy migration ( $h_m^-$ )	0.33	1.2	0.36
Interstitialcy (average)	0.15		0.15
Collinear interstitialcy		-0.1 to 0.4	0.078
Non-collinear interstitialcy		0.8	0.23
Direct interstitial		3.0	

experiments that the silver ions move by interstitialcy or indirect interstitial mechanism and not by direct jumps between interstitial sites. Such calculations have yet to be carried out for AgI.

Comparable experimental data on AgI were also contradictory and scarce. One reason is that the phase transition at 147° C made it difficult to grow quality single crystals from the melt as for AgCl and AgBr. Further, results on polycrystalline AgI were ambiguous due to the coexistence of the  $\beta$  and  $\gamma$  phases in varying proportions, depending on the method of preparation.

The ionic conductivity in the anisotropic  $\beta$ -phase posed an interesting problem. In a hexagonal crystal there will be two independent components of conductivity, along and perpendicular to the  $c$ -axis. Thus the activation energies will be  $E_{1c}$ ,  $E_{2c}$  and  $E_{1a}$  and  $E_{2a}$  which could all be different. However,  $h_F$ , the Frenkel defect formation energy is expected to be a bulk property independent of direction of measurement. Thus it is expected from equations (1) and (2) that  $h_F = 2(E_{1c} - E_{2c}) = (E_{1a} - E_{2a})$ . However Cochrane and Fletcher (1971) who first performed detailed anisotropy experiments found that this did not follow from their results.

The problem of single crystal growth was partly solved by using the gel method near room temperature. Comparatively small pyramids and platelets were grown by Cochrane and Fletcher, Suri and later by Govindacharyulu and Bose. However the  $c$  dimension was always much less than the other dimensions and it was realised that non-uniform space-charge effects may have an important effect on the measurements. Although electronic space-charge is common place in semiconductors, ionic space-charge is also readily observable in insulating polar crystals and has been discussed theoretically by Kliever (1966). It arises from the fact that at the surface of such a crystal positive and negatively charged defects, interstitials and vacancies can and do occur separately and not necessarily in pairs. Thus the Frenkel defect energy can be divided into  $h_F = h_i + h_v$ , the energies of interstitial and vacancy formation respectively. The one with higher formation energy will predominate at the surface as fixed charge with a compensating distri-

buted space-charge region of the opposite type following immediately below. Thus for crystals thin in the  $c$  direction, measurements along the  $a$  direction would give  $E_{1a} - E_{2a} = h_i$  if  $h_i < h_v$ . The fact that positively charged species dominated at the surface space-charge region was shown by Bardeen-Brattain type ambient experiments with ozone and water-vapour. It was thus possible for the first time (Govindacharyulu *et al* 1978), using the anisotropy of AgI, to determine that the interstitial and vacancy formation energies are  $h_i = 0.20$  eV and  $h_v = 0.44$  eV respectively giving  $h_F = h_i + h_v = 0.64$  eV. Values for AgCl and AgBr are given in table 1 for comparison, but have been determined in a different manner using measurements on thin films.

From the values of  $h_F$  for AgBr and AgI (table 1) it follows that the concentration of Frenkel defects at  $300^\circ$  K is much higher in AgI than in AgBr,  $2.9 \times 10^{18}/\text{cm}^3$  compared to  $8 \times 10^{14}/\text{cm}^3$ . However, due to much lower  $h_m^+$ , ionic mobility is much lower in AgI,  $4.0 \times 10^{-7} \text{ cm}^2/\text{V sec}$  for silver interstitials compared to  $4 \times 10^{-3} \text{ cm}^2/\text{V sec}$  for AgBr.

#### 4. Surface properties

An immediate consequence of the surface space-charge in ionic crystals is that considerable fields are present and the surface potential, which deviates from the bulk, is given by  $e\phi = \frac{1}{2}(W_o - W_i)$ . Impurities may thus segregate at the surface depending on their charge as found recently by Bartzold (1975). Values of surface potential for the three common halides are also given in table 1. It is apparent that (i) surface fields are largest in AgBr and the fixed negative surface charge is neutralised by a positive space-charge region of interstitials, the net field being such as that to impel electrons into the bulk while holes are trapped at the surface and (ii) due to increased band-bending the net concentration of interstitials at the surface is at least 2 orders of magnitude larger in AgBr than in AgI. These conclusions have important implications for the photographic process.

