

## Particle-size effects on the value of $T_c$ of $\text{MnFe}_2\text{O}_4$ : Evidence for finite-size scaling

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$^{57}\text{Fe}$  Mössbauer measurements show that the  $T_c$  of  $\langle 12 \rangle$ -nm particles of  $\text{MnFe}_2\text{O}_4$  is higher than that of the  $\langle 33 \rangle$ -nm particles, which exhibit the bulk  $T_c$  value (573 K). The  $^{57}\text{Fe}$  isomer shift as well as the internal field are, however, independent of the particle size, as is the cation distribution determined by Mn and Fe extended x-ray absorption fine structure and Mn x-ray absorption near-edge structure measurements at different temperatures. The results confirm that the higher  $T_c$  of the  $\langle 12 \rangle$ -nm particles arises from finite-size scaling.

Size-dependent properties of nanoscale particles have attracted considerable attention in recent years.<sup>1-4</sup> Of particular interest is the question of what happens to magnetic properties when the particle size is of the order of a domain. Experimental studies of small magnetic particles have by and large been limited to superparamagnetism in very small particles.<sup>5,6</sup> Tang *et al.*<sup>1</sup> have recently measured the Curie temperature of ferrimagnetic  $\text{MnFe}_2\text{O}_4$  particles and found a significant increase in  $T_c$  (relative to the bulk value) in particles slightly larger than the superparamagnetic critical size. The authors explain this interesting observation on the basis of the finite-size scaling model,<sup>7</sup> according to which the relative enhancement in  $T_c$  is related to the particle diameter raised to an exponent. This explanation has been questioned by some workers<sup>8,9</sup> who consider that particle-size-dependent changes in the cation distribution over tetrahedral and octahedral sites could as well be responsible for the increase in  $T_c$ . Considering the importance of the problem,<sup>10</sup> we have reexamined the properties of  $\text{MnFe}_2\text{O}_4$  particles employing atom-specific probes. We have probed the magnetic environment around the Fe ions by  $^{57}\text{Fe}$  Mössbauer spectroscopy, as particles of different sizes undergo the ferrimagnetic-paramagnetic transition. We have monitored the local environment of the Mn and Fe ions as a function of size as well as temperature by means of the Mn  $K$  and Fe  $K$  extended x-ray absorption fine structure (EXAFS). In addition, we have examined the Mn x-ray absorption near-edge structure (XANES) in order to study the distribution of Mn in the different valence states. The present study has shown that the cation distribution remains essentially the same in  $\text{MnFe}_2\text{O}_4$  particles of different sizes and that there is indeed a ferrimagnetic phase comprising of small ferrite particles ( $\sim 12$  nm) with a  $T_c$  higher than that of the bulk spinel.

$\text{MnFe}_2\text{O}_4$  particles of varying sizes (6–33 nm) were prepared following the procedure in the literature.<sup>11</sup> The samples were annealed in helium at 773 K for 1 h and slowly cooled to room temperature in order to ensure equilibrium cation distribution. The cation stoichiometry was checked by energy dispersive analysis of x rays and the single crystal nature of the particles was verified by

high-resolution electron microscopy. Mean sizes of the particles was obtained from the Brunauer-Emmett-Teller surface areas. Samples with mean particle diameters of 9, 12, and 33 nm were chosen for detailed studies. Mössbauer and EXAFS measurements were carried out by taking the samples in the form of wafers. For the EXAFS measurements, a Ge(220) crystal was used as the monochromator on a Rigaku spectrometer attached to a rotating anode x-ray source (RU 200B, Rigaku, Japan). Fourier transforms (FTs) of the EXAFS data were obtained with  $k_{\min} \sim 3$  and  $k_{\max} \sim 12 \text{ \AA}^{-1}$  after weighing the data by  $k^3$ .

