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SUBDWARFS AND ABSORPTION DUE TO QUASI-MOLECULAR HYDROGEN

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A quasi-molecule of two colliding hydrogen atoms in the unstable ${}^3\Sigma_u^+$ state goes to the excited ${}^3\Sigma_g^+$ state, with the absorption of light. The free-bound absorption continuum produced by quasi- H_2 was considered as early as 1947 by Wildt¹ and found to be unimportant in late-type stars. Recent work by Erkovich² suggested that this source of opacity may in fact be astrophysically important (see *e.g.* Zwaan³). The work of Soloman⁴ and of Doyle⁵ showed, however, that Erkovich cross-sections have been overestimated by about two orders of magnitude. Hence, the absorption due to quasi- H_2 is negligible compared to H^- and H in stars with chemical composition similar to that in the solar photosphere. However, Varsavsky⁶ has recently suggested that absorption due to quasi- H_2 may not be negligible in subdwarfs, which have rather low metal abundance. We have tried in this note to examine the role of quasi- H_2 in metal-deficient stars.

We have assumed that the relative abundance of metals is similar to that in the solar photosphere⁷, with the ratio of metal to hydrogen being one-hundredth that in the Sun, and the ratio of helium to hydrogen, by number, being 0.0625. We have used the absorption cross-sections for quasi- H_2 as given by Soloman⁴ and have followed Vardya⁸ for the continuous opacity due to H , H^- , H^-_2 , He and He^- . The relevant range in pressure for a given temperature can be obtained from the work of Swihart and Fischel⁹.

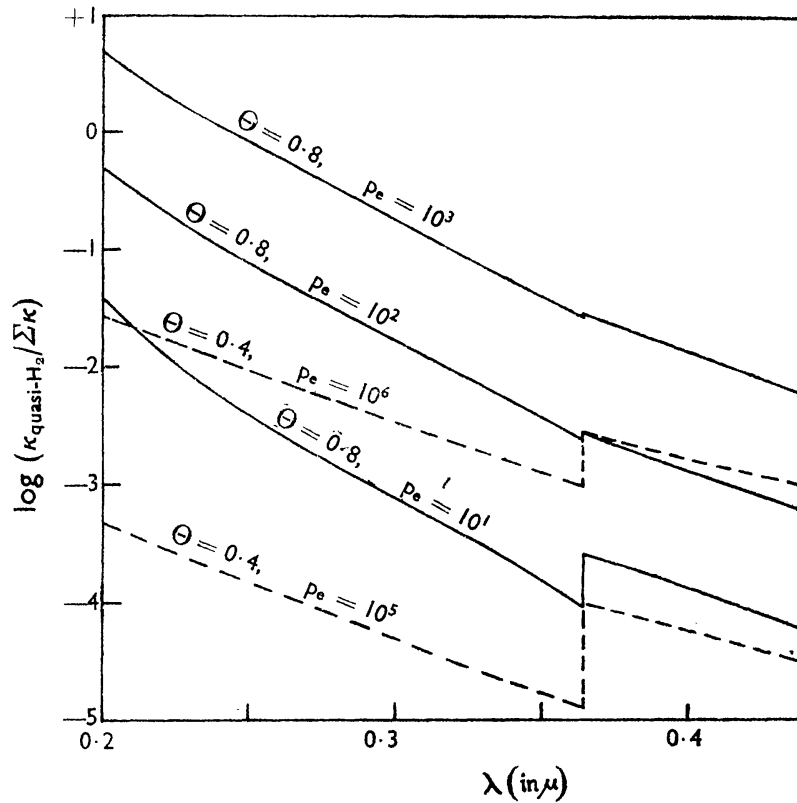


FIG. 1

Logarithm of the ratio of absorption due to quasi- H_2 to sum of other sources of continuous opacity as a function of wavelength.

Table I gives the values of $\theta = 5040.39/T$ (where T is the absolute temperature), electron pressure, p_e , and the corresponding gas pressure, P_g , at which comparisons have been made of the relative importance of quasi- H_2 absorption to other sources of continuous opacity. (P_g and p_e are in dyn/cm².)