#### 5. Electronic and optical properties

Distinguishing between electronic and ionic conductivity is a long-standing problem in many solids including silver halides. Ionic conduction involves mass transfer and hence correlation with diffusion studies has been useful in determining ionic properties. Electronic conductivity in these materials cannot be studied by conventional resistivity, Hall effect and mobility measurements due to ionic polarisation and charge motion, which reduces the field within the solid setting up fields only across the electrodes. Two methods have been used to circumvent this problem: (i) Wagner cell measurements (Wagner and Wagner 1957) utilising an asymmetric d.c. cell with one reversible and one blocking contact such as Ag/AgI/C. The steady-state current-voltage relation can be used to determine the electronic (or hole) component of conduction in a predominantly ionic conductor; (ii) pulsed field measurements in which the duration of the applied field is less than the time required for ionic polarisation to develop. This technique was developed by Brown at the University of Illinois to measure drift mobility in AgCl and AgBr and later used by Govindacharyulu and Bose (1977) to determine drift mobility of holes in  $\beta$ -AgI. The results are given in table 3.

Table 3. Mobility of charge carriers and band-gaps in silver halides.

Com- pound	$\mu_{nc}$	$\mu_{pc}$	$\mu_{ng}$	$\mu_{pg}$	$\mu_{Ag^+}$	$\mu_{Ag^-}$	$E_g$ (eV)	
							direct	Indirect
AgCl	50	$2 \times 10^{-4}$			$5 \times 10^{-3}$	$6 \times 10^{-6}$	5.15	3.25
AgBr	75	1	$2 \times 10^{-1}$	$10^{-3}$	$4 \times 10^{-3}$	$4 \times 10^{-6}$	4.29	2.69
AgI	30	12			$4 \times 10^{-7}$ (   c)	$2 \times 10^{-8}$ (   c)	2.83	

$\mu_{nc}$ ,  $\mu_{pc}$  electron and hole mobilities in crystals.  $\mu_{ng}$ ,  $\mu_{pg}$  electron and hole mobilities in grains.  $\mu_{Ag^+}$ ,  $\mu_{Ag^-}$  interstitial and vacancy mobilities in crystals. All values in  $\text{cm}^2/\text{Vsec}$  at  $300^\circ \text{K}$ .

In the experiment shown in figure 2a a pulse of absorbed light or electrons is incident on one surface of a crystal while a d.c. field is applied across it in synchronism. The arrival of the generated carriers at the second electrode is measured on an oscilloscope as in the well-known Shockley-Haynes drift mobility experiment. The drift mobility is given by  $\mu = L^2/V\tau_t$ , where  $L$  is the length of the crystal,  $V$  is the applied voltage and  $\tau_t$  is the transit time of the carrier. One notable distinction between the Shockley-Haynes and the Brown experiments is that while only the minority carrier mobility can be measured by the first method in the latter, the

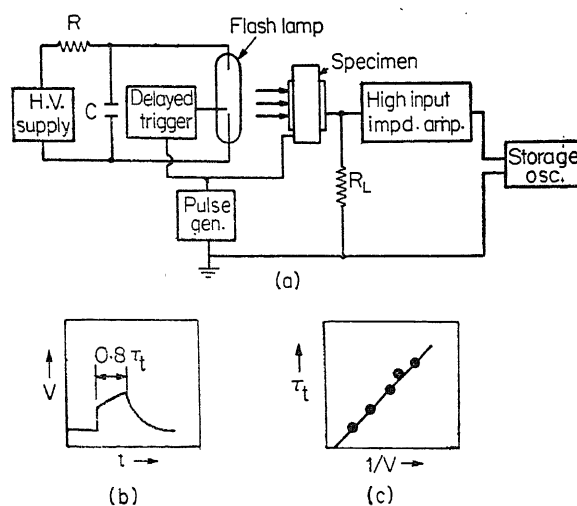


Figure 2. (a) Experiment for drift mobility measurement in ionic materials. (b) Output voltage vs. time, giving transit time. (c) Transit time  $\tau_t$  vs.  $1/V$  where  $V$  = applied drift voltage.

