ON DEVIATIONS OF THE SAMPLE MEAN

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1. Introduction. Let X_1 , X_2 , \cdots be a sequence of independent and identically distributed random variables. Let a be a constant, $-\infty < a < \infty$, and for each $n = 1, 2, \cdots$ let

$$(1) p_n = P\left(\frac{X_1 + \cdots + X_n}{n} \ge a\right).$$

It is assumed throughout the paper that the distribution of X_1 and the given constant a satisfy the conditions stated in the following paragraph. These conditions imply that $p_n > 0$ for each n, and that $p_n \to 0$ as $n \to \infty$. The object of the paper is to obtain an estimate of p_n , say q_n , which is precise in the sense that

$$(2) q_n/p_n = 1 + o(1) as n \to \infty.$$

Let t be a real variable, and let $\varphi(t)$ denote the moment generating function (m.g.f.) of X_1 , i.e., $\varphi(t) = E(e^{tX_1})$, $0 < \varphi \le \infty$. Define

$$\psi(t) = e^{-at}\varphi(t).$$

Let T denote the set of all values t for which $\varphi(t) < \infty$. We suppose that $P(X_1 = a) \neq 1$, that T is a non-degenerate interval, and that there exists a positive τ in the interior of T such that $\psi(\tau) = \inf_t \{\psi(t)\} = \rho$ (say). These conditions are satisfied if, for example, $\varphi(t) < \infty$ for all t, $E(X_1) = 0$, a > 0, and $P(X_1 > a) > 0$. In any case, τ and ρ are uniquely determined by

(4)
$$\frac{\varphi'(\tau)}{\varphi(\tau)} = a \quad \text{and} \quad \rho = \psi(\tau),$$

where $\varphi' = d\varphi/dt$, and we have $0 < \rho < 1$.

There are three separate cases to be considered.

Case 1: The distribution function (d.f.) of X_1 is absolutely continuous, or, more generally, this d.f. satisfies Cramér's condition (C) [1, p. 81].

Case 2: X_1 is a lattice variable, i.e., there exist constants x_0 and d > 0 such that X_1 is confined to the set $\{x_0 + rd : r = 0, \pm 1, \pm 2, \cdots\}$ with probability one

Case 3: Neither Case 1 nor Case 2 obtains.

We can now state

THEOREM 1. There exists a sequence b_1 , b_2 , \cdots of positive numbers b_n such that

(5)
$$p_n = \frac{\rho^n}{(2\pi n)^{\frac{1}{2}}} b_n [1 + o(1)], \quad \log b_n = O(1)$$

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as $n \to \infty$. In Cases 1 and 3, b_n is independent of n. This last also holds in Case 2 if $P(X_1 = a) > 0$.

The proof of Theorem 1, and of Theorem 2 below, is given in Sections 2-5. The present determination of b_n is given by (4), (9) and (33) in Cases 1 and 3, and by (4), (8), (37), (38) and (46) in Case 2. The following refinements of Theorem 1 are available in Cases 1 and 2:

THEOREM 2. (Cases 1 and 2). For each $j=1, 2, \cdots$ there exists a bounded (possibly constant) sequence $c_{j,1}, c_{j,2}, \cdots$ such that, for any given positive integer k,

(6)
$$p_n = \frac{\rho^n}{(2\pi n)^{\frac{1}{2}}} b_n \left[1 + \frac{c_{1,n}}{n} + \frac{c_{2,n}}{n^2} + \cdots + \frac{c_{k,n}}{n^k} \right] \left[1 + O\left(\frac{1}{n^{k+1}}\right) \right]$$

as $n \to \infty$.

The sequences $\{c_{j,n}\}$ are given explicitly for Cases 1 and 2 in Sections 3 and 4 respectively. It would be interesting to know whether (6) holds in Case 3 as well, perhaps with the $\{c_{j,n}\}$ determined according to the formula for Case 1.

Estimates in the form (5) or (6) were first obtained by Cramér [2, pp. 20-21] in the case when X_1 has an absolutely continuous component (so that Case 1 obtains). Cramér showed that in the latter case (6) holds for every k (with b_n and each $c_{j,n}$ independent of n), and determined b_n . Our method of proof in the general case (cf. Sections 2-5) is essentially a variant or extension of Cramér's method. Case 2 was treated recently by Blackwell and Hodges [3] by a different method. It is shown in [3] that (6) holds for k = 1 in Case 2, under the restriction on n and n that n0 and n1 that n2 and n3 are determined explicitly. Some other references bearing on the problem under consideration are [4], [5] and [6].

In the following Section 2 it is shown that p_n can be expressed as $\rho^n I_n$, where I_n is a certain integral; $0 < I_n < 1$, and $I_n = O(n^{-\frac{1}{2}})$ as $n \to \infty$. I_n can be estimated by application of certain refinements [1], [7] of the central limit theorem. This estimation of I_n is carried out in Sections 3, 4 and 5 for Cases 1, 2 and 3 respectively. It may be added here that, as was pointed out in [2], direct application of the central limit theorem (or refinements thereof) to p_n defined by (1) does not, in general, yield approximations q_n which satisfy (2).

In Section 6 we describe certain numerical approximations to p_n which are suggested by Theorems 1 and 2 and their proofs.

