

Exchange Interaction in Binuclear Complexes with Rare Earth and Copper Ions: A Many-Body Model Study

Indranil Rudra , C. Raghu and S. Ramasesha

Solid State and Structural Chemistry Unit, Indian Institute of Science,
Bangalore 560 012, India

ABSTRACT

We have used a many-body model Hamiltonian to study the nature of the magnetic ground state of hetero-binuclear complexes involving rare-earth and copper ions. We have taken into account all diagonal repulsions involving the rare-earth 4f and 5d orbitals and the copper 3d orbital. Besides, we have included direct exchange interaction, crystal field splitting of the rare-earth atomic levels and spin-orbit interaction in the 4f orbitals. We have identified the inter-orbital 4f repulsion, U_{ff} and crystal field parameter, Δ_f as the key parameters involved in controlling the type of exchange interaction between the rare earth 4f and copper 3d spins. We have explored the nature of the ground state in the parameter space of U_{ff} , Δ_f , spin-orbit interaction strength λ and the 4f filling n_f . We find that these systems show low-spin or high-spin ground state depending on the filling of the 4f levels of the rare-earth ion and ground state spin is critically dependent on U_{ff} and Δ_f . In case of half-filling (Gd(III)) we find a reentrant low-spin state as U_{ff} is increased, for small values of Δ_f , which explains the recently reported apparent anomalous anti-ferromagnetic behaviour of Gd(III)-radical complexes. By varying U_{ff} we also observe a switch over in the ground state spin for other fillings. We have introduced a spin-orbit coupling scheme which goes beyond L-S or j-j coupling scheme and we find that spin-orbit coupling does not significantly alter the basic picture.

1 Introduction

Electron spin pairing and the associated electron exchange between atoms is the cornerstone of chemical bond formation. Hence most of the molecules have closed shell ground state [1] and the synthesis of high spin molecules as well as molecular magnets has been a challenge. The simplest idea to synthesize a high spin molecule is to start with a high spin carrier. In this regard lanthanides, and in particular, Gd^{3+} which has a $4f^7$ configuration is an obvious choice. Mononuclear complexes of Gd^{3+} are already known. Systems with ground state spin higher than $\frac{7}{2}$ can be built up with multinuclear complexes involving Gd^{3+} and other rare-earth or transition metal ions. Studies of such heteropolymetallic

compounds have also been of great interest in modeling metalloenzymes and in understanding their magnetic properties [2]. Indeed, the interaction between two nonequivalent metal centers can lead to some unique features which are not encountered in the complexes with same kind of spin carriers.

In the last few years, there have been several studies of heteropolymetallic molecular coordination complexes of lanthanides and Cu(II) or other ions because of their luminescence and interesting magnetic properties [3]. Some of these complexes exhibit a high spin ground state corresponding to a ferromagnetic alignment of magnetic moments of the individual ions, a fact which is important for designing building blocks for molecular magnets. Gatteschi and coworkers [4–6] found that in case of Gd(III)-Cu(II) binuclear and Cu(II)-Gd(III)-Cu(II) trinuclear complexes, the Gd(III)-Cu(II) interaction is ferromagnetic in nature, independent of the structure of the complex or the ligand molecule. Subsequently, a systematic effort to synthesize and study binuclear systems between Cu(II) and other rare-earth ions was undertaken [7–9]. These studies indicate a general trend for the ground state spin of the complex as we traverse the lanthanide series. The interaction between Cu(II)-Ln(III) is ferromagnetic for Gd(III) and other rare-earth ions with greater than half-filled $4f$ shell such as Dy(III) and Er(III). Below half filling, this interaction is antiferromagnetic and leads to low-spin ground state for systems such as Cu(II)-Ce(III) and Cu(II)-Eu(III) [7]. The situation remains the same when the Cu(II) is replaced by organic ligands like nitronyl nitroxide triazole which are essentially spin-1/2 carriers.

As a possible mechanism for the observed ferromagnetic alignment of the spin of the rare-earth ion with the spin of the copper ion, it was pointed out [10] that the stabilization of the high spin state arises from the coupling between the ground state configuration and the excited state configuration arising out of the transfer of the electron from the singly occupied copper $3d$ orbital to an empty $5d$ orbital of the rare-earth ion. This mechanism is in the spirit of the superexchange mechanisms proposed by Goodenough [13] in 1963. Similar mechanism is invoked recently by Tchougreff [14] to explain the ferromagnetic ordering in decamethylferrocenium tetracyanoethenide, by Kinoshita and coworkers [15] to explain the intermolecular ferromagnetic coupling in paranitrophenyl nitronyl nitroxide and by Girerd and coworkers [16] to interpret the magnetic properties of μ -oxo Mn(II) compounds. Indeed, the interplay of direct exchange which favours parallel spin alignment on a given site (Hund's rule) and kinetic exchange which favours antiparallel alignment of spins on bonded pair of sites (by virtue of Pauli principle) is what determines the spin state of a system [17].

Kahn and coworkers modeled the superexchange process using an extended Hückel model and obtained a value for the septet-nonet gap from the computed $J_{\text{Gd-Cu}}$ in fairly good agreement with the experimental result. But in this model strong electron-electron

repulsions among the $4f$ electrons arising out of the contraction of the lanthanide $4f$ orbitals as well as between $4f$ electrons and the $5d$ electrons have not been taken into account. These electron repulsions, according to the recent reports seems to play an important role in explaining the photoemission spectra of rare-earth compounds [18]. Another factor which was not accounted for in the above model is the splitting of the free ion Ln(III) ground state due to the ligand field effect, on the grounds that spin-orbit coupling term is an order of magnitude larger than crystal field term [19]. The need of a detailed microscopic model taking these factors into account has also been felt since some exceptions to ferromagnetic coupling in Gd(III) complexes were recently reported for Cu(II) complexes and nitronyl nitroxide radicals [20–22]. To explain this unexpected antiferromagnetic coupling, suggestions were put forward that the observed coupling is actually the sum of two contributions [22], one from the direct overlap of the magnetic orbitals of the ligand with the f orbitals, which leads to antiferromagnetism and the other from the overlap with the s and d orbitals, which results in ferromagnetism. But the rare-earth f orbitals are highly contracted to the core which prompted Kahn *et. al.* [10] to neglect the transfer between the $4f$ orbitals of Ln(III) and $3d$ orbitals of Cu(II).

In this paper we present our studies of a detailed microscopic model in which we have taken into account explicit electron-electron correlations, crystal field effects and spin-orbit effects in the Lanthanide ion. We have spanned the parameter space of the model and have obtained phase diagrams for different fillings of the $4f$ orbitals to elucidate the role played by the different parameters in controlling the ground state spin. The paper is organized as follows. In section 2, we introduce the many-body model Hamiltonian. We also briefly describe the computational schemes employed in solving the model Hamiltonian. In section 3 we discuss the results obtained by our computations on the model. Finally, we summarize our results in section 4.