drift mobility of *both* types of carriers can be measured by merely changing the polarity of the applied drift field. This is because in this case the transit-time is less than the dielectric relaxation time and space-charge neutrality conditions are not maintained. In  $\beta$ -AgI for example, the drift mobility of electrons and holes at 300° K have thus been found to be 30 cm<sup>2</sup>/V sec and 12 cm<sup>2</sup>/V sec respectively (Bose and Govindacharyulu 1976). In all the silver halides the mobility is controlled by shallow traps at lower temperatures, the effective mobility decreasing with decrease in temperature due to carriers remaining immobilised at these localised states. The identity of these traps is a subject for continuing investigation. At higher temperatures the mobility goes through a maximum and then decreases as usual due to phonon scattering.

Of crucial importance as far as optical properties are concerned is the nature of the electronic band-structure. In semiconductor devices it has been realised over the last 20 years that the operation of devices such as lasers, LEDs or Gunn diodes is directly dependent on the band-structure of a material. As is well known materials can be divided into direct band-gap materials like GaAs, InSb, CdS for which the minimum of the conduction band as well as the maximum of the valence band have the wave vector  $k = 0$  (figure 3a) and (ii) indirect band-gap materials like silicon and germanium (figure 3c) in which the conduction band minimum is displaced from  $k = 0$ . Optical absorption and emission involve vertical electronic transition with no change in  $k$  vector since a photon has negligible momentum on this scale. Thus while direct band-gap materials have high probability of absorption or emission of radiation, indirect band-gap materials require the participation of a phonon which is either absorbed or emitted, to conserve momentum.

It had been shown by Bassani *et al* (1965) from band-structure calculations that AgCl and AgBr represent another type of indirect band-gap material (figure 3b) in which the maximum of the valence band is displaced from  $k = 0$  although the minimum of the conduction band is still at this position. However, considerable controversy existed as to the nature of the band-gap in  $\beta$ -AgI. Govindacharyulu and Bose (1979) have recently shown by measurements of photo-conductivity excitation spectra and anisotropy that  $\beta$ -AgI is a direct band-gap material very much like CdS which it resembles in structure. This is supported by pseudo-

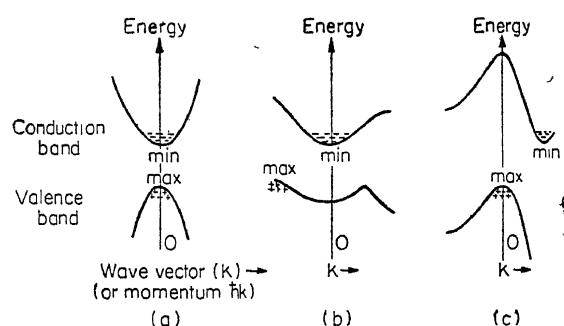


Figure 3. Energy band diagram of three types of semiconductors. (a) Direct band gap like GaAs, InSb, AgI, etc. (b) Indirect band gap like AgCl, AgBr, etc. (c) Indirect band gap like Si, Ge, GaP, etc.

potential calculations of band-structure by Smith (1976). Thus while the recombination of excess carriers in  $\beta$ -AgI will be rapid, it will be comparatively slow in AgCl and AgBr. In fact the anomaly of AgBr having a smaller band-gap of 2.69 eV compared with 2.83 eV of AgI is now explained, since the former is indirect while the latter is direct, and the direct band-gaps are found to scale according to Welker's rule. It is thus expected that AgI could be made to exhibit stimulated emission at low temperatures. The study of electron-hole droplet formation in polar materials is also of current interest and the theory suggests the possibilities of stable droplets in AgBr (Beni and Rice 1978) while AgI is yet to be examined.