2. Lemmas. Let $Y_1 = X_1 - a$, and let F be the (left-continuous) distribution function (d.f.) of Y_1 , $F(y) = P(Y_1 < y)$. Let G be defined by $G(z) = \int_{-\infty < y < z} \rho^{-1} e^{\tau y} dF(y)$. Since $E(e^{\tau Y_1}) = \psi(\tau) = \rho$, it is clear that G is a probability d.f.. Let Z_1 be a random variable distributed according to G.

Lemma 1. The m.g.f. of Z_1 exists in a neighborhood of the origin. We have

(7)
$$E(Z_1) = 0, \quad 0 < \operatorname{Var}(Z_1) < \infty.$$

Proof. Let $\xi(t)$ denote the m.g.f. of Z_1 . Then $\xi(t) = \psi(\tau + t)/\rho$ for all t,

by (1)

by (3) and the definition of Z_1 . Since $\psi(t) < \infty$ in a neighborhood of $t = \tau$, it follows that $\xi(t) < \infty$ in a neighborhood of t = 0. Consequently, $E \mid Z_1 \mid^{\tau} < \infty$ for $r = 1, 2, 3, \dots$ and $E(Z_1^{\tau}) = \{d^{\tau}\xi/dt\}_{t=0}$. In particular, $E(Z_1) = \{d\xi/dt\}_{t=0} = \psi'(\tau)/\rho = 0$, since $\psi(t)$ is minimum at $t = \tau$, and τ is in the interior of T. It remains to show that $\text{Var}(Z_1) > 0$. Suppose to the contrary that $\text{Var}(Z_1) = 0$; then $P(Z_1 = 0) = 1$; hence $P(Y_1 = 0) = 1$, i.e., $P(X_1 = a) = 1$, which is contrary to our assumptions. This completes the proof.

Let $Var(Z_1)$ be denoted by σ^2 . It follows from the preceding paragraph and (4) that

(8)
$$\sigma^2 = \frac{\varphi''(\tau)}{\varphi(\tau)} - a^2.$$

Define

(9)
$$\alpha = \sigma \tau, \quad (0 < \alpha < \infty).$$

Let Z_1 , Z_2 , \cdots be a sequence of independent and identically distributed random variables. For each n, let

$$(10) U_n = \frac{Z_1 + \dots + Z_n}{n^{\frac{1}{2}\sigma}}$$

and

(11)
$$H_n(x) = P(U_n < x), \quad (-\infty < x < \infty).$$

LEMMA 2. $p_n = \rho^n I_n$, where

(12)
$$I_n = n^{\frac{1}{2}} \alpha \int_0^\infty e^{-n^{\frac{1}{2}} \alpha x} \left[H_n(x) - H_n(0) \right] dx.$$

PROOF. Let $Y_j = X_j - a$ for $j = 1, 2, \dots, n$. Then

 $p_n = P(Y_1 + \cdots + Y_n \ge 0)$

(13)
$$= \int \cdots \int_{y_1 + \cdots + y_n \ge 0} dF(y_1) \cdots dF(y_n)$$

$$= \rho^n \int \cdots \int_{z_1 + \cdots + z_n \ge 0} e^{-\tau(z_1 + \cdots + z_n)} dG(z_1) \cdots dG(z_n)$$

$$= \rho^n \int_{0 \le x < \infty} e^{-n\frac{1}{2}\alpha x} dH_n(x)$$
 by (9), (10), (11)
$$= \rho^n I_n^* \text{ say.}$$

It follows by integration by parts that I_n^* defined in (13) is equal to I_n , and this completes the proof.

A theorem of Chernoff [4] states that $p_n \leq \rho^n$ for every n, and that for any

given positive $\rho_0 < \rho$, we have $p_n \ge \rho_0^n$ for all sufficiently large n. A simple proof of Chernoff's theorem can be given as follows. Since $0 \le H_n(x) - H_n(0) \le 1$ for every n and $x \ge 0$, we have $I_n \le 1$ and hence $p_n \le \rho^n$ for every n, by Lemma 2. To establish the second part of the theorem, we note first that $\lim_{n\to\infty} H_n(x) = \Phi(x)$ for every x, where

(14)
$$\Phi(x) = \int_{-\infty}^{x} (2\pi)^{-\frac{1}{2}} e^{-\frac{1}{2}t^2} dt \qquad (-\infty < x < \infty),$$

by (7), (10), (11) and the central limit theorem. Let ϵ be a positive constant. Then

$$I_n \ge n^{\frac{1}{2}} \alpha \int_{\epsilon}^{\infty} e^{-n\frac{1}{2}\alpha x} \left[H_n(x) - H_n(0) \right] dx$$

$$\ge \left[H_n(\epsilon) - H_n(0) \right] n^{\frac{1}{2}} \alpha \int_{\epsilon}^{\infty} e^{-n\frac{1}{2}\alpha x} dx$$

$$= \left[H_n(\epsilon) - H_n(0) \right] e^{-n\frac{1}{2}\alpha \epsilon}.$$

Hence $\lim \inf_{n\to\infty} \{n^{-\frac{1}{2}} \log I_n\} \ge -\alpha\epsilon$. Since $I_n \le 1$ for every n, and since ϵ is arbitrary, it follows that $n^{-\frac{1}{2}} \log I_n = o(1)$. Hence $n^{-1} \log p_n = \log \rho + o(1)$, by Lemma 2, and this is equivalent to the conclusion desired.