$^{57}\text{Fe}$  Mössbauer spectra show the sextet due to the ferrimagnetic species or/and a doublet due to the para and/or superparamagnetic species, depending on the size and temperature (Fig. 1). The 300-K spectrum of the  $\langle 33 \rangle$ -nm particles show a six finger pattern indicating that all iron ions are in the magnetic phase. The internal magnetic field ( $H_{\text{int}}$ ) was found to be 463 KOe when the spectrum was fitted with a single sextet. The spectrum could also be fitted with two sextets (75:25) with  $H_{\text{int}}$  values of 456 and 484 kOe corresponding to two Fe sites in the spinel lattice, but with nearly the same isomer shift ( $\text{IS} \sim 0.4 \text{ mm s}^{-1}$ ). At 573 K ( $T_c$  of bulk  $\text{MnFe}_2\text{O}_4$ ), the  $\langle 33 \rangle$ -nm particles show a doublet in the Mössbauer spectrum along with a broad singlet. The  $\langle 9 \rangle$ -nm particles exhibit Mössbauer spectra characteristic of small superparamagnetic particles (Fig. 1), showing the ferrimagnetic sextet at very low temperatures ( $H_{\text{int}}$  490 kOe at 90 K). While there can be no sharp  $T_c$  for such a sample, ferrimagnetism does not manifest itself down to  $\sim 200$  K. However, the isomer shift in the  $\langle 9 \rangle$ -nm sample is the same ( $\sim 0.4 \text{ mm s}^{-1}$ ) as that of the  $\langle 33 \rangle$ -nm particles.

The behavior of the  $\langle 12 \rangle$ -nm sample is interesting in that the Mössbauer spectrum which is composed of a sextet (ferrimagnetic) and a doublet (superparamagnetic) at 300 K continues to show the sextet even at 573 K (Fig. 1). The sextet disappears only above 620 K, clearly indicating that the  $T_c$  of the  $\langle 12 \rangle$ -nm particles is higher than that of the  $\langle 33 \rangle$ -nm particles by  $\sim 50$  K. We present the results of the Mössbauer studies in terms of the variation of the ferrimagnetic fraction at 573 K as well as of  $T_c$

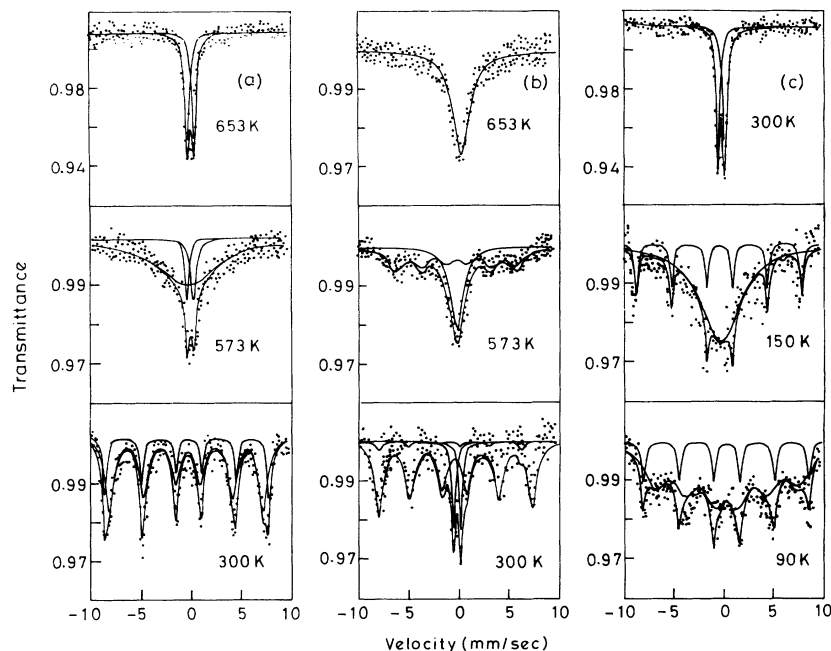


FIG. 1.  $^{57}\text{Fe}$  Mössbauer spectra of the ferrite samples at various temperatures. Particles with mean diameters of (a) 33 nm, (b) 12 nm, and (c) 9 nm.