The peculiar valence band structure of the silver halides is due to the close proximity of bands originating on the halogen  $p$  levels and the silver  $d$  states. Because the rock salt structure of AgCl and AgBr has inversion symmetry, these states do not mix at  $\Gamma$ , the centre of the Brillouin zone, but away from  $\Gamma$  the repelling of bands causes the uppermost levels to bend upwards rapidly. This leads to widely separated direct and indirect gaps. The wurzite structure of  $\beta$ -AgI without inversion symmetry does not force the electronic states to form an indirect edge in  $\beta$ -AgI which is thus unique among the silver halides. This is borne out by the photoemission experiments of Bauer and Spicer (1976). The detailed band-structures are given in figure 4.

## 6. Photographic properties

Interestingly enough the first photosensitive material used in the famous Daguerrotype process of 1839 by Louis Daguerre was silver iodide. It was discarded for AgBr shortly after because of the latter's higher sensitivity and wider spectral response. Although microscopic theories of the photographic process were developed by Gurney and Mott (1938) to be modified by Mitchell in 1946, the difference in material properties was not understood until recently. Wei and Brown (1973) remarked on the fortuitous combination of electronic and ionic

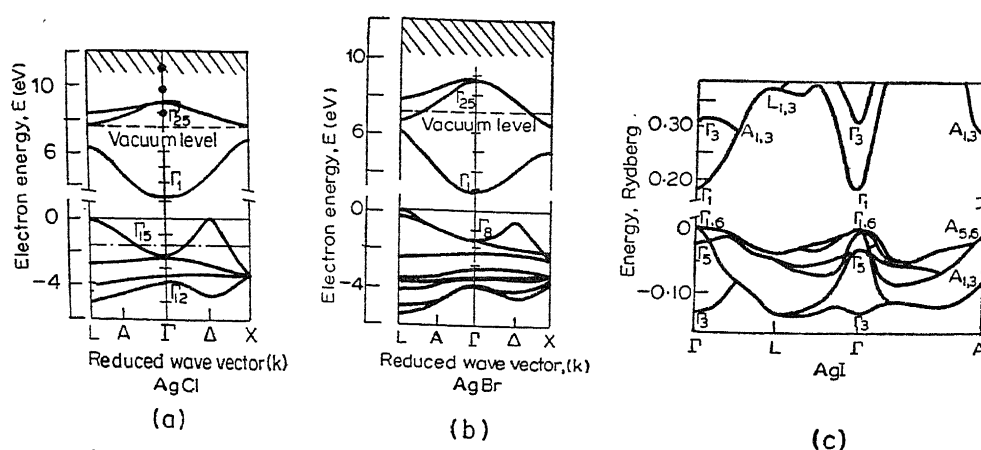


Figure 4. (a) Band-structure of AgCl. (b) Band-structure of AgBr. (c) Band-structure of AgI.



properties which make silver bromide still pre-eminent among photographic materials.

According to the Mott-Gurney theory when a photon is incident on silver bromide, it is absorbed if its energy  $h\nu > E_g$ , the band-gap of the semiconductor, creating an electron-hole pair. The spectral response of pure AgBr extends only to wavelength of  $0.46 \mu\text{m}$  and is widened by using suitable sensitising dyes. The hole is trapped near the surface of the grain which is typically  $0.02$  to  $4 \mu$  in size, while the electron migrates into the grain aided by the surface field to be trapped by a previously neutral electron trap called a sensitivity speck. This may consist of a sulphur atom. Being charged negatively on acquiring an electron, the speck attracts a mobile positive charge, which, in AgBr, is a positive silver ion. The speck thus collects a neutral silver atom which forms a latent image capable of being developed. In practice it is found that about 4 photons are necessary to create a developable latent image.