TABLE I

θ	p_e	P_g	θ	p_e	P_g
0.4	10^5	7.0×10^5	1.2	10^{-1}	1.5×10^5
0.4	10^6	5.1×10^7	1.2	10^1	9.9×10^7
0.8	10^1	7.6×10^3	1.6	10^{-2}	1.3×10^5
0.8	10^2	7.6×10^5	1.6	10^0	4.6×10^7
0.8	10^3	8.5×10^7			

Fig. 1 is a plot of $\log_{10} (\kappa_{\text{quasi-}H_2} / \Sigma \kappa)$ versus λ (wavelength in μ) for $\theta = 0.4$ and 0.8 and Fig. 2 is a similar plot for $\theta = 1.2$ and 1.6. Here $\Sigma \kappa$ is the sum of the continuous absorption coefficients due to H , H^- , H_2^- , He and He^- . These curves indicate that the importance of quasi- H_2 absorption, compared to other sources of opacity in the ultra-violet region of subdwarfs, increases as one proceeds from high temperatures to low temperatures, reaches a maximum around $\theta \simeq 1.2$, and then its importance decreases. The effect of metal abundance, for a given (θ, P_g) , is rather small at high temperatures, because there most of the electrons are donated by hydrogen and atomic hydrogen is the main source of opacity. In fact, the ratio $\kappa_{\text{quasi-}H_2} / \kappa_H$, where

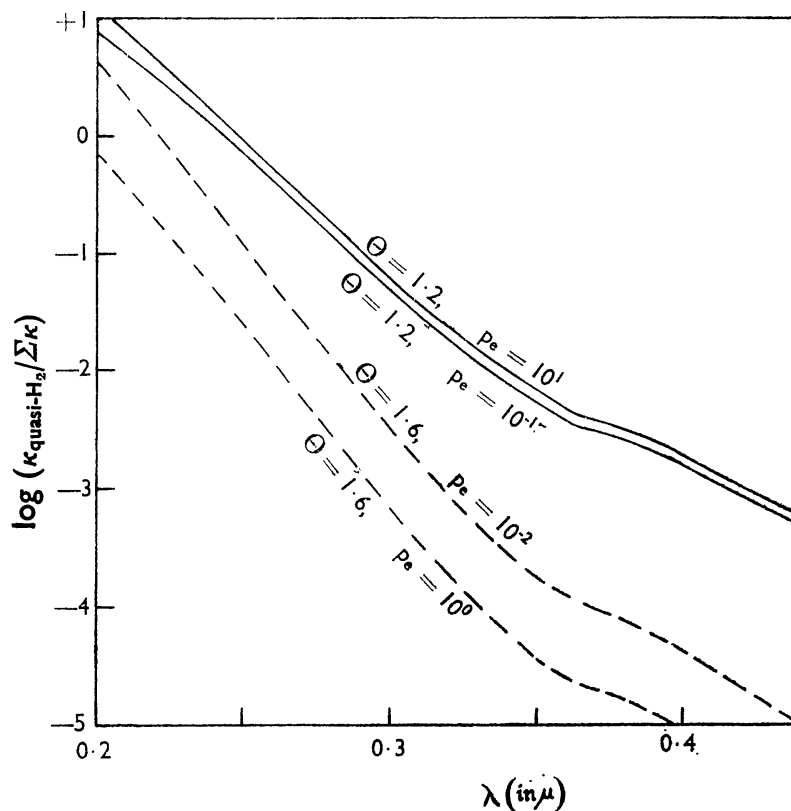


FIG. 2

Logarithm of the ratio of absorption due to quasi- H_2 to sum of other sources of continuous opacity as a function of wavelength.

κ_H is the absorption coefficient due to atomic hydrogen, is independent of metal abundance. At lower temperatures, where the metals donate most of the electrons, the importance of quasi- H_2 is greatly enhanced with decrease in metal abundance, as is evident from a comparison of our figures with those of Soloman⁴. It is the increasing formation of molecular hydrogen that reduces the relative importance of quasi- H_2 absorption as one proceeds to values of θ beyond $\theta \simeq 1.2$. Needless to say, the importance of quasi- H_2 absorption decreases with decrease in hydrogen to metal ratio. Though quasi- H_2 is important in subdwarfs, it should not be forgotten that, in the ultra-violet region of these stars, Rayleigh scattering is very important in redistributing radiation.

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