2 Many Body Model and Computational Scheme

The remarkable feature of these systems is that the parallel or anti-parallel alignment of spins in the ground state is independent of the details of structure. This led Kahn *et. al.* to propose a general mechanism by the virtue of which these systems attain either the high spin or the low spin state. The contraction of the $4f$ orbitals towards the nucleus means that they are very weakly delocalized towards the Cu $3d$ orbitals and thus, there is a vanishingly small overlap between the Gd $4f$ and Cu $3d$ orbitals. This also means that any interaction between the ground state configuration and the charge transfer configuration corresponding to a $3d - 4f$ transfer cannot stabilize the low spin $S = 3$ state. So, Kahn *et. al.* neglected these $3d - 4f$ interactions and considered charge transfer to the higher $5d$ level, ie, the $3d$ electron being transferred to an empty $5d$ orbital of the

Gd(III) ion. This process results in the $S = 4$ high spin state being stabilized due to a strong Hund's coupling between the $4f$ and $5d$ levels. This process is illustrated in Fig. 1. The mechanism outlined by Kahn *et. al.* is however, still in a non-interacting picture since it does not take into account explicit Coulomb repulsions among the electrons in different orbitals. Another factor missing in the model is the effect of the crystal field interaction on the various orbitals involved. In order to take these effects into account and understand how they modify the ground state spin in these systems, we have studied these systems by employing an interacting Hamiltonian. This Hamiltonian is an extension of the Pariser-Parr-Pople model [11] to more than one orbital per atom and includes direct exchange interaction terms between electrons in orbitals on the same site. This model has been greatly successful in describing ground state and low-lying excited states of various strongly correlated systems.

When we have more than one orbital per site, the ZDO (zero differential overlap) approximation is not valid and we should include electron repulsion integrals when two electrons are in different orbitals on the same atom. In this case, we will have several kinds of nonzero two electron integrals on the same site. The interaction part of the Hamiltonian in this modified ZDO approximation is given by,

$$H_{int} = 1/2 \sum_{i,j} \sum_{\mu,\mu',\nu,\nu'} [i_{\mu}i_{\mu'}|j_{\nu}j_{\nu'}](\hat{E}_{i\mu,i\mu'}\hat{E}_{j\nu,j\nu'} - \delta_{ij}\delta_{\mu'\nu}\hat{E}_{i\mu,j\nu'}) \quad (1)$$

Here, i, j are site indices and μ, ν correspond to the orbital on a particular site. The two electron integrals $[i_{\mu}i_{\mu'}|j_{\nu}j_{\nu'}]$ contains contributions from various kinds of situations. The case $\mu = \mu' = \nu = \nu'$ and $i = j$ corresponds to the electron-repulsion integral for the electrons in the same orbital of a single site. The most general case for the repulsion integral on the same site is realized when there are at least four orbitals per site and $\mu \neq \mu' \neq \nu \neq \nu'$ and $i = j$. For the f orbitals on the same site, we have assumed that (i) the intra-orbital repulsion integral $[f_{\mu}f_{\mu}|f_{\mu}f_{\mu}]$ is the same for all the seven f orbitals, independent of μ , and given by U_f , (ii) the inter f orbital repulsion integral $[f_{\mu}f_{\mu}|f_{\nu}f_{\nu}]$ is also independent of the particular pair of f orbitals μ, ν and given by U_{ff} and (iii) that there is also a single direct exchange integral J_{ff} between two f orbitals corresponding to the integral $[f_{\mu}f_{\nu}|f_{\mu}f_{\nu}]$. We have further assumed that all other electron repulsion integrals involving the f orbitals are zero. The direct exchange integral although smaller than the electron repulsion integrals in (i) and (ii) by about an order of magnitude, is crucial to the study of magnetic gaps. Indeed, the small value of the exchange integral is also reflected in the small value of the magnetic gaps. For our rare-earth system, in most cases, we have considered all the seven $4f$ orbitals and one $5d$ orbitals for the rare-earth metal ion and a single $3d$ orbital for the Cu ion. Inclusion of $5d$ orbital introduces corresponding electron repulsion parameters U_{5d} , U_{fd} and J_{fd} . Electron repulsion integrals of the form $[i_{\mu}i_{\mu'}|j_{\nu}j_{\nu}]$ which in our case corresponds to repulsion integral V_{fd} between an electron in the $4f$

orbital of the rare-earth and an electron in the $3d$ orbital of the Cu ion is parametrized using Ohno interpolation scheme [12]. All other inter-site electron repulsion integrals are neglected.

Using the above parameters, the model Hamiltonian is written as ,

$$\begin{aligned}
\hat{H} = & \sum_{f=1}^7 \Delta_f f \hat{n}_f + \Delta_{3d} \hat{n}_{3d} + \Delta_{5d} \hat{n}_{5d} \\
& + t_{3d-4f} \sum_{f=1}^7 (\hat{E}_{f,3d} + H. c.) + t_{3d-5d} (\hat{E}_{3d,5d} + H. c.) \\
& + \frac{U_f}{2} \sum_{f=1}^7 \hat{n}_f (\hat{n}_f - 1) + \frac{U_{3d}}{2} \hat{n}_{3d} (\hat{n}_{3d} - 1) + \frac{U_{5d}}{2} \hat{n}_{5d} (\hat{n}_{5d} - 1) \\
& + U_{ff} \sum_{f>f'=1}^7 \hat{n}_f \hat{n}_{f'} + U_{fd} \sum_{f=1}^7 \hat{n}_f \hat{n}_{5d} \\
& + V_{3d-4f} \sum_{f=1}^7 \hat{n}_f \hat{n}_{3d} + V_{3d-5d} \hat{n}_{3d} \hat{n}_{5d} \\
& + J_{ff} \sum_{f>f'=1}^7 [(\hat{E}_{f,f'} \hat{E}_{f,f'} + \hat{E}_{f',f} \hat{E}_{f',f} + H. c.) - \hat{n}_f - \hat{n}_{f'}] \\
& + J_{fd} \sum_{f=1}^7 [(\hat{E}_{f,5d} \hat{E}_{f,5d} + \hat{E}_{5d,f} \hat{E}_{f,5d} + H. c.) - \hat{n}_f - \hat{n}_{5d}] \tag{2}
\end{aligned}$$

where we define the electron hop operator \hat{E}_{ij} which hops an electron of either spin from orbital j to orbital i as $\hat{E}_{ij} = \sum_{\sigma} \hat{a}_{i,\sigma}^{\dagger} \hat{a}_{j,\sigma}$; $\hat{a}_{i,\sigma}^{\dagger}$ ($\hat{a}_{i,\sigma}$) creates (annihilates) an electron with spin σ in the i^{th} orbital. The first two lines in Eqn. 2 corresponds to the noninteracting part of the Hamiltonian. The first line corresponds to the site energy contribution. Here, we assume that the degeneracy of the $4f$ levels are lifted uniformly by the nonsymmetric crystalline field, leading to equally spaced levels with an energy level spacing Δ_f . The $5d$ orbital of the rare-earth and the $3d$ orbital of the Cu are located at energies Δ_{5d} and Δ_{3d} respectively. The second line in the Hamiltonian corresponds to the transfer part between the individual $4f$ orbitals and the Cu $3d$ orbital, with a transfer parameter t_{3d-4f} and a transfer between the rare-earth $5d$ orbital and the Cu $3d$ orbital with a transfer parameter t_{3d-5d} . Here, we assume that the orbitals on the same site are eigenstates of the site Hamiltonian and therefore do not mix. The third line in the Hamiltonian corresponds to Hubbard electron repulsion terms, U_f , U_{5d} and U_{3d} being the repulsion parameters for the orbitals we have considered. The fourth line involves inter-orbital interactions on the