The important differences between silver bromide and iodide in the photographic process are the following (i) the indirect band-gap of the former reduces the probability of direct electron-hole recombination (ii) higher surface field in AgBr which aids electron motion, (iii) lower activation energy for migration  $0.14 \text{ eV}$  for AgBr compared to  $0.38\text{--}0.41 \text{ eV}$  for AgI leads to much higher silver interstitial mobility,  $4 \times 10^{-3} \text{ cm}^2/\text{V sec}$  and  $4 \times 10^{-7} \text{ cm}^2/\text{V sec}$  respectively, and (iv) the difference in hole mobility  $\mu_p$ ,  $1 \text{ cm}^2/\text{V sec}$  for AgBr compared with  $12 \text{ cm}^2/\text{V sec}$  for AgI in large single crystals as given in table 3, the ratio being similar in emulsion grains. Since  $h_F = 0.64 \text{ eV}$  for AgI compared with  $1.06 \text{ eV}$  for AgBr, the concentration of Frenkel defects at  $300^\circ \text{ K}$  are  $2.9 \times 10^{18}/\text{cm}^3$  and  $8 \times 10^{14}/\text{cm}^3$  respectively. Assuming uniform distribution of these defects and the interstitial mobilities mentioned above, the time for capture of an interstitial in AgI and AgBr are  $5.7 \times 10^{-4} \text{ sec}$  and  $2.7 \times 10^{-6} \text{ sec}$  respectively for an electric field of  $10^3 \text{ V/cm}$ . Thus interstitial capture by trapped electrons is favoured in AgBr while hole capture is more probable in AgI. The serendipitous combination of an indirect gap which necessitates the participation of a phonon in recombination, the asymmetry between electron and hole mobilities and the high interstitial silver ion mobility transpire to give AgBr its unique photographic properties.

### 7. Fast ion conduction

As mentioned earlier  $\beta$ -AgI undergoes a phase transition at  $147^\circ \text{ C}$  into  $\alpha$ -AgI which is a fast ion conductor. AgCl and AgBr also undergo similar transition at  $455^\circ \text{ C}$  and  $425^\circ \text{ C}$  respectively. The high temperature phases have conductivities of the order of  $1 (\Omega \text{ cm})^{-1}$ , as high as liquid electrolytes. At the  $\beta \rightarrow \alpha$  transition of AgI the conductivity increases by more than 3 orders of magnitude, and remains almost constant decreasing somewhat on melting. The activation energy in  $\alpha$ -AgI is only  $0.10 \text{ eV}$  and is typical of a series of derived compounds such as  $\text{RbAg}_4\text{I}_5$  which is a fast ion conductor at  $300^\circ \text{ K}$ . CuBr and CuI also display similar transitions and fast ion conduction in the high temperature phase. The phenomenon has been reviewed in detail recently (Mahan and Roth 1976) and hence

will only be briefly discussed here. Theoretical models used have been (i) hopping conduction (ii) cooperative interaction model and (iii) free ion like model but none have satisfactorily explained all the experimental data. For example, all hopping conduction models predict an increase in conductivity with frequency unlike the decrease observed experimentally for  $\alpha$ -AgI at microwave frequencies by Funke (1976). The free-ion like model predicts a decrease with frequency but cannot explain conductivity peaks obtained in the microwave range.

An interesting calculation based on motion of cations along tunnels in the  $\langle 100 \rangle$  direction of the b.c.c. anion lattice of  $\alpha$ -AgI has been developed by Flygare and Huggins (1973). The activation energy for motion has been calculated as a function of cationic size and shows a minimum for cations of radii  $0.83 \text{ \AA}$  which is very close to the radius of the highly mobile silver ion. Smaller ions face larger energy barriers due to electrostatic attraction while larger cations are hindered by the repulsive term in potential energy.

Two experimental techniques that have provided interesting insight are the microwave measurements of Funke (1976) which show that the jump distances are of the order of  $5 \text{ \AA}$  and that the mean time of flight is of the same order as residence time at sites  $\approx 10^{-11}$  sec. Thus cationic motion in AgI is intermediate between diffusion and hopping mechanisms.