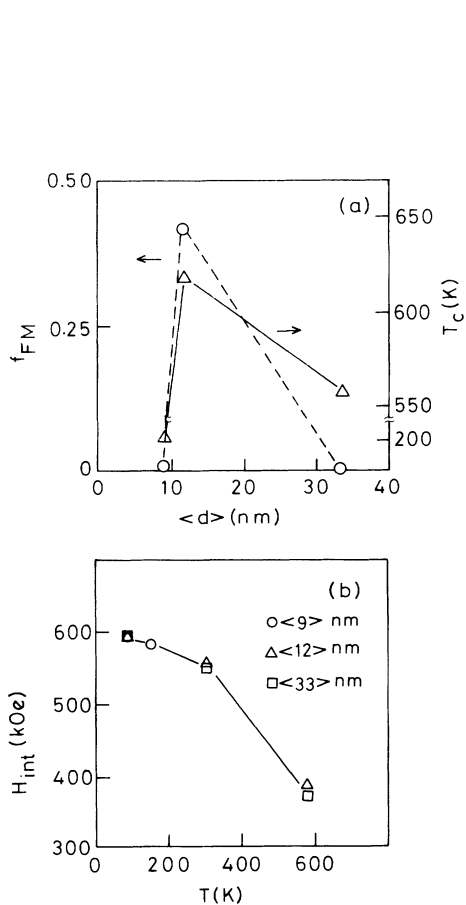


FIG. 2. (a) Variation of the fraction of ferrimagnetic particles (circles) at 573 K (bulk  $T_c$  value) and  $T_c$  of  $\text{MnFe}_2\text{O}_4$  (triangles) with particle size. In case of the  $\langle 9 \rangle$ -nm particles, we have shown  $T_c$  as the temperature at which ferrimagnetism first manifests itself in the Mössbauer spectrum. (b) Variation of the internal magnetic field,  $H_{int}$  with temperature for the  $\langle 33 \rangle$ -,  $\langle 12 \rangle$ -, and  $\langle 9 \rangle$ -nm particles.

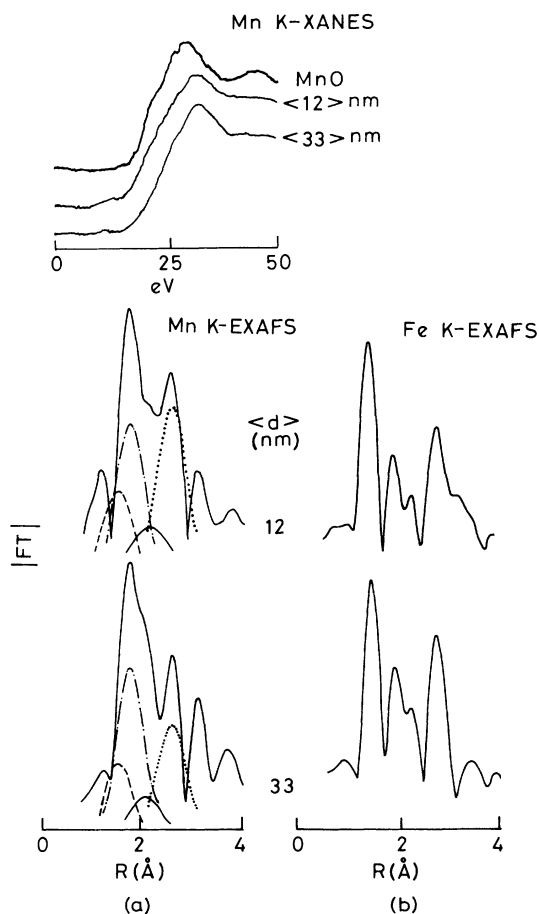


FIG. 3. Fourier transforms (FTs) of the EXAFS of the ferrite samples at room temperature: (a) Mn K EXAFS and (b) Fe K EXAFS. Fourier deconvoluted peaks corresponding to different oxygen coordinations are also shown. The peak due to metal-metal coordination is shown by the dotted line. The inset in (a) shows Mn K near-edge spectra of the ferrite samples along with that of MnO.

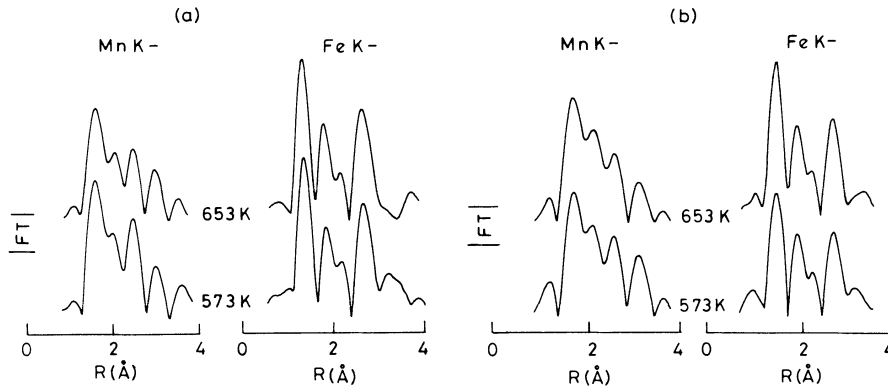


FIG. 4. Fourier transforms of Mn *K* and Fe *K* EXAFS of the ferrite samples at 573 and 653 K. (a) The  $\langle 33 \rangle$ -nm particles and (b) the  $\langle 12 \rangle$ -nm particles.