rare-earth site with U_{ff} for repulsion between electrons in different $4f$ orbitals and U_{fd} for repulsion between electrons in the $4f$ and $5d$ orbitals at the rare-earth site. The fifth line in the Hamiltonian corresponds to the inter-site interaction between electrons on the rare-earth site and those in the Cu $3d$ orbital. The parameters V_{3d-4f} and V_{3d-5d} are obtained using Ohno parametrization with intersite distance of 2\AA . The last two lines in the Hamiltonian correspond to the exchange term involving orbitals on the rare-earth site. The exchange term has both diagonal and off-diagonal contributions, the latter coming from two successive electron hops.

We set up the matrix for the above Hamiltonian in the Valence Bond (VB) basis. The VB basis being eigenstates of the total spin operator is ideal for this study since the total spin of the targeted state is known exactly. The VB method is discussed extensively in [23]. The resulting Hamiltonian matrix is non-symmetric and is diagonalized using Rettrup's algorithm [24]. The low energy states are determined for a range of values of the parameters in the Hamiltonian. These studies were carried out at various $4f$ orbital fillings, corresponding to different lanthanide ions.

We have also carried out calculations that includes spin-orbit interactions between the $4f$ orbital angular momentum and the electron spin. The spin-orbit contribution is given by,

$$\hat{H}_{S.O.} = \sum_i \xi(r_i) l(i) \cdot s(i) \quad (3)$$

where $\xi(r_i)$ is given by

$$\xi(r_i) = \frac{\hbar}{2m^2c^2r} \frac{\partial V}{\partial r}, \quad (4)$$

with usual notations. Integrating over the radial part leads to the spin-orbit interaction parameter λ and $\hat{H}_{S.O.}$ can be written in the second quantized form as,

$$\begin{aligned} \hat{H}_{S.O.} = & \frac{\lambda}{2} \sum_m [\sqrt{l(l+1) - m(m+1)} \hat{a}_{m+1,\beta}^\dagger \hat{a}_{m,\alpha} \\ & + \sqrt{l(l+1) - m(m-1)} \hat{a}_{m-1,\alpha}^\dagger \hat{a}_{m,\beta} \\ & + m(\hat{a}_{m,\alpha}^\dagger \hat{a}_{m,\alpha} - \hat{a}_{m,\beta}^\dagger \hat{a}_{m,\beta})] \end{aligned} \quad (5)$$

where, l takes the value 3, corresponding to the f orbitals and m varies from -3 to 3. In the presence of spin-orbit interactions, the total spin of the system is not conserved. Thus, we cannot employ the VB method for their solution. Instead we employ a constant M_s basis, typified by Slater determinants, to set-up the Hamiltonian matrix. The Hamiltonian matrix in this case is symmetric and Davidson algorithm [25] gives the low-lying states.

Since the states are now labelled by the total angular momentum $J = L + S$, we compute the expectation value of \hat{S}^2 and use the nearest integer or half-integer value S such that the expectation value is nearest to $S(S + 1)$. This allows identifying the approximate total spin of the states.

3 Results and Discussion

The Hamiltonian in Eq. (2) has several parameters and exploring the phase space for all possible filling of the $4f$ rare-earth orbital is very compute intensive. In order to meaningfully explore the parameter space, we have firstly identified the most sensitive parameters and fixed the remaining parameters at generally accepted values. We have considered three different models for the Ln-Cu system. These correspond to (a) ignoring the $5d$ levels on the rare-earth ion, (b) including all the five $5d$ orbitals on the rare-earth ion and (c) including only one $5d$ orbital on the rare-earth ion. In all these cases, we have found that the inter-orbital electron repulsion strength, U_{ff} , in the $4f$ orbital and the crystal-field splitting of the $4f$ orbitals Δ_f are the parameters to which the spin of the ground state is most sensitive. The remaining parameters have been held fixed at the following values. The value of the exchange constants J_{ff} and J_{fd} are fixed at $0.1U_f$ and $0.2U_f$, respectively. The U_{3d} on copper is fixed at 5eV, t_{3d-4f} at 0.1eV and in all the models and in models (b) and (c) t_{3d-5d} was fixed at 0.5eV, U_{5d} at 5eV and U_{fd} at 4 eV. In what follows, we discuss our results, first for the half-filled $4f$ shell and then for rare-earth ions with other $4f$ occupancies, for various values of the parameters U_{ff} and Δ_f . In the last subsection we explore the effect of spin-orbit interactions on the "spin" of the ground state for these parameter values.

3.1 $4f^7$ system

For the half-filled Ln ion in model (a) where the $5d$ orbitals on the rare-earth site is ignored, we always obtain a low-spin ground state, for all values of U_f and Δ_f that we scanned. Changing the parameters t_{3d-4f} and U_{3d} within reasonable limits still leaves the system in a low-spin ground state. This is because, the path way for delocalization of the electrons in the high-spin state is blocked (Fig. 2a), depriving the system of kinetic stabilization. Thus, this model does not produce a high-spin ground state for any value of Δ_f or U_{ff} , although experimentally most $4f^7$ -radical or Cu(II) systems are found in high-spin ground state.

In the presence of the empty $5d$ orbitals (Fig. 2b), the Hunds rule stabilizes the state in which the $4f$ and $5d$ electrons have the same spin alignment. This in turn implies that the

high-spin state in which the unpaired electron in the Cu $3d$ orbital has the same orientation as the spins in the $4f$ orbital has lower energy than the low-spin state. Our studies on model (b) is therefore expected to produce a high-spin ground state. However, what is quite unexpected is that we find a very interesting variation in the ground state spin; the ground state is a high-spin state for small Δ_f ($\leq 0.1eV$) and $1.5eV \leq U_{ff} \leq 4.5eV$ at $U_f = 10eV$. For other values of U_{ff} , the ground state is in a low-spin state. Thus, in the quantum phase diagram (Fig. 3), there is a re-entrant low-spin phase. For small values of U_{ff} the virtual low-spin state in which the intermediate rare-earth ion configuration is $4f^8$ has lower energy than the configuration $4f^7 5d^1$ and this would lead to a low-spin ground state of the complex. At intermediate values of U_{ff} , the virtual state $4f^7 5d^1$ has a lower energy than the $4f^8$ virtual state and we should expect the ground state to be a high spin state. For large values of U_{ff} , the ground state configuration of the rare-earth ion is no longer $4f^7$ but $4f^6 5d^1$. In this case, the hopping between the $5d$ and $3d$ orbitals favours antiparallel alignment of spins. The $3d$ to $4f$ hop prefers parallel alignment on energy consideration but there is only one channel for this process as against seven channels for antiparallel alignment. Therefore, the ground state switches back to low-spin state for high values of U_{ff} .