The other technique is neutron scattering to determine the lattice dynamics of  $\beta$  and  $\alpha$ -AgI. The results of Bührer *et al* (1978) show the presence of very low lying dispersionless  $T_g$  modes. These provide a strong contribution to thermal motion because they have a high density of states. The r.m.s. displacement of Ag ions is thus quite large increasing from  $0.164 \text{ \AA}$  in AgCl,  $0.173 \text{ \AA}$  in AgBr to  $0.292 \text{ \AA}$  in AgI, thus favouring promotion of a cation to an interstitial site. However no mode softening has been observed between  $80^\circ \text{ K}$  and  $295^\circ \text{ K}$  and the zone-centre low frequency mode is stable up to the transition point. The phase transition thus appears to be driven by the entropy change due to cation disordering, the transition temperature being given by  $T_c = \Delta H / \Delta S$ , where  $\Delta H$  is the change in enthalpy.

A number of solids related to AgI also exhibit similar properties, among them  $\text{Ag}_3\text{SI}$ ,  $\text{Ag}_2\text{HgI}_4$ . The most interesting from the point of view of applications as solid electrolytes in batteries (Funke 1976) are the group of isostructural compounds  $\text{MAg}_4\text{I}_5$  where ( $M = \text{NH}_4$ , K or Rb). These compounds are virtually pure ionic conductors at room temperature with conductivities as high as  $0.2 \Omega^{-1} \text{ cm}^{-1}$  and have been used in cells such as  $\text{Ag/RbAg}_4\text{I}_5/\text{RbI}_3$ , C. The main drawback of these batteries is their low output voltage of  $0.6 \text{ V}$  and hence low energy densities. The search is thus continuing for better materials such as those based on lithium ion conductors.

It is thus evident that the fundamental properties now being determined such as band-structure, defect formation and migration energies of silver halides are of considerable importance in determining their applications. Better understanding of this relatively simple class of solids should thus help the development of new materials, for example less expensive ones for silver-less photography in the infra-red region and other materials for high-energy density solid-state batteries.

## References

- Baetzold R C 1975 *Appl. Phys. Lett.* **12** 709  
Barr L W and Lidiard A B 1970 *Physical Chemistry* **10** 152  
Bassani F, Knox R S and Fowler W B 1965 *Phys. Rev.* **137** A1217  
Bauer R S and Spicer W E 1976 *Phys. Rev.* **B14** 10, 4539  
Beni G and Rice T M 1978 *Phys. Rev.* **B18** 2, 768  
Bose D N and Govindacharyulu P A 1976 *Physics of Semiconductors* (ed.) F G Fumi (Amsterdam : North Holland) p. 1180  
Brown F C 1967 *Physics of Solids* (New York : W A Benjamin, Inc.)  
Buhner W, Nicklow R M and Bruesch P 1978 *Phys. Rev.* **B17** 8, 3362  
Cochrane G and Fletcher N H 1971 *J. Phys. Chem. Solids* **32** 2557  
Flygare W H and Huggins R A 1973 *J. Phys. Chem. Solids* **34** 1199  
Friauf R T 1962 *J. Phys. Chem. Solids* **66** 2376  
Funke K 1976 *Superionic Conductors* (New York : Plenum Pub. Co.)  
Funke K 1976 *Prog. in S.S. Chemistry* **11** 345  
Govindacharyulu P A and Bose D N 1977 *J. Appl. Phys.* **48** 3, 1981  
Govindacharyulu P A, Bose D N and Suri S K 1978 *J. Phys. Chem. Solids* **39** 961  
Govindacharyulu P A and Bose D N 1979 *Phys. Rev.* **137** 6532  
Gurney R W and Mott N F 1938 *Proc. R. Soc.* **A164** 151; 1964 *Electronic processes in ionic crystals* (Dover: New York)  
Kliewer K L 1966 *J. Phys. Chem. Solids* **27** 705, 719  
Liang C C 1974 "Solid State Batteries," *Appl. Sol. St. Science* **4** 95  
Mahan G D and Roth W L (eds.) 1976 *Superionic Conductors* (New York : Plenum)  
Neblette C B 1964 *Photography* (New York : Van Nostrand)  
Phillips J C 1973 *Bonds and bands in semiconductors* (New York : Academic Press)  
Smith P V 1976 *J. Phys. Chem. Solids* p. 37765  
Wagner J B and Wagner C 1957 *J. Chem. Phys.* **26** 1597  
Wei J S and Brown F C 1973 *Photog. Sci. and Engg.* **17** 2, 197