The preceding argument depends only on the central limit theorem. In the following sections we estimate I_n more accurately by substituting the expansions of $H_n(x)$ due to Cramér [1] and Esseen [7] in the right side of (12). The remainder of this section is concerned with preparations for this application of the Cramér-Esseen expansions. Almost all the considerations of the following paragraphs are well known, and we include them here only for the sake of completeness.

Let $\eta(w)$ denote the m.g.f. of Z_1/σ . According to Lemma 1, $\eta < \infty$ in a neighborhood of w = 0. For $j = 2, 3, \cdots$ let λ_j be defined by

(15)
$$\lambda_2 = \frac{1}{2}; \qquad \lambda_j = (j!\sigma^j)^{-1} (d^j/dt^j) \{ \log \varphi(t) \}_{t=\tau} \quad (j=3,4,\cdots).$$

It should be noted that $j!\lambda_j$ is the jth cumulant of the distribution of Z_1/σ . The m.g.f. of U_n , with U_n defined by (10), is

$$[\eta(w/n^{\frac{1}{2}})]^n = \exp [n \sum_{j=0}^{\infty} \lambda_j (w/n^{\frac{1}{2}})^j].$$

Clearly, $[\eta(w/n^{\frac{1}{2}})]^n$ exp $(-w^2/2)$ is analytic in a domain independent of n, and can be expanded there as a power series in w. By regrouping the terms of this series according to powers of n we shall have

$$[\eta(w/n^{\frac{1}{2}})]^n e^{-w^2/2} = \sum_{j=0}^{\infty} n^{-\frac{1}{2}j} P_j(w)$$

where the P_j are polynomials. P_j is of degree 3j, and P_j is even or odd according as j is even or odd. The first few polynomials are

$$P_{0}(w) = w^{0} \equiv 1,$$

$$P_{1}(w) = \lambda_{3}w^{3},$$

$$P_{2}(w) = \lambda_{4}w^{4} + \frac{1}{2}\lambda_{3}^{2}w^{6},$$

$$P_{3}(w) = \lambda_{5}w^{5} + \lambda_{3}\lambda_{4}w^{7} + \frac{1}{6}\lambda_{3}^{3}w^{9},$$

$$P_{4}(w) = \lambda_{6}w^{6} + (\frac{1}{2}\lambda_{4}^{2} + \lambda_{3}\lambda_{5})w^{8} + \lambda_{3}^{2}\lambda_{4}w^{10} + \frac{1}{24}\lambda_{3}^{4}w^{12}.$$

Write $\Phi^{(0)}(x) = \Phi(x)$ and $\Phi^{(r)}(x) = (d^r/dx^r)\Phi(x)$ for $r = 1, 2, \dots$, where Φ is given by (14). Let $P_j(-\Phi)$ denote the function of x obtained by replacing w^r with $(-1)^r\Phi^{(r)}(x)$ in the polynomial $P_j(w)$. It is clear that each $P_j(-\Phi)$ is absolutely continuous and of bounded variation in $(-\infty, \infty)$. It should also be noted that $P'_j(-\Phi)$ is square integrable with respect to Lebesgue measure.

In the following, for any function K(x) of bounded variation in $(-\infty, \infty)$, we denote the c.f. of K by $\chi(t \mid K)$, i.e.,

(18)
$$\chi(t \mid K) = \int_{-\infty}^{\infty} e^{itx} dK(x)$$

for every real t. If K is absolutely continuous, χ is, of course, $(2\pi)^{\frac{1}{2}}$ times the Fourier transform of K'. The reader may refer to [8, Chapters I–III] for such elements of Fourier transform theory as are used in this paper.

LEMMA 3. For every j, t, and x

(19)
$$\chi(t \mid P_{j}(-\Phi)) = P_{j}(it) e^{-\frac{1}{2}t^{2}}$$

and

(20)
$$P'_{j}(-\Phi) = (2\pi)^{-\frac{1}{2}} \int_{-\infty}^{\infty} e^{-itx} P_{j}(it) \ d\Phi(t).$$

PROOF. As is pointed out in [1, p. 49], we have

(21)
$$\chi(t \mid \Phi^{(r)}) = (-it)^r e^{-\frac{1}{2}t^2}$$

for $r=0,1,\cdots$. Suppose, for given j, that $P_j(w)=\sum_{r=0}^N a_r\,w^r$, where the a_r and N are constants (depending on j). Then $P_j(-\Phi)=\sum_{r=0}^N a_r\,(-1)^r\,\Phi^{(r)}(x)$; hence the left side of (19) equals $\sum_0^N a_r\,(-1)^r\,\chi(t\,|\,\Phi^{(r)})$; (19) now follows from (21). The relation (20) follows from (19) by the inversion formula for the Fourier transform, since $dP_j(-\Phi)=P_j'(-\Phi)\,dx$, and $d\Phi(t)=(2\pi)^{-\frac{1}{2}}\,e^{-\frac{1}{2}t^2}\,dt$.

A probability d.f. K(x) is said to satisfy condition (C) if

$$\lim \sup_{|t|\to\infty} |\chi(t|K)| < 1.$$

In the following lemma the F_i are arbitrary probability d.fs.