We also find that the high-spin state is quite stable and continues to be the ground state, even when the t_{3d-4f} transfer integral is increased to upto twice its value of $0.1eV$ we have used in all our studies. Furthermore, we find that the Δ_f parameter which corresponds to the crystal field splitting of the $4f$ levels also plays a significant part; when Δ_f is small, the cross over from high-spin ground state to low-spin ground state occurs at a higher value of U_{ff} . Above a critical value of Δ_f , the system is found only in the low-spin ground state (Fig. 3).

This picture is not changed in any significant way, in going to model (c) where only one $5d$ orbital is used instead of the five $5d$ orbitals on the rare-earth considered in model (b). In view of this, in cases away from half-filled $4f$ levels, we have considered only a single $5d$ orbital (model (c)) for reasons of computational feasibility.

3.2 Systems with non-half-filled $4f$ shell

We have studied systems with various $4f$ orbital fillings using the Hamiltonian in Eq. 2 and the general result we find is that when U_{ff} is small, systems with fewer than seven $4f$ electrons have a high-spin ground state while those with more than seven $4f$ electrons have a low-spin ground state (Fig. 6). This result follows even in model (a) where we have neglected the participation of the $5d$ orbitals in delocalization. This result is quite simply understood from the direct exchange interaction term. In the less than half-filled case (Fig. 4), an electron hop from the Cu $3d$ orbital to the rare-earth $4f$ orbital can result

in a high-spin or a low-spin rare-earth ion. If the unpaired electrons on the rare-earth have the same alignment as the hopping electron, we obtain the high-spin state, otherwise the resulting state of the rare-earth ion is a low-spin state. The direct exchange interaction within $4f$ orbitals ensures that the virtual high-spin state of the ion has lower energy than the virtual low-spin state. Therefore, the ferromagnetic alignment of the spins on the Ln-Cu system is favoured over antiferromagnetic alignment in the less than half-filled case. In the more than half-filled case, electron-hop is only permitted when the hopping electron has an alignment opposite that of the rare-earth spin (Fig. 5), because of Pauli exclusion principle. Thus, in the more than half-filled case, we find a low-spin ground state. Indeed, this picture is not significantly changed when we go to model (c) where we have the copper $3d$ electron having an additional delocalization path-way via the rare-earth $5d$ orbital. What we also note is that the t_{3d-5d} does not have a significant effect, unlike in the half-filled case.

We should expect the crystal field splitting parameter, Δ_f , of the $4f$ orbitals to play an important role in determining the ground state of the Ln-Cu system, since in case of large Δ_f , the rare earth ion ground state would not be a high-spin ground state. For a similar reason, we should also expect the inter-orbital repulsion, U_{ff} also to determine the magnetic state of the dimer. In Fig. 6, we present some typical quantum phase diagrams for the less than half-filled case as well as the more than half-filled case. In the less than half-filled case, for large U_{ff} , we obtain a low-spin ground state since the energy difference between high-spin and low-spin rare-earth virtual states obtained on electron transfer, is not significant. However, number of low-spin virtual states are more than the high-spin virtual states and we obtain a low-spin ground state for large U_{ff} . The experimental system with less than half-filled $4f$ shell are all known to be in the low-spin (afm) ground state. Thus, these systems all correspond to the large U_{ff} limit. The critical value of U_{ff} decreases with increase in $4f$ occupancy from 1 to 7.

In Fig. 6, we show some typical quantum phase diagrams for the more than half-filled case. Here again we note a quantum phase transition, from the low-spin ground state to the high-spin ground state, at large values of U_{ff} . This transition is also qualitatively understood in the following way. At large U_{ff} , the low-energy exchange path way is via the empty $5d$ orbitals. The direct exchange interaction between the $5d$ and $4f$ orbitals on the rare-earth ensure that the virtual states corresponding to high-spin state has a lower energy than those corresponding to low-spin state. Thus, the dimer switches to a high-spin ground state with increasing U_{ff} . All the experimental systems known in this case correspond to high-spin ground state. Thus, we are in the large U_{ff} limit in real systems even in systems in which the rare earth ion has more than half-filled $4f$ shell. This is also intuitively expected since the $4f$ shell is chemically well isolated in rare earth compounds. The critical value of U_{ff} in this case increases with increase in the $4f$ occupancy beyond

half-filling.

3.3 Role of spin-orbit interactions

It is well known that the spin-orbit interactions in rare-earth ions are important and often larger than the crystal field effects. We have considered only the spin-orbit interaction for electrons in the $4f$ orbital of the rare-earth and have ignored it in the $5d$ orbital of the rare-earth as well as in the $3d$ orbital of the copper. Our method of calculation takes into account spin-orbit interactions at the same level as all other interactions in the Hamiltonian in Eq. 2, since we set-up and solve the total Hamiltonian (Hamiltonians in Eq. 2 plus Eq. 3). Thus, our studies go beyond the usual L-S or j-j coupling schemes in which either the spin-orbit interaction is treated perturbatively within a given manifold of L and S or the electron repulsions are treated perturbatively over the individual spin-orbit coupled states. We have used values of λ in the range 0 to 1.0 eV, for the U_{ff} and Δ_f values [7]. In the presence of spin-orbit interactions, total spin is no longer a good quantum number. Thus we cannot strictly classify the eigenstates by their spins as low-spin or high-spin states. However, we find that the expectation value of \hat{S}^2 in the low-lying states, in all the cases we have studied, is close enough to allowed S values in the absence of spin-orbit interactions. This aids in labelling the states as ‘high-spin’ or ‘low-spin’ states.

What we find is that the role of the spin-orbit interaction in determining the ground state magnetic moment of the Ln-Cu dimers is negligible. In Fig. 7, we show the quantum phase diagram for three representative cases corresponding to less than half-filled, half-filled and more than half-filled $4f$ shells, for $\Delta_f = 0$ as a function of λ and U_{ff} , with all other parameters being the same as in Figs. 3 and 6. We find that the U_{ff} value for which the magnetic transition occurs in the ground state is hardly affected by spin-orbit interactions. Only the re-entrant low-spin state occurs for a slightly smaller value of U_{ff} when spin-orbit interaction is sufficiently strong. We have also explored the role of spin-orbit interaction when the crystal field splitting of the $4f$ orbitals Δ_f is varied near a critical value of U_{ff} . Here we assume that the splitting of the f -orbital manifold by the crystal field does not quench the orbital angular momentum of the resulting state. In Fig. 8 is shown the phase diagram as a function of λ for the case of the rare-earth electron configuration $4f^{10}$. We have chosen this configuration as we found the strongest effect due to spin-orbit interactions in this case. We find in Fig. 7 that high-spin state crosses over to the low-spin state at a smaller value of Δ_f when the spin-orbit interaction strength λ is increased. Thus, we conclude from our studies that the magnetic state of the cluster is only weakly influenced by spin-orbit interactions.

