

## Asymptotic analysis of some nonlinear problems using Hopf-Cole transform and spectral theory

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**Abstract.** We consider initial boundary value problems for certain nonlinear scalar parabolic equations. A formula for the unique classical solution by Hopf-Cole transformations is obtained and the asymptotic behaviour of the solution as time goes to  $\infty$  is studied.

**Keywords.** Nonlinear parabolic equation; initial boundary value problem; eigenvalue problem; asymptotic behaviour.

### 1. Introduction

Consider the following initial boundary value problem

$$u_t + \frac{1}{2}u_x^2 = \frac{1}{2}u_{xx} + q(x) + \mu, \quad (1.1)$$

$$u(x, 0) = u_0(x), \quad (1.2)$$

$$u_x(0, t) = a, \quad (1.3)$$

$$u_x(1, t) = b, \quad (1.4)$$

strip  $D = \{(x, t): 0 \leq x \leq 1, t \geq 0\}$ .  $\mu$  is a real parameter and  $a$  and  $b$  are real constants. We assume  $q(x)$  is continuous in  $0 \leq x \leq 1$ ,  $u_0(x)$  is twice continuously differentiable in  $0 \leq x \leq 1$ , and  $(u_0)_x(0) = a$ ,  $(u_0)_x(1) = b$ .

Denote by  $D^0 = \{(x, t): 0 \leq x \leq 1, t > 0\}$ . By a classical solution of (1.1)–(1.4) we mean a function  $u(x, t)$  with  $u(x, t)$  and  $u_x(x, t)$  continuous in  $D$  and  $u_t$  and  $u_{xx}$  continuous in  $D^0$  which satisfies the partial differential equation (1.1) in  $D^0$  and the initial and boundary conditions (1.2)–(1.4) in the usual sense.

Using the Hopf-Cole transformation, (see Hopf [2]) we linearize the problem (1.1)–(1.4) and obtain an expression for the solution in terms of the eigenvalues and eigenfunctions of the eigenvalue problem

$$\frac{1}{2}\phi_{xx} = (q(x) + \lambda)\phi, \quad (1.5)$$

$$\phi_x(0) + a\phi(0) = 0, \quad (1.6)$$

$$\phi_x(1) + b\phi(1) = 0. \quad (1.7)$$

We prove the validity of the expression for  $u(x, t)$  and study its asymptotic behaviour

as  $t \rightarrow \infty$  using the following facts concerning the eigenvalues and eigenfunctions of (1.5)–(1.7), (see Birkhoff and Rota [1]).

(a) The spectrum of (1.5)–(1.7) is discrete and can be ordered

$$\lambda_0 > \lambda_1 > \dots, \quad (1.8)$$

$$\lambda_n = -\frac{1}{2}(n^2\pi^2) + O(1) \quad \text{as } n \rightarrow \infty. \quad (1.9)$$

(b) Let  $\phi_n(x)$  be the normalized eigenfunctions corresponding to  $\lambda_n$ , the set  $\{\phi_n(x), n = 0, 1, 2, \dots\}$  is a complete set for  $L^2[0, 1]$ .

Also  $\phi_n(x)$  has the following estimates uniformly in  $x \in [0, 1]$

$$\phi_n(x) = \sqrt{2} \cos \frac{n\pi}{\sqrt{2}}x + \frac{O(1)}{n} \quad \text{as } n \rightarrow \infty, \quad (1.10)$$

$$\phi'_n(x) = -n\pi \sin \frac{n\pi}{\sqrt{2}}x + O(1) \quad \text{as } n \rightarrow \infty. \quad (1.11)$$

(c)  $\phi_0(x) \neq 0 \quad \forall x \in (0, 1)$ .

Since  $\phi_0(x) \neq 0 \quad \forall x \in (0, 1)$  we can assume  $\phi_0(x) > 0$  for  $x \in (0, 1)$ . We claim that

$$\phi_0(x) > 0 \quad \forall x \in [0, 1]. \quad (1.12)$$

If  $\phi_0(0) = 0$ , then by the boundary condition (1.6)  $(\phi_0)_x(0) = 0$ , then  $\phi_0(x) \equiv 0$ , because  $\phi_0(x)$  solves

$$\frac{1}{2}\phi_{xx} = (q(x) + \lambda)\phi,$$

$$\phi(0) = 0,$$

$$\phi_x(0) = 0,$$

and this has only one solution  $\phi(x) \equiv 0$ . By the same argument  $\phi_0(1)$  is also not equal to 0. The claim is proved. In our discussion we always take  $\phi_0(x)$  normalized so that

$$\int_0^1 \phi_0^2(x) dx = 1$$

and

$$\phi_0(x) > 0 \quad \forall x \in [0, 1].$$

This paper is organized as follows. In §2 we prove the uniqueness of classical solution of (1.1)–(1.4) and obtain a valid expression for it. The asymptotic behaviour of the solution as  $t \rightarrow \infty$  is analysed in §3. In §4 are given some comments on the stationary problem and the Dirichlet problem for the Burger's equation is studied in §5.

## 2. An expression for classical solution of (1.1)–(1.4)

First we prove the uniqueness of classical solution of (1.1)–(1.4) by standard energy estimates. Let  $u_1(x, t)$  and  $u_2(x, t)$  be two solutions and let  $Z(x, t) = u_1(x, t) - u_2(x, t)$ .

Then from (1.1)–(1.4) we get,  $Z(x, t)$  which is solved as

$$Z_t + \frac{1}{2}[(u_1)_x + (u_2)_x]Z_x = \frac{1}{2}Z_{xx}, \quad (2.1)$$

$$Z_x(0, t) = 0, \quad (2.2)$$

$$Z_x(1, t) = 0, \quad (2.3)$$

$$Z(x, 0) = 0. \quad (2.4)$$

Multiplying (2.1) by  $Z$  and integrating by parts w.r.t  $x$ , we get using (2.2) and (2.3)

$$\frac{1}{2} \frac{d}{dt} \int_0^1 Z^2(x, t) dx + \frac{1}{2} \int_0^1 (u_1 + u_2)_x Z_x Z dx = -\frac{1}{2} \int_0^1 Z_x^2(x, t) dx. \quad (2.5)$$

Let

$$C_T = \sup_{\substack{0 \leq x \leq 1 \\ 0 \leq t \leq T}} |(u_1 + u_2)_x|.$$

From (2.5) we obtain, for  $0 \leq t \leq T$

$$\begin{aligned} & \frac{d}{dt} \int_0^1 Z^2(x, t) dx \\ & \leq - \int_0^1 Z_x^2(x, t) dx + 2C_T \left( \int_0^1 Z_x^2(x, t) dx \right)^{1/2} \left( \int_0^1 Z^2(x, t) dx \right)^{1/2} \\ & \leq - \int_0^1 Z_x^2(x, t) dx + \int_0^1 Z_x^2(x, t) dx + C_T^2 \int_0^1 Z^2(x, t) dx \\ & = C_T^2 \int_0^1 Z^2(x, t) dx. \end{aligned}$$

By Grownwall's lemma and (2.4) we get

$$\int_0^1 Z^2(x, t) \leq 0 \quad \forall 0 \leq t \leq T.$$

Thus  $Z^2(x, t) \equiv 0$  for  $0 \leq t \leq T$ ; since  $T$  is arbitrary we get  $Z(x, t) \equiv 0$  i.e.  $u_1(x, t) \equiv u_2(x, t) \forall (x, t) \in D$ .

Next we construct the unique classical solution of (1.1)–(1.4). First we need the following:

*Lemma 2.1. Let  $v(x, t)$  be the solution of*

$$v_t = \frac{1}{2}v_{xx} - (q(x) + \mu)v, \quad (2.6)$$

$$v_x(0, t) + av(0, t) = 0, \quad (2.7)$$

$$v_x(1, t) + bv(1, t) = 0, \quad (2.8)$$

$$v(x, 0) = \exp[-u_0(x)]. \quad (2.9)$$

Then

$$(i) \quad v(x, t) > 0.$$

(ii) *The function*

$$u(x, t) = -\log v(x, t)$$

is a solution of (1.1)–(1.4)

*Proof.* To prove (i), set

$$v(x, t) = \exp(-Cx + Mt)W(x, t), \quad (2.10)$$

where

$$C = \min(a, b)$$

$$M = 1 + \sup_{0 \leq x \leq 1} \left| \frac{C^2}{2} - q(x) - \mu \right|.$$

To prove  $v(x, t) > 0$ , it is enough to prove  $W(x, t) > 0$ . We prove this for the case  $C = a$ , the other case being a similar one.

Assume the contrary, i.e.  $\min W(x, t) \leq 0$ . Now from (2.6)–(2.10) we have  $W(x, t)$  which is solved as

$$\frac{1}{2}W_{xx} - CW_x + \left( \frac{C^2}{2} - q(x) - \mu - M \right)W - W_t = 0, \quad (2.11)$$

$$W_x(0, t) + (a - C)W(0, t) = 0, \quad (2.12)$$

$$W_x(1, t) + (b - C)W(1, t) = 0, \quad (2.13)$$

$$W(x, 0) = \exp[Cx - u_0(x)]. \quad (2.14)$$

Notice that the coefficient of  $W$ ,  $(C^2/2) - q(x) - \mu - M < 0$  by the definition of  $M$  so that the maximum principle can be applied (see Smoller [3]). The maximum principle min of  $W(x, t)$  has to occur at the boundaries and if min is achieved at  $x = 0$ , then  $W_x(0, t) > 0$  and if it occurs at  $x = 1$ ,  $W_x(1, t) < 0$ .

By assumption  $\min W(x, t) \leq 0$ . So evidently min is not achieved for  $t = 0$ . Since  $a = C$ ,  $W_x(0, t) = 0$  (by (2.12)) so again min cannot be achieved at  $x = 0$ . From (2.13)

$$W_x(1, t) = (a - b)W(1, t).$$

But  $a - b \leq 0$ , which implies that  $W(x, t)$  cannot achieve a non-positive minimum at  $x = 1$ . So we have  $\min W(x, t) > 0$ .

To prove (ii) notice that from (i)

$$u(x, t) = -\log v(x, t) \quad (2.15)$$

is well defined. A simple calculation gives

$$u_t = \frac{-\frac{1}{2}v_{xx} + (q(x) + \mu)v}{v},$$

where we used (2.6)

$$\begin{aligned}u_x &= -v_x/v \\ u_{xx} &= (-vv_{xx} + v_x^2)/v^2.\end{aligned}\tag{2.16}$$

It follows that  $u(x, t)$  given by (2.15) is solved as

$$u_t + \frac{1}{2}u_x^2 = \frac{1}{2}u_{xx} + q(x) + \mu.$$

From (2.16), (2.7) and (2.8) we get

$$\begin{aligned}u_x(0, t) &= a, \\ u_x(1, t) &= b.\end{aligned}$$

The proof of lemma is complete.

Let  $\lambda_n$  and  $\phi_n(x)$  be as in §1. In the next lemma we get an expression for the solution of (1.1)–(1.4).

*Lemma 2.2.* Let  $u(x, t)$  be the classical solution of (1.1)–(1.4) then

$$u(x, t) = -\log \left[ \sum_0^\infty \exp [(\lambda_n - \mu)t] a_n \phi_n(x) \right] \tag{2.17}$$

where

$$a_n = \int_0^1 \exp [-u_0(x)] \phi_n(x) dx. \tag{2.18}$$

*Proof.* It is easy to check that  $v(x, t)$  is a solution of (2.6)–(2.9) iff

$$v(x, t) = \exp (-\mu t) V^*(x, t), \tag{2.19}$$

where  $V^*(x, t)$  is the solution of

$$\begin{aligned}V_t^* &= \frac{1}{2}v_{xx}^* - q(x)V^*, \\ V_x^*(0, t) + aV^*(0, t) &= 0, \\ V_x^*(1, t) + bV^*(1, t) &= 0, \\ V^*(x, 0) &= \exp [-u_0(x)] = V_0^*(x).\end{aligned}$$

By separation of variable we obtain

$$V^*(x, t) = \sum_0^\infty \exp (\lambda_n t) a_n \phi_n(x) \tag{2.20}$$

where

$$a_n = \int_0^1 \exp [-u_0(x)] \phi_n(x) dx.$$

From (1.5), (1.9), (1.10) and (1.11) it follows that for  $n > N$ ,  $N$  sufficiently large

$$|\phi_n^{(k)}(x)| \leq Cn^k, \quad k = 0, 1, 2$$

$$|\exp(\lambda_n t)| \leq 1 \quad \forall t \geq 0 \quad (2.21)$$

$$|\exp(\lambda_n t)| \leq C \exp[-\frac{1}{2}(n^2 \pi^2) \delta] \quad \forall t \geq \delta > 0.$$

$C > 0$  is a constant independent of  $x, t$  and  $n$  and  $\phi_n^{(k)}$  denote the  $k$ th derivative w.r.t.  $x$  of  $\phi_n(x)$ . To get an estimate for  $a_n$ , notice that  $V_0^*(x) = \exp[-u_0(x)]$  satisfies

$$\begin{aligned} V_x^*(0) + aV^*(0) &= 0 \\ V_x^*(1) + bV^*(1) &= 0. \end{aligned} \quad (2.22)$$

Multiplying (1.5) by  $V_0^*(x)$ , integrating w.r.t.  $x$ , integrating by parts and using (2.22), we get

$$\frac{1}{2} \int_0^1 \phi_n(x) (V_0^*)''(x) dx = \lambda_n \int_0^1 \phi_n(x) V_0^*(x) dx.$$

for  $n \geq N$

$$a_n = \int_0^1 \phi_n(x) V_0^*(x) dx = \frac{1}{2\lambda_n} \int_0^1 (V_0^*)''(x) \phi_n(x) dx$$

so that, for  $n > N$

$$|a_n| \leq \frac{C}{n^2} |b_n|, \quad (2.23)$$

where

$$b_n = \int_0^1 (V_0^*)''(x) \phi_n(x) dx.$$

From (2.21) and (2.23) it follows that  $V^*(x, t)$  is continuous in  $D$ . Also

$$\begin{aligned} V_x^*(x, t) &= \sum_0^\infty \exp(\lambda_n t) a_n \phi_n'(x) \\ &= \sum_0^{N-1} \exp(\lambda_n t) a_n \phi_n'(x) + \sum_N^\infty \exp(\lambda_n t) a_n \phi_n'(x). \end{aligned} \quad (2.24)$$

Now

$$\left| \sum_N^\infty \exp(\lambda_n t) a_n \phi_n'(x) \right| \leq C^2 \sum_N^\infty \frac{1}{n^2} |b_n| \cdot n \quad (2.25)$$

and

$$\sum_0^\infty |b_n|^2 = \int_0^1 (V_0^*)''(x)^2 dx \quad (2.26)$$

by completeness of eigenfunctions. From (2.24)–(2.26) it follows that

$$|V_x^*(x, t)| \leq \sum_0^{N-1} \exp(\lambda_n t) a_n \phi_n'(x) + C^2 \left( \sum_N^\infty \frac{1}{n^2} \right)^{1/2} \left( \sum_N^\infty |b_n|^2 \right)^{1/2}$$

so that the series is absolutely convergent and hence  $V_x^*(x, t)$  is continuous in  $D$ .

Using the estimates (2.21), (2.23) and (2.26) and using the same argument as before,  $v_{xx}^*(x, t)$  and  $V_t^*(x, t)$  are also continuous in  $0 \leq x \leq 1, t \geq \delta > 0$ . Now by (2.19)

$$v(x, t) = \exp(-\mu t) V^*(x, t) = \sum_0^\infty \exp[(\lambda_n - \mu)t] a_n \phi_n(x)$$

and by lemma 2.1

$$u(x, t) = -\log \left( \sum_0^{\infty} \exp [(\lambda_n - \mu)t] a_n \phi_n(x) \right)$$

is the classical solution of (1.1)–(1.4). The proof of lemma (2.2) is complete.

Next we study the asymptotic behaviour of the solution  $u(x, t)$  constructed in this section.

### 3. Asymptotic behaviour of the classical solution of (1.1)–(1.4)

In this section  $u(x, t)$  denotes the unique classical solution of (1.1)–(1.4) constructed in §2. We shall prove the following.

**Theorem.** (i) Let  $\mu = \lambda_0$ ; then

$$\sup_{0 \leq x \leq 1} |u(x, t) + \log(a_0 \phi_0(x))| \leq C \exp [(\lambda_1 - \lambda_0)t].$$

(ii) Let  $\mu \neq \lambda_0$ ; then for  $t \geq 1$

$$\sup_{0 \leq x \leq 1} \left| \frac{u(x, t)}{t} + (\lambda_0 - \mu) + \frac{\log(a_0 \phi_0(x))}{t} \right| \leq \frac{C}{t} \exp [(\lambda_1 - \lambda_0)t].$$

where  $\lambda_0$  and  $\phi_0(x)$  are as in §1,  $a_0$  is given by (2.18) and  $C$  a positive constant independent of  $x$  and  $t$ .

*Proof.* From lemma 2.2 we have

$$u(x, t) = -\log \left\{ \sum_0^{\infty} \exp [(\lambda_n - \mu)t] a_n \phi_n(x) \right\}.$$

To prove (i), notice that when  $\mu = \lambda_0$ , the above expression for  $u(x, t)$  becomes

$$\begin{aligned} u(x, t) &= -\log \left[ a_0 \phi_0(x) + \sum_1^{\infty} \exp [(\lambda_n - \lambda_0)t] a_n \phi_n(x) \right] \\ &= -\log [a_0 \phi_0(x)] - \log \left[ 1 + \sum_1^{\infty} \exp [(\lambda_n - \lambda_0)t] \left( \frac{a_n}{a_0} \right) \left( \frac{\phi_n(x)}{\phi_0(x)} \right) \right] \end{aligned}$$

$a_0 > 0$  and  $\inf_{0 \leq x \leq 1} \phi_0(x) > 0$  by (1.12). Using the estimate (2.21) we have, for  $t \geq 1$

$$\left| \sum_1^{\infty} \exp [(\lambda_n - \lambda_0)t] \left( \frac{a_n}{a_0} \right) \frac{\phi_n(x)}{\phi_0(x)} \right| \leq C \exp [(\lambda_1 - \lambda_0)t]$$

so that we get

$$|u(x, t) + \log [a_0 \phi_0(x)]| \leq \log \{ 1 + C \exp [(\lambda_1 - \lambda_0)t] \}.$$

But for  $0 < y < 1$

$$\log(1 + y) \leq y$$

we get

$$|u(x, t) + \log(a_0 \phi_0(x))| \leq C \exp[(\lambda_1 - \lambda_0)t].$$

To prove (ii) arguing as before, we get

$$\begin{aligned} u(x, t) = & -\log \{ \exp[(\lambda_0 - \mu)t] a_0 \phi_0(x) \} \\ & - \log \left[ 1 + \sum_1^\infty \exp[(\lambda_n - \lambda_0)t] \left( \frac{a_n}{a_0} \right) \left( \frac{\phi_n(x)}{\phi_0(x)} \right) \right] \end{aligned} \quad (3.4)$$

As before, for  $t \geq 1$ ,

$$\left| \log \left[ 1 + \sum_1^\infty \exp[(\lambda_n - \lambda_0)t] \left( \frac{a_n}{a_0} \right) \left( \frac{\phi_n(x)}{\phi_0(x)} \right) \right] \right| \leq C \exp[(\lambda_1 - \lambda_0)t]. \quad (3.5)$$

Also

$$\log \{ \exp[(\lambda_0 - \mu)t] a_0 \phi_0(x) \} = (\lambda_0 - \mu)t + \log(a_0 \phi_0(x)). \quad (3.6)$$

Using (3.5) and (3.6) in (3.4) we get for  $t \geq 1$

$$\sup_{0 \leq x \leq 1} \left| \frac{u(x, t)}{t} + (\lambda_0 - \mu) + \frac{\log(a_0 \phi_0(x))}{t} \right| \leq \frac{C}{t} \exp[(\lambda_1 - \lambda_0)t].$$

The proof of the theorem is complete.

#### 4. Some remarks on the stationary problem

Consider the stationary problem

$$p_x^2/2 = \frac{1}{2} p_{xx} + q(x) + \lambda_0, \quad (4.1)$$

$$p_x(0) = a, \quad (4.2)$$

$$p_x(1) = b, \quad (4.3)$$

where  $\lambda_0$  is as in §1. Consider the one-parameter family of functions

$$p_\alpha(x) = -\log(\alpha \phi_0(x)), \quad \alpha > 0 \quad (4.4)$$

$\phi_0(x)$  is as in §1 with condition (1.12). By a direct calculation it is easy to verify that  $p_\alpha(x)$  is a solution of (4.1)–(4.3) for each  $\alpha > 0$ . The theorem in §3, part (i) says that, in the case  $\mu = \lambda_0$ , the solution  $u(x, t)$  of (4.1)–(4.4) converges to  $p_{a_0}(x)$ .

The second part of the theorem says that if  $\mu \neq \lambda_0$ ,  $u(x, t) \rightarrow \infty$  as  $t \rightarrow \infty$ . The stationary problem

$$p_x^2/2 = \frac{1}{2} p_{xx} + q(x) + \mu, \quad (4.5)$$

$$p_x(0) = a, \quad (4.6)$$

$$p_x(1) = b \quad (4.7)$$

does not have any solution if  $\mu \neq \lambda_0$ . In fact as in lemma 2.1 one can easily check

that  $p(x)$  is a solution to (4.5)–(4.7) iff the function

$$h = \exp(-p) \quad (4.8)$$

is a solution to

$$\frac{1}{2}h_{xx} = [q(x) + \mu]h, \quad (4.9)$$

$$h_x(0) + ah(0) = 0, \quad (4.10)$$

$$h_x(1) + bh(1) = 0. \quad (4.11)$$

But (4.9)–(4.11) has a non-zero solution iff  $\mu = \lambda_n$ ,  $n = 0, 1, 2, \dots$ , and  $\lambda_n$  is as in §1. Further the corresponding solution has to be positive, by (4.8). This happens iff  $\mu = \lambda_0$ .

### 5. Burger's equation

The method presented in previous sections can be used to get a closed form expression and asymptotic behaviour of the solution of the Burger's equation in  $D = \{ |x, t| : a \leq x \leq 1, t \geq 0 \}$ . Let  $u(x, t)$  be the unique solution to

$$u_t + (u^2/2)_x = \frac{1}{2}u_{xx}, \quad (5.1)$$

$$u(x, 0) = u_0(x), \quad (5.2)$$

$$u(0, t) = a, \quad (5.3)$$

$$u(1, t) = b, \quad (5.4)$$

$a$  and  $b$  are constants,  $u_0(x)$  is  $C^1[0, 1]$  and  $u_0(0) = a$ ,  $u_0(1) = b$ . Then it can be easily seen as in lemmas 2.1 and 2.2 that

$$u(x, t) = - \frac{\sum_{n=0}^{\infty} \exp(\lambda_n t) a_n \phi_n^1(x)}{\sum_{n=0}^{\infty} \exp(\lambda_n t) a_n \phi_n(x)} \quad (5.5)$$

$\lambda_0 > \lambda_1 > \dots \rightarrow -\infty$  are the eigenvalues of

$$\frac{1}{2}\phi_{xx} = \lambda\phi,$$

$$\phi_x(0) + a\phi(0) = 0,$$

$$\phi_x(1) + b\phi(1) = 0$$

and  $\phi_n(x)$ ,  $n = 0, 1, \dots$  are the normalized eigenfunctions corresponding to  $\lambda_n$  with  $\phi_0(x) > 0$  and

$$a_n = \int_0^1 \exp\left(-\int_0^x u_0(y) dy\right) \phi_n(x) dx. \quad (5.6)$$

As in §3, it is easy to prove the following.

**Theorem.** Let