LEMMA 4. If F_1 satisfies (C), and if F_1 is absolutely continuous with respect to F_2 , then F_2 also satisfies (C).

Proof. In this proof, for any probability d.f. K let K^* denote the symmetrized d.f. defined by $K^*(x) = \int_{-\infty}^{\infty} K(x+y) dK(y)$. We then have

$$\chi(t \mid K^*) = \int_{-\infty}^{\infty} \cos(tx) dK^* = |\chi(t \mid K)|^2$$

for all t.

Suppose, contrary to the lemma, that there exists a sequence $\{t_j: j=1,2,\cdots\}$ such that $|t_j| \to \infty$ and $|\chi(t_j|F_2)| \to 1$ as $j \to \infty$. It then follows from the above paragraph with $K = F_2$ that $\int_{-\infty}^{\infty} \cos(t_j x) dF_2^* \to 1$. Hence $\cos(t_j x) \to 1$ in F_2^* -measure. Since F_2 -measure dominates F_1 -measure, it is easily seen that F_2^* -measure dominates F_1^* -measure. Consequently, $\cos(t_j x) \to 1$ in F_1^* -measure. It now follows from the above paragraph with $K = F_1$ that $|\chi(t_j|F_1)|^2 \to 1$ as $j \to \infty$, which is impossible. This completes the proof.

We conclude this section with a description of the functions $S_1(x)$, $S_2(x)$ which occur in the Euler-Maclaurin sum formulae, and which are required in the analysis of Case 2. It is convenient to define S_1 as follows:

(22)
$$S_1(x) = \frac{1}{2} - x$$
 for $0 < x \le 1$; $S_1(x+1) \equiv S_1(x)$.

For $j \geq 2$, S_j may be defined as

(23)
$$S_{j}(x) = \begin{cases} \frac{1}{2^{j-1}} \sum_{r=1}^{\infty} \frac{\cos(2\pi rx)}{(\pi r)^{j}} & (j \text{ even}) \\ \frac{1}{2^{j-1}} \sum_{r=1}^{\infty} \frac{\sin(2\pi rx)}{(\pi r)^{j}} & (j \text{ odd}). \end{cases}$$

Each S_j is a bounded and periodic function; S_j is absolutely continuous for $j \ge 2$; and at each non-integral x we have

(24)
$$S'_{1}(x) = -1, S'_{j+1}(x) = (-1)^{j} S_{j}(x) \qquad (j = 1, 2, \cdots).$$

3. I_n in Case 1. Suppose that the d.f. of X_1 satisfies (C). Since $Y_1 = X_1 - a$, it is plain that F, the d.f. of Y_1 , also satisfies (C). It is easily seen that F and G (the d.f. of Z_1) are absolutely continuous with respect to each other. It therefore follows from Lemma 4 with $F_1 = F$ and $F_2 = G$ that G also satisfies (C).

Let k be an arbitrary but fixed positive integer. It follows from the conclusion of the preceding paragraph by Cramér's theorem [1, p. 81] that $H_n(x) = K_n(x) + R_n(x)$, where

(25)
$$K_n(x) = \sum_{i=0}^k n^{-\frac{1}{2}i} P_i(-\Phi)$$

and $R_n(x)$ is of the order $n^{-(k+1)/2}$ uniformly in x. It follows hence from (12) that

(26)
$$I_n = n^{\frac{1}{2}} \alpha \int_0^\infty e^{-n^{\frac{1}{2}} \alpha x} \left[K_n(x) - K_n(0) \right] dx + O(n^{-\frac{1}{2}k - \frac{1}{2}}).$$

We have

(27)
$$\chi(t \mid K_n) = \sum_{j=0}^k n^{-\frac{1}{2}j} P_j(it) e^{-\frac{1}{2}t^2}$$

by (19) and (25). Let $f_n(x) = \exp(-n^{\frac{1}{2}}\alpha x)$ for $x \ge 0$ and $f_n(x) = 0$ other-

wise. Then $\int_{-\infty}^{\infty} e^{itx} f_n(x) dx = 1/(n^{\frac{1}{2}}\alpha - it) = g_n(t)$ say. Consequently, by first using integration by parts and then Parseval's formula, it follows that

(28)
$$n^{\frac{1}{2}}\alpha \int_{\mathbf{0}}^{\infty} e^{-n^{\frac{1}{2}}\alpha x} \left[K_{n}(x) - K_{n}(0) \right] dx = \int_{\mathbf{0}}^{\infty} e^{-n^{\frac{1}{2}}\alpha x} K'_{n}(x) dx \\ = \int_{-\infty}^{\infty} f_{n}(x) K'_{n}(x) dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} \overline{g_{n}(t)} \chi(t \mid K_{n}) dt.$$

It follows from (26), (27) and (28) that

(29)
$$\alpha (2\pi n)^{\frac{1}{2}} I_n = \int_{-\infty}^{\infty} \left(1 + \frac{it}{n^{\frac{1}{2}}\alpha}\right)^{-1} \left(\sum_{j=0}^{k} n^{-\frac{1}{2}j} P_j(it)\right) d\Phi(t) + O(n^{-\frac{1}{2}k}).$$

Define

(30)
$$\mu_{r,s} = \int_{-\infty}^{\infty} (it)^r P_s(it) d\Phi(t) \qquad (r, s = 0, 1, 2, \cdots).$$

Since P_s is an even [odd] polynomial if s is even [odd], and since $\int_{-\infty}^{\infty} t^{2j+1} d\Phi(t) = 0$ for $j = 0, 1, 2, \cdots$, it follows that each $\mu_{r,s}$ is a real constant, and that

(31)
$$\mu_{r,s} = 0 \quad \text{if } r + s \text{ is odd.}$$

Now, $(1 + itn^{-\frac{1}{2}}\alpha^{-1})^{-1} = \sum_{0 \le r < k} (-itn^{-\frac{1}{2}}\alpha^{-1})^r + n^{-\frac{1}{2}k} t^k \omega_n(t)$, where $|\omega|$ is bounded in n and t. Since Φ has finite moments of all orders, it therefore follows from (29), (30) and (31) that

(32)
$$\alpha (2\pi n)^{\frac{1}{2}} I_n = \sum_{0 \le j < \frac{1}{2}k} n^{-j} \left\{ \sum_{r+s=2j} \left(-\frac{1}{\alpha} \right)^r \mu_{r,s} \right\} + O(n^{-\frac{1}{2}k}).$$

Since $p_n = \rho^n I_n$, and since $\mu_{0,0} = 1$, it follows by replacing k with 2k + 2 in (32) that (6) holds for any given k, with

$$(33) b_n = \alpha^{-1}$$

and

(34)
$$c_{j,n} = \sum_{r+s=2j} \left(-\frac{1}{\alpha}\right)^r \mu_{r,s} \qquad (j=1,2,\cdots)$$

for every n. This establishes Theorem 2, and hence also Theorem 1, in Case 1. It follows from (17) and (30) that the coefficients $\mu_{r,s}$ required to compute $c_{1,n}$ according to (34) are

(35)
$$\mu_{2,0} = -1$$

$$\mu_{1,1} = 3\lambda_3$$

$$\mu_{0,2} = 3\lambda_4 - \frac{15}{2}\lambda_3^2,$$

where the λ_i are given by (15). Similarly, $c_{2,n}$ can be computed from

$$\mu_{4,0} = 3$$

$$\mu_{3,1} = -15\lambda_3$$

$$(36) \qquad \mu_{2,2} = -15\lambda_4 + 105\lambda_3^2$$

$$\mu_{1,3} = -15\lambda_5 + 105\lambda_3\lambda_4 - \frac{315}{2}\lambda_3^3$$

$$\mu_{0,4} = -15\lambda_6 + 105(\frac{1}{2}\lambda_4^2 + \lambda_3\lambda_5) - \frac{945}{2}\lambda_3^2\lambda_4 + \frac{10395}{24}\lambda_3^4.$$

We conclude this section with a remark concerning the role of Cramér's theorem [1, p. 81] in the preceding argument. Suppose that H_n is absolutely continuous, and that H'_n is square integrable over $(-\infty, \infty)$. It then follows, by integrating (12) by parts and using Parseval's formula, that

(29*)
$$\alpha (2\pi n)^{\frac{1}{2}} I_n = \int_{-\infty}^{\infty} \left(1 + \frac{it}{n^{\frac{1}{2}}\alpha} \right)^{-1} \left\{ \left[\eta \left(\frac{it}{n^{\frac{1}{2}}} \right) \right]^n e^{\frac{1}{2}t^2} \right\} d\Phi(t)$$

where η is, as before, the m.g.f. of Z_1/σ . (The square integrability condition is imposed here for the validity of Parseval's formula, and can be replaced by others, e.g., that $(1+t^2)^{-\frac{1}{2}} \mid \eta(it) \mid$ be integrable). According to (16), the function in curly brackets on the right side of (29*) can be expressed as $\sum_{j=0}^{\infty} n^{-\frac{1}{2}j} P_j(it)$. By comparing (29) and (29*) it is seen that, from a technical point of view, the role of Cramér's theorem in the present special case is to guarantee that when $\sum_{j=0}^{\infty} n^{-\frac{1}{2}j} P_j$ is replaced by $\sum_{j=0}^{k-1} n^{-\frac{1}{2}j} P_j$ on the right side of (29*), the error introduced is indeed of the order $n^{-\frac{1}{2}k}$. The same remark, but with (29*) replaced by a rather different formula for I_n , applies to the role of Esseen's theorem in the argument of the following section.

4. I_n in Case 2. Suppose that X_1 is a lattice variable. Let d be the maximum span of X_1 , i.e., d>0 is the g.c.d. of the differences between consecutive possible values of X_1 . Let x_0 be the number such that $a \le x_0 < a + d$, and such that the possible values of X_1 are included in the set $\{x_0 + rd : r = 0, \pm 1, \pm 2, \cdots\}$. Let

(37)
$$\beta = d/\sigma, \qquad \gamma = \tau d, \qquad \kappa = (x_0 - a)/d$$

It should be noted that $0 \le \kappa < 1$. For each n, let

$$\theta_n = n\kappa - [n\kappa], \qquad 0 \le \theta_n < 1,$$

where [x] denotes the greatest integer contained in x.

Let k be an arbitrary but fixed positive integer. It follows from Esseen's theorem for the lattice case [7, p. 61] that $H_n(x) = K_n(x) + L_n(x) + R_n(x)$, where $K_n(x)$ is given by (25), R_n is of the order $n^{-(k+1)/2}$ uniformly in x, and L_n

is defined as follows. For any $j=1, 2, \cdots$ let $h_j=1$ if $j\equiv 1$ or 2 (mod 4) and $h_j=-1$ if $j\equiv 0$ or 3 (mod 4). Then

(39)
$$L_n(x) = \sum_{j=1}^k n^{-\frac{1}{2}j} h_j \beta^j S_j(n^{\frac{1}{2}}\beta^{-1}x - \theta_n) K_n^{(j)}(x) = \sum_{j=1}^k M_{j,n}(x) \text{ say,}$$

where $K_n^{(j)}$ is the jth derivative of K_n . It follows hence from (12) that

(40)
$$I_{n} = n^{\frac{1}{2}} \alpha \int_{0}^{\infty} e^{-n^{\frac{1}{2}} \alpha x} \left[K_{n}(x) - K_{n}(0) \right] dx + \sum_{j=1}^{k} n^{\frac{1}{2}} \alpha \int_{0}^{\infty} e^{-n^{\frac{1}{2}} \alpha x} \left[M_{j,n}(x) - M_{j,n}(0) \right] dx + O(n^{-\frac{1}{2}k - \frac{1}{2}}).$$

The first term on the right side of (40) is (cf., (28)) equal to $\int_0^\infty e^{-n^{\frac{1}{2}}\alpha x} K_n^{(1)}(x) dx$. We observe next that, for $j \geq 2$,

$$n^{\frac{1}{2}}\alpha \int_{0}^{\infty} e^{-n^{\frac{1}{2}}\alpha x} \left[M_{j,n}(x) - M_{j,n}(0) \right] dx = \int_{0}^{\infty} e^{-n^{\frac{1}{2}}\alpha x} M'_{j,n}(x) dx$$

$$= n^{-\frac{1}{2}j} h_{j} \beta^{j} \int_{0}^{\infty} e^{-n^{\frac{1}{2}}\alpha x} \left[S_{j}(y_{n}) K_{n}^{(j+1)}(x) + (-1)^{j-1} n^{\frac{1}{2}} \beta^{-1} S_{j-1}(y_{n}) K_{n}^{(j)}(x) \right] dx$$

$$= n^{-\frac{1}{2}j} h_{j} \beta^{j} \int_{0}^{\infty} e^{-n^{\frac{1}{2}}\alpha x} S_{j}(y_{n}) K_{n}^{(j+1)}(x) dx$$

$$- n^{-\frac{1}{2}(j-1)} h_{j-1} \beta^{j-1} \int_{0}^{\infty} e^{-n^{\frac{1}{2}}\alpha x} S_{j-1}(y_{n}) K_{n}^{(j)}(x) dx$$

$$= N_{j,n} - N_{j-1,n} \text{ (say)}.$$

In (41), we have put $n^{i}\beta^{-1}x - \theta_{n} = y_{n}$, and used integration by parts, (24), and the identity $(-1)^{i}h_{i} = h_{i-1}$. In order to evaluate the contribution of $M_{1,n}$ to the right side of (40), suppose for the moment that $0 < \theta_{n} < 1$, and let

(42)
$$\zeta_0 = 0, \qquad \zeta_r = (r - 1 + \theta_n)\beta/n^{\frac{1}{2}} \qquad (r = 1, 2, \cdots).$$

Let A_r denote the open interval (ζ_r, ζ_{r+1}) . Then $S_1(y_n)$ is linear in x over each A_r (cf. (22)), and its derivative there equals $-n^{\frac{1}{2}}\beta^{-1}$. By writing $\int_0^\infty = \sum_{r=0}^\infty \int_{A_r}^\infty$, and applying integration by parts to \int_{A_r} , it follows without difficulty that

(43)
$$n^{\frac{1}{2}\alpha} \int_{0}^{\infty} e^{-n^{\frac{1}{2}\alpha x}} M_{1,n}(x) dx = -\int_{0}^{\infty} e^{-n^{\frac{1}{2}\alpha x}} K_{n}^{(1)}(x) dx + N_{1,n} + M_{1,n}(0) + \beta n^{-\frac{1}{2}} \sum_{r=1}^{\infty} e^{-\gamma(r-1+\theta_{n})} K_{n}^{(1)}(\zeta_{r}),$$

where $\gamma = \alpha\beta = \tau d$ (cf. (37)). Now, $S_1(x)$ is a left-continuous function of x. It follows hence that, for given n, the left and right sides of (43) are right-

continuous in θ_n . Since (43) holds for each θ_n in (0, 1), we conclude that (43) is valid for $\theta_n = 0$ also.

Since S_k and $K_n^{(k+1)}$ are bounded functions, it is plain from the definition of $N_{j,n}$ (cf. (41)) that $N_{k,n}$ is of the order $n^{-\frac{1}{2}k-\frac{1}{2}}$. It therefore follows from (40), (41) and (43) that

(44)
$$I_n = \beta n^{-\frac{1}{2}} \sum_{r=1}^{\infty} e^{-\gamma (r-1+\theta_n)} K_n^{(1)}(\zeta_r) + O(n^{-\frac{1}{2}k-\frac{1}{2}}).$$

Now, according to (20) and (25),

(45)
$$K_n^{(1)}(\zeta_r) = (2\pi)^{-\frac{1}{2}} \int_{-\infty}^{\infty} e^{-it\zeta_r} \left(\sum_{j=0}^k n^{-\frac{1}{2}j} P_j(it) \right) d\Phi(t)$$

for every r. Let us write

(46)
$$z = e^{-\gamma}, \quad b_n = [\beta/(1-z)]z^{\theta_n}.$$

It follows from (44) and (45) that

$$(47) b_n^{-1} (2\pi n)^{\frac{1}{2}} I_n = \int_{-\infty}^{\infty} \frac{(1-z) \exp\left[-it\beta\theta_n/n^{\frac{1}{2}}\right]}{(1-z \exp\left[-it\beta/n^{\frac{1}{2}}\right])} \cdot \left(\sum_{0}^{k} n^{-\frac{1}{2}i} P_j(it)\right) d\Phi(t) + O(n^{-\frac{1}{2}k}).$$

For any θ and any $j = 0, 1, 2, \cdots$ let

(48)
$$\ell_{j}(\theta) = \frac{1}{j!} \left\{ \frac{d^{j}}{dw^{j}} \left(\frac{(1-z)e^{-\theta w}}{(1-ze^{-w})} \right) \right\}_{w=0}.$$

It then follows easily from (31) and (47) that

(49)
$$b_n^{-1}(2\pi n)^{\frac{1}{2}}I_n = \sum_{0 \le j < k/2} n^{-j} \{ \sum_{r+s=2j} \beta^r \ell_r(\theta_n) \mu_{r,s} \} + O(n^{-\frac{1}{2}k}).$$

By replacing k with 2k + 2 in (49) we see that (6) holds for any given k, with b_n given by (46), and

(50)
$$c_{j,n} = \sum_{r+s=2j} \beta^r \, l_r(\theta_n) \, \mu_{r,s} \, .$$

This establishes Theorem 2 in Case 2, and hence also the first part of Theorem 1. To complete the proof of Theorem 1 in Case 2, we see from (37), (38) that $P(X_1 + \cdots + X_n = na) > 0$ implies $\theta_n = 0$. Consequently if $P(X_1 = a) > 0$ then $\theta_n = 0$ for every n, and hence $b_n = \beta/(1-z)$ for every n.

It may be worthwhile to note that in the present case b_n can be expressed as $\alpha^{-1}[\gamma e^{\gamma(1-\theta_n)}/(e^{\gamma}-1)]$, which shows that, in general, b_n oscillates about the value α^{-1} (cf. (33)) as $n \to \infty$ through the sequence 1, 2, \cdots .

An alternative formula for the coefficients l_i required in (50) is

(51)
$$\ell_{j}(\theta) = (-1)^{j} \sum_{r+s=j} \frac{\theta^{r}}{r!s!} \left\{ (1-z) \left(z \frac{d}{dz} \right)^{s} (1-z)^{-1} \right\}.$$

From (51) it is easily seen that, with u = z/(1-z),

$$l_{0} \equiv 1$$

$$l_{1} = -(\theta + u)$$

$$(52) \quad l_{2} = \frac{1}{2} \{ (\theta + u)^{2} + u(1 + u) \}$$

$$l_{3} = -\frac{1}{6} \{ (\theta + u)^{3} + 3u(1 + u)\theta + u(1 + u)(1 + 5u) \}$$

$$l_{4} = \frac{1}{24} \{ (\theta + u)^{4} + 6u(1 + u)\theta^{2} + 4u(1 + u)(1 + 5u)\theta + 23u^{4} + 36u^{3} + 14u^{2} + u \}.$$

The coefficients $c_{1,n}$ and $c_{2,n}$ can be computed from (35), (36), (50) and (52). The formulae for b_n and $c_{1,n}$ with $\theta = 0$ agree with the results of [3].

5. I_n in Case **3.** If X_1 is not a lattice variable, then neither is Z_1 . It follows hence from a theorem of Esseen [7, p. 49] that $H_n(x) = \Phi(x) + n^{-\frac{1}{2}}f(x) + n^{-\frac{1}{2}}r_n(x)$, where $f(x) = (\text{const.}) (1 - x^2) \exp(-\frac{1}{2}x^2)$, and $r_n(x) \to 0$ uniformly in x as $n \to \infty$. The contribution of $n^{-\frac{1}{2}}f$ to I_n is $n^{-\frac{1}{2}}\int_0^\infty e^{-(n)^{\frac{1}{2}}ax}f'(x) dx$, which is easily seen to be of the order $n^{-\frac{1}{2}}$. It follows that

(53)
$$I_{n} = n^{\frac{1}{2}} \alpha \int_{0}^{\infty} e^{-n^{\frac{1}{2}} \alpha x} \left[\Phi(x) - \Phi(0) \right] dx + o(n^{-\frac{1}{2}})$$

$$= \int_{0}^{\infty} e^{-n^{\frac{1}{2}} \alpha x} \Phi'(x) dx + o(n^{-\frac{1}{2}})$$

$$= e^{\frac{1}{2} n \alpha^{2}} \left[1 - \Phi(n^{\frac{1}{2}} \alpha) \right] + o(n^{-\frac{1}{2}})$$

$$= (2\pi n)^{-\frac{1}{2}} \alpha^{-1} + o(n^{-\frac{1}{2}}).$$

In (53), we have used integration by parts, a linear change of variable, and the leading term of the asymptotic formula [9, p. 179]

(54)
$$1 - \Phi(x) = (2\pi)^{-\frac{1}{2}} e^{-\frac{1}{2}x^2} \{x^{-1} - x^{-3} + 3x^{-5} + O(x^{-7})\} \text{ as } x \to \infty.$$

It follows from (53) that (5) holds, with $b_n = \alpha^{-1}$ for every n. This completes the proof of Theorem 1.

Since $\Phi(x) + n^{-1}f(x) = K_n(x)$, where K_n is defined by (25) with k = 1, the conclusion of the preceding paragraph is also available from the argument of Section 3. We have used a direct calculation instead because this calculation suggests the form of the numerical approximations described in the following section.

6. Concluding remarks. Suppose, in a given case, and for given n and a, that it is required to compute the numerical value of p_n defined by (1). In this section we consider approximations of the form

(55)
$$q_n = \rho^n e^{\frac{1}{2}v_n^2} ((1 - \Phi(v_n)),$$

where ρ and Φ are defined by (4) and (14), and v_n is a suitably chosen number.

We shall describe four choices of v_n , called v_n^* , $v_n^{(0)}$, $v_n^{(1)}$, and $v_n^{(2)}$. The resulting values of q_n are denoted by q_n^* , $q_n^{(0)}$, etc.

First consider

$$(56) v_n^* = n^{\frac{1}{2}} \alpha$$

where α is given by (4), (8) and (9). This choice of v_n amounts (cf. (53)) to approximating I_n by replacing H_n with Φ on the right side of (12). It therefore follows from the Esseen-Berry theorem that we always have

(57)
$$|p_{n} - q_{n}^{*}| \leq 2C \frac{\rho^{n}}{n^{\frac{1}{2}}} \frac{E |Z_{1}|^{3}}{\sigma^{3}}$$

where C is a universal constant. Wallace [10, p. 637] states that $C \leq 2.05$. Next, consider

$$(58) v_n^{(0)} = n^{\frac{1}{2}}/b_n$$

where b_n is defined by (33) in Cases 1 and 3, and by (46) in Case 2. (Of course, $q_n^{(0)} = q_n^*$ in Cases 1 and 3). Then $q_n^{(0)}$ satisfies (2), and the o(1) term in (2) is known to be of the order n^{-1} in Cases 1 and 2. Finally, let $c_{j,n}$ be defined according to Section 4 in Cases 1 and 3, and according to Section 5 in Case 2. Define

$$v_n^{(1)} = v_n^{(0)} \left[1 - (b_n^2 + c_{1,n})/n \right]$$

if the expression within the square brackets is positive and $v_n^{(1)} = 0$ otherwise; and

(60)
$$v_n^{(2)} = v_n^{(1)} \left[1 + (b_n^4 + c_{1,n}^2 - b_n^2 c_{1,n} - c_{2,n}) / n^2 \right]$$

if the expression in square brackets is positive and $v_n^{(2)}=0$ otherwise. Then $q_n^{(j)}$ also satisfies (2), and $o(1)=O(n^{-j-1})$ in Cases 1 and 2 (j=1, 2). The stated theoretical properties of the approximations $q_n^{(j)}$ are easy consequences of (5), (6), (54), and (58).

Although (unlike q_n^*) the approximations $q_n^{(j)}$ are derived from asymptotic expansions corresponding to the case when $n \to \infty$ and a is held fixed, the usefulness of these approximations may be wider than is suggested by the derivation. Some evidence to this effect is provided by the fact that if X_1 is normally distributed then $p_n = q_n^{(0)} = q_n^{(1)} = q_n^{(2)}$ for every admissible a and every n.

REFERENCES

- H. Cramér, Random Variables and Probability Distributions, Cambridge University Press, 1937.
- [2] H. Cramér, "Sur un nouveau théorème-limite de la théorie des probabilités," Actualités Scientifiques et Industrielles, No. 736, Hermann Cie, Paris, 1938.
- [3] DAVID BLACKWELL AND J. L. HODGES, "The probability in the extreme tail of a convolution," Ann. Math. Stat., Vol. 30 (1959), pp. 1113-1120.
- [4] HERMAN CHERNOFF, "A measure of asymptotic efficiency for tests of a hypothesis based on the sum of observations," Ann. Math. Stat., Vol. 23 (1952), pp. 493-507.
- [5] V. V. Petrov, "Generalization of Cramér's limit theorem," Uspekhi Mat. Nauk. Vol. 9, No. 4, (1955), pp. 195-202. (In Russian).

- [6] R. R. Bahadur, "Some approximations to the binomial distribution function," Ann. Math. Stat., Vol. 31 (1960), pp. 43-54.
- [7] Carl Gustav Esseen, "Fourier analysis of distribution functions," Acta Mathematica, Vol. 77 (1945), pp. 1-125.
- [8] E. C. Titchmarsh, Introduction to the theory of Fourier Integrals, Oxford University Press, 1937.
- [9] WILLIAM FELLER, An Introduction to Probability and its Applications, Vol. I, 2nd Ed., John Wiley and Sons, New York, 1957.
- [10] David L. Wallace, "Asymptotic approximations to distributions," Ann. Math. Stat., Vol. 29 (1958), pp. 635-654.