4 Summary and Conclusions

In this paper, we have studied the magnetic ground state of the hetero bi-metallic organic molecular magnets, with a rare-earth ion and a copper ion. These systems have been modeled with a many-body Hamiltonian, which includes the direct exchange interaction and the repulsion between the orbitals on the same site. The systems are shown to exhibit a low-spin or a high-spin ground state depending on the filling of electrons in the rare earth $4f$ levels, and also critically on the inter-orbital $4f$ electron repulsions, U_{ff} and the splitting of the degenerate $4f$ orbitals in the presence of the crystal field, Δ_f . We have obtained a "phase diagram" in the space of these parameters. We find at half-filling a re-entrant low-spin state as U_{ff} is increased for small values of Δ_f . There is also a transition in the ground state spin at other fillings, again for small values of Δ_f , when U_{ff} is increased. The general phase diagram of the ground state spin is shown in Fig. 9, as a function of filling and U_{ff} . This picture is not altered by the inclusion of spin-orbit interactions. To obtain the experimentally observed interaction, corresponding to low-spin ground state below half-filling and high-spin ground state above half-filling of the $4f$ shell, we need to assume fairly large inter-orbital repulsions in the $4f$ shell. Our studies show that in systems where the parameters are near the phase boundaries, small changes brought about by the modification of the ligands can lead to drastic changes in the magnetic properties of the complex.

Acknowledgement

We thank Prof. Diptiman Sen for many useful discussions . We thank the Council of Scientific and Industrial Research, India and Board of Research in Nuclear Sciences (BRNS) for financial support.

References

- [1] W. Heitler and F. London, *Z. Phys.*, **44**, 455 (1927).
- [2] For a review, see G. Blondin and J. J. Girerd, *Chem. Rev.* **90**, 1359 (1990).
- [3] O. Kahn, *Molecular Magnetism*, VCH, New York (1993).
- [4] A. Bencini, C. Benelli, A. Caneschi, A. Dei and D. Gatteschi, *Inorg. Chem.* **25**, 572 (1986).
- [5] C. Benelli, A. Caneschi, D. Gatteschi, O. Guillou and L. Pardi, *Inorg. Chem.* **29**, 1750 (1990).

- [6] O. Guillou, O. Kahn, R. L. Oushoorn, Boubekeur and P. Batail, *Inorg. Chim. Acta*, **198-200**, 119 (1992).
- [7] M. L. Kahn, J. P. Sutter, S. Golhen, P. Guionneau, L. Ouahab, O. Kahn and D. Chasseau, *J. Am. Chem. Soc.* **122**, 3413 (2000).
- [8] J. P. Sutter, M. L. Kahn and O. Kahn, *Adv. Mater.* **11**, 863 (1999).
- [9] J. P. Sutter, M. L. Kahn, *Polyhedron* **20**, 1593 (2001).
- [10] M. Andruh, I. Ramade, E. Codjvoi, O. Guillou, O. Kahn and J. C. Trombe, *J. Am. Chem. Soc.* **115**, 1822 (1993).
- [11] R. Pariser and R. G. Parr, *J. Chem. Phys.* **21**, 446 (1953);
- [12] K. Ohno, *Theor. Chem. Acta* **2**, 219 (1964); G. Klopman, *J. Am. Chem. Soc.* **86**, 4550 (1964). J. A. Pople, *Trans. Faraday. Soc.* **49**, 1375 (1953).
- [13] J. B. Goodenough, *Magnetism and Chemical Bond*, p. 165, Interscience, New York (1963).
- [14] A. L. Tchougreeff, *J. Chem. Phys.* **96**, 6026 (1992).
- [15] P. Turek, K. Nozawa, D. Shiomi, K. Awaga, T. Inabe, Y. Maruyama and M. Kinoshita, *Chem. Phys. Lett.* **180**, 327 (1991).
- [16] R. Hotzelmann, K. Wiegardt, U. Flörke, H. J. Haupt, D. C. Weatherburn, J. Bonvoisin, G. Blondin and J. J. Girerd, *J. Am. Chem. Soc.* **114**, 1681 (1992).
- [17] Bhabadyuti Sinha and S. Ramasesha, *Phys. Rev.*, **B48**, 16410 (1993).
- [18] H. Ogasawara, A. Kotani, P. L. Fevre, D. Chandesris and H. Magnan, *Phys. Rev. B* **62**, 7970 (2000).
- [19] C. J. Ballhausen, *Introduction to Ligand Field Theory*, McGraw-Hill, New York (1962).
- [20] C. Lescop, D. Luneau, E. Belorizky, P. Fries, M. Guillot and P. Rey, *Inorg. Chem.* **38**, 5472 (1999).
- [21] J. P. Costes, F. Dahan and A. Dupuis, *Inorg. Chem.* **39**, 169 (2000).
- [22] A. Caneschi, A. Dei, D. Gatteschi, L. Sorace and K. Vostrikova, *Angew. Chem. Int. Ed.* **39**, 246 (2000).

- [23] For a review, see Z. G. Soos and S. Ramasesha in *Valence Bond Theory and Chemical Structure*, Ed. D. J. Klein and N. Trinajstić, Elsevier (1990).
- [24] S. Rettrup, *J. comput. Phys.* **45**, 100 (1982).
- [25] E.R. Davidson, *J. Comput. Phys.*, **17**, 87 (1975).

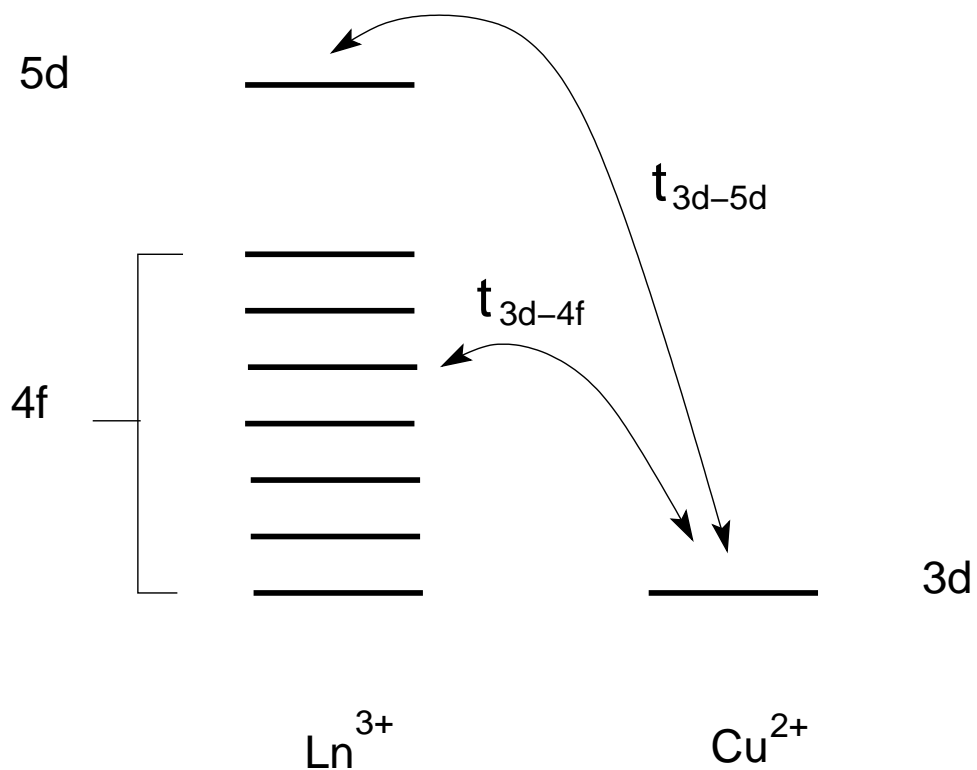
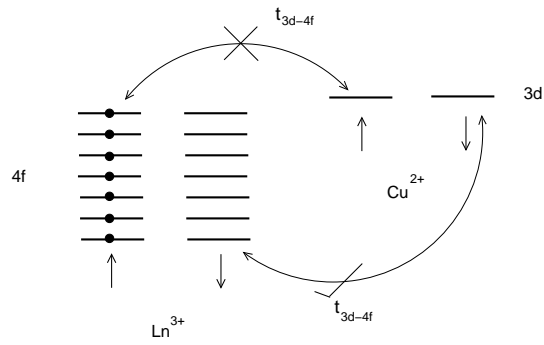
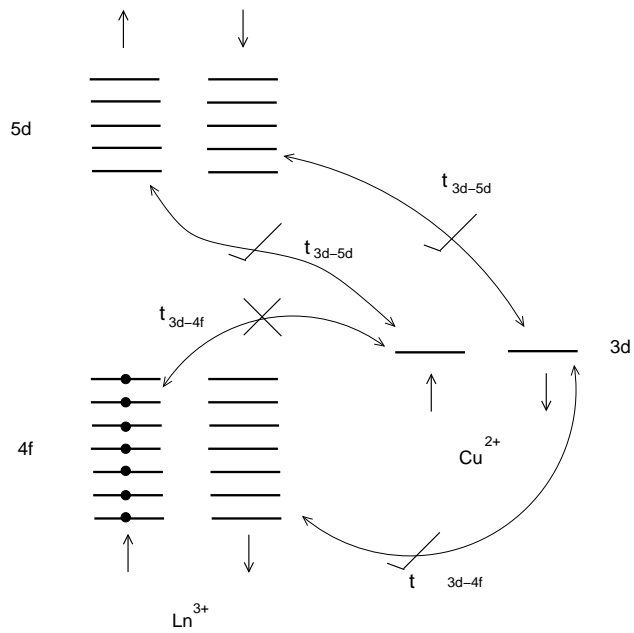


Figure 1: A schematic for the mechanism which leads to the stabilization of the high spin state in the system. Ln^{3+} is the rare-earth (Lanthanide) ion. Transfers between the Cu^{2+} $3d$ and the Ln $4f$ and $5d$ levels are indicated by the arrows. The $3d \rightarrow 4f$ transfer leads to a higher energy virtual state than $3d \rightarrow 5d$ transfer.



(a)



(b)

Figure 2: A schematic diagram showing the various possible virtual states obtained by electron transfer from $3d$ orbital of Cu(II) to orbitals of Ln(III) , when $4f$ is half-filled - (a) without $5d$ levels, (b) with all the $5d$ levels taken into consideration.

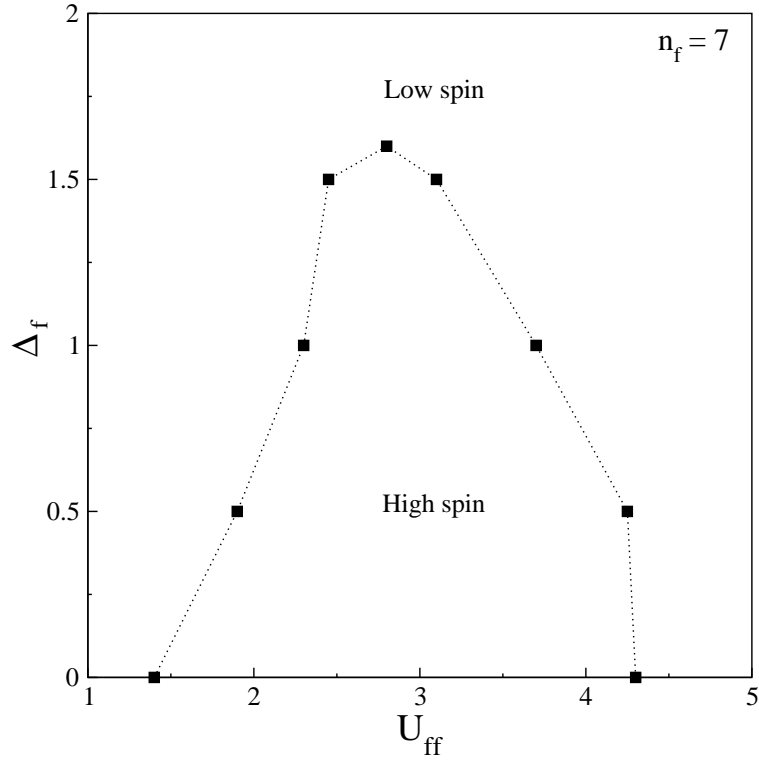


Figure 3: Phase diagram in the space of parameters U_{ff} and Δ_f , for the case of rare-earth $4f$ levels being half-filled. Other parameter values are as follows: $U_f=10$ ev, $t_{3d-4f}=0.1$ ev, $t_{3d-5d}=0.5$ ev, $U_{5d}=5$ ev, $U_{3d}=5$ ev, $U_{fd}=4$ ev, $J_{ff}=0.1$ ev and $J_{fd}=0.2$ ev

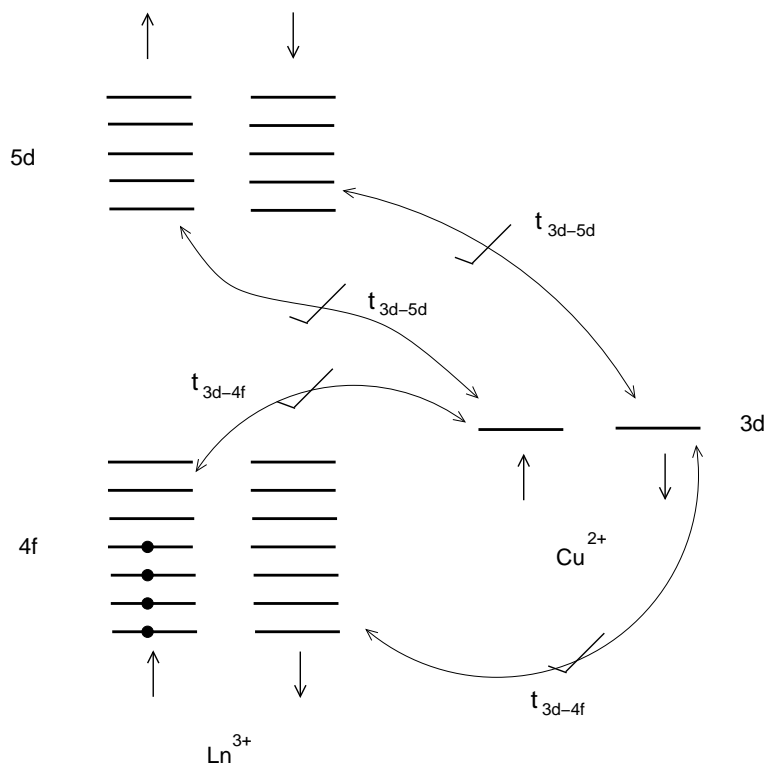


Figure 4: A schematic diagram showing the various possible virtual states obtained by electron transfer from $3d$ orbital of $\text{Cu}(\text{II})$ to orbitals of $\text{Ln}(\text{III})$, when $4f$ is less than half-filled.

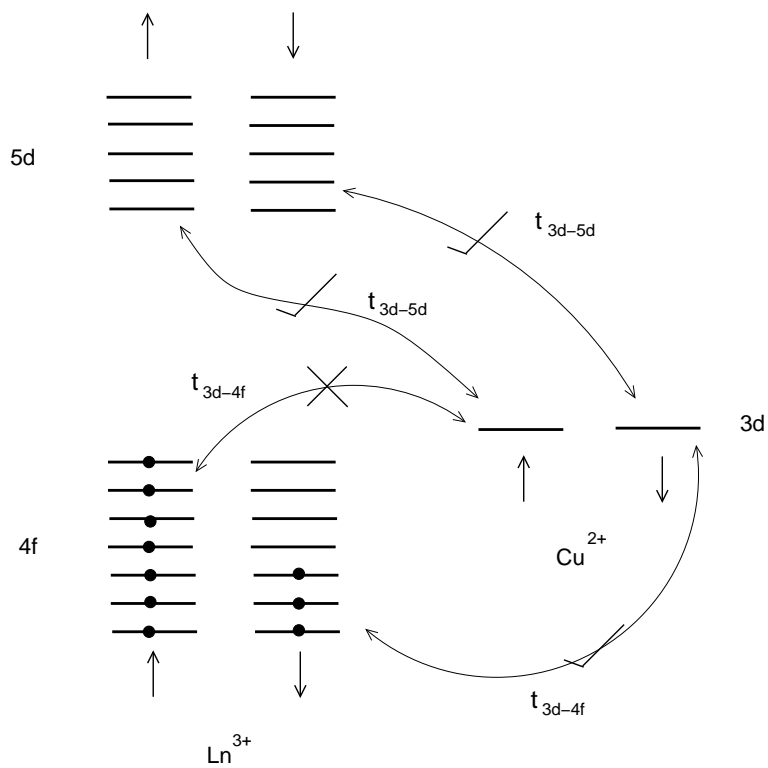


Figure 5: A schematic diagram showing the various possible virtual states obtained by electron transfer from 3d orbital of Cu(II) to orbitals of Ln(III), when 4f is more than half-filled.

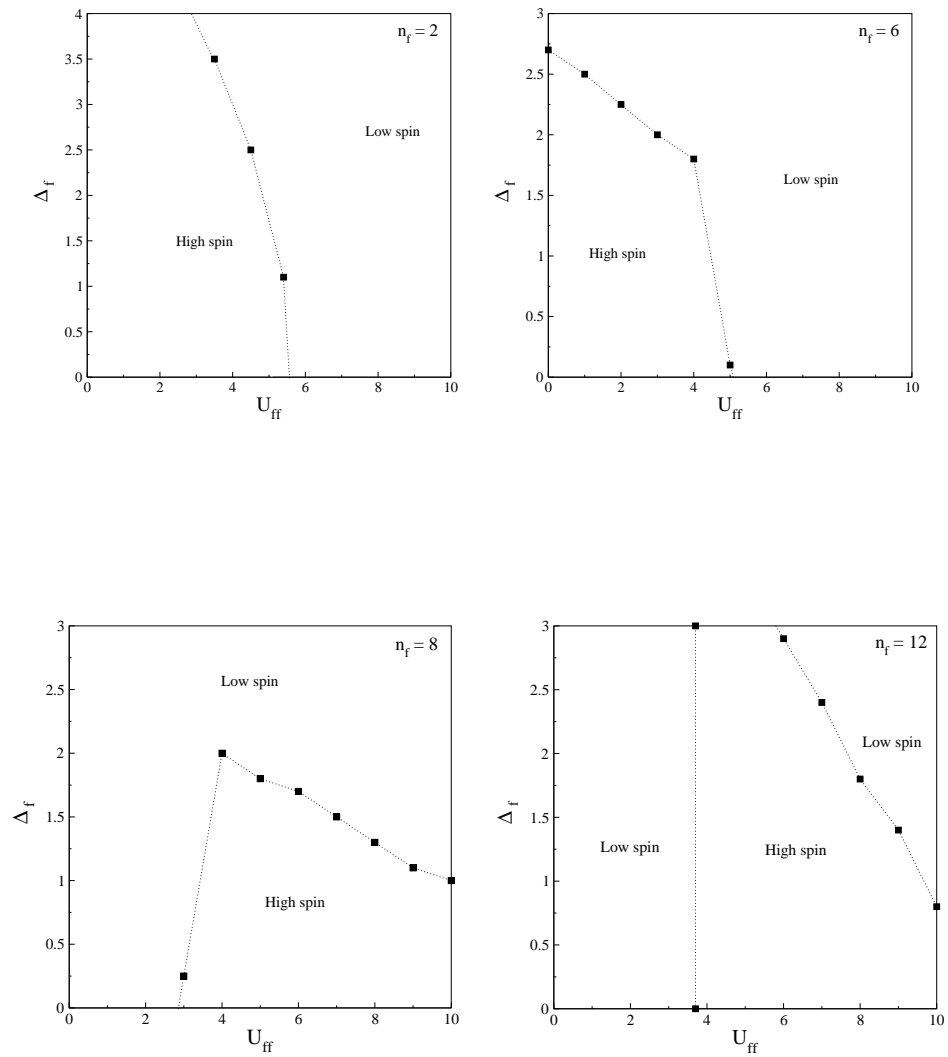


Figure 6: Phase diagram in the space of parameters U_{ff} and Δ_f , for the case of rare-earth $4f$ levels being away from half-filling. Other parameter values are same as in Fig. 3

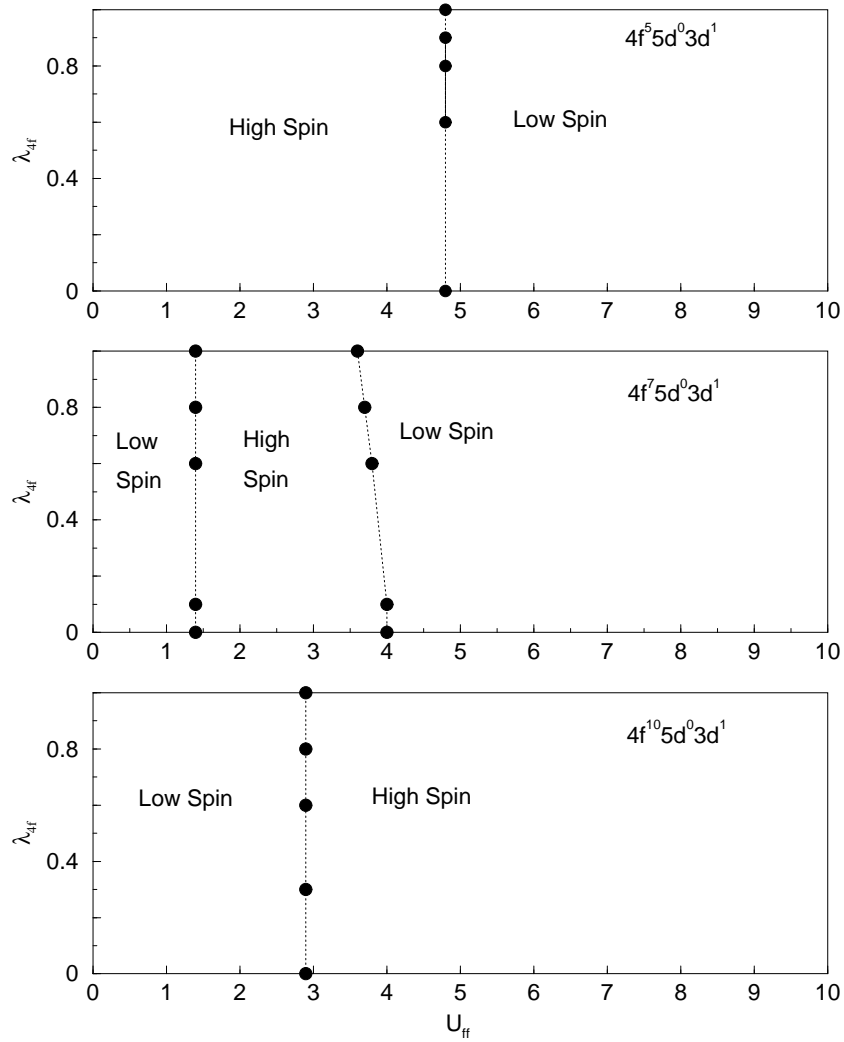


Figure 7: Phase diagram in the space of parameters U_{ff} and λ_{4f} , for the case of rare-earth $4f$ levels at less than half-filling, half-filling and more than half-filling. Other parameter values are same as in Fig. 3

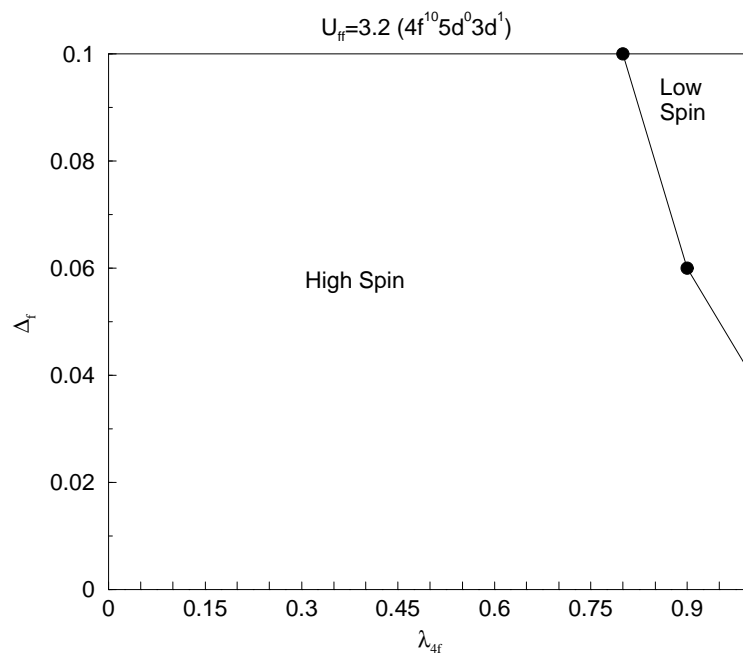


Figure 8: Phase diagram in the space of parameters λ_{4f} and Δ_f for the case of rare-earth $4f$ levels being away from half-filling. Other parameter values are same as in Fig. 3

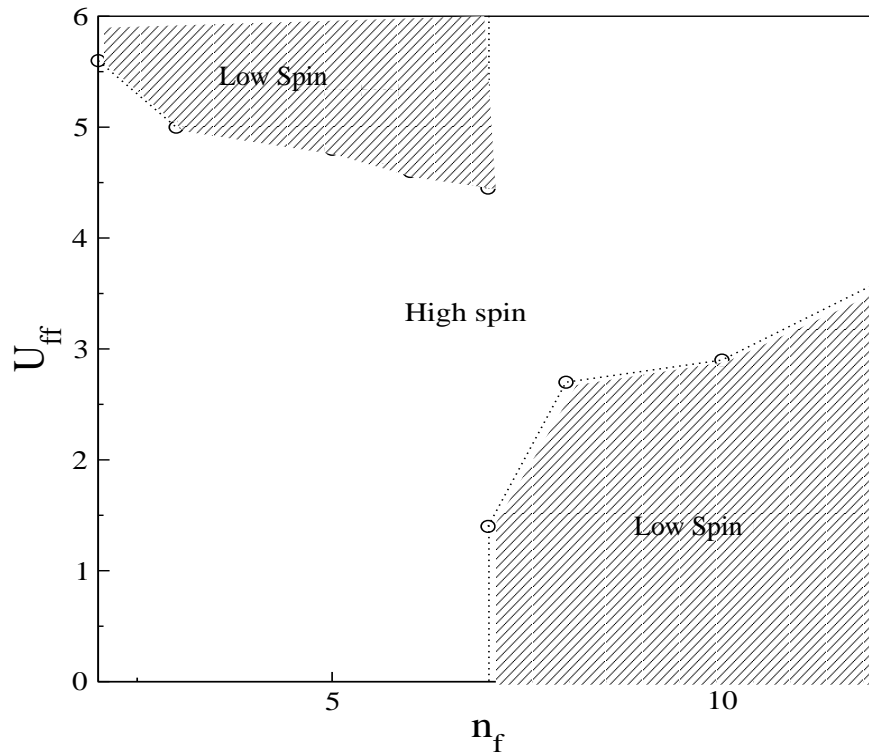


Figure 9: A phase diagram for the Ln(III)-Cu(II) system, between parameters U_{ff} and n_f , n_f being the 4f level filling. Δ_f value was fixed at 0.1 eV. Other parameter values are same as in Fig. 3