with the particle size. We show in Fig. 2(b) how the  $H_{\text{int}}$  remains essentially the same in samples of different particle sizes ( $\sim 490$ – $495$  kOe at 90 K). The near constancy of the isomer shift and  $H_{\text{int}}$  in samples of different particle sizes indicates that the cation distribution does not change with particle size. It, therefore, appears that the higher  $T_c$  of the  $\langle 12 \rangle$ -nm sample is due to finite-size scaling. By employing the finite-size scaling formula we estimate the  $T_c$  of the  $\langle 12 \rangle$ -nm particles to be 617 K, which agrees with the experimental value (620 K) found by us. The correlation length exponent is 1.42 and the amplitude is 2 nm.

In order to ensure that the increase in  $T_c$  of the  $\langle 12 \rangle$ -nm sample relative to the  $\langle 33 \rangle$ -nm sample is not due to changes in cation distribution and related factors, we have examined the local environment of Fe and Mn ions by EXAFS. The FT of the Mn *K* EXAFS of the  $\langle 33 \rangle$ -nm sample at 300 K (Fig. 3) shows the main feature at 1.8 Å with a shoulder around 2.1 Å due to oxygen coordination. By appropriate analysis of the data by the residual FT method<sup>12</sup> ( $r$  window: 1.4–2.9 Å) by employing MnO parameters, we obtain Mn–O distances of 2.0, 2.22, and 2.52 Å with coordinations of 2.0, 4.5, and 1.2. The 2.22-Å distance is due to the tetrahedral site, the other two distances arising from the John-Teller distorted octahedra. The proportion of octahedral Mn is estimated to be  $\sim 32\%$  consistent with the degree of inversion reported in the literature.<sup>13</sup> The FT of the  $\langle 12 \rangle$ -nm sample (Fig.

3) gave Mn–O distances and coordination close to those of the  $\langle 33 \rangle$ -nm sample. Mn *K* XANES is similar in ferrite samples of different sizes (see inset of Fig. 3) indicating identical average oxidation states of Mn (estimated to be 2.32) and hence the degree of inversion. FTs of the Mn *K* EXAFS of the  $\langle 33 \rangle$  and  $\langle 12 \rangle$  nm were nearly identical at 573 and 653 K (Fig. 4). The first main feature in the FTs at these temperatures was at a lower distance since the proportion of octahedral Mn was higher ( $\sim 60\%$ ) than at 300 K.

Fourier transforms of the Fe *K* EXAFS of the  $\langle 33 \rangle$ -nm sample at 300 K show distinct features due to oxygen coordination at 1.85, 2.24, and 2.70 Å (Fig. 3) which on curve-fitting analysis gave Fe–O distances of 1.88, 2.33, and 2.73 Å. The 2.33-Å distance is ascribed to the tetrahedral site and the other two to the octahedral site. The ratio of octahedral to tetrahedral Fe sites was 0.57. This ratio as well as the Fe–O distances are found to be the same in the  $\langle 12 \rangle$ -nm sample. The FTs at 573 and 653 K were also nearly identical for the  $\langle 12 \rangle$ - and  $\langle 33 \rangle$ -nm samples (Fig. 4) but the proportion of the tetrahedral species was higher than at 300 K.

The present EXAFS study establishes that the cation distribution and the associated features remain essentially the same in the  $\langle 12 \rangle$ - and  $\langle 33 \rangle$ -nm samples (Table I), thus reinforcing the suggestion that the higher  $T_c$  of the  $\langle 12 \rangle$ -nm sample is likely to be due to finite-size scaling.

TABLE I. Structural parameters of  $\text{MnFe}_2\text{O}_4$  particles.

$T$ (K)	Coordination number and distances for Mn–O				Coordination numbers and distances for Fe–O			
	$N$	$R$ (Å)	$N$	$R$ (Å)	$N$	$R$ (Å)	$N$	$R$ (Å)
300	4	1.97	4	2.00	4	1.90	4	1.88
	4	2.20	4	2.22	4	2.32	4	2.33
	2	2.54	2	2.52	2	2.74	2	2.74
573	3	2.05	3	2.05	3	1.89	4	1.90
	2	2.20	2	2.20	3	2.33	4	2.35
	3	2.53	3	2.53	3	2.74	2	2.75
653	3	2.06	4	2.06	4	1.89	4	1.89
	3	2.19	2	2.20	3	2.34	3	2.33
	3	2.50	2	2.50	2	2.73	2	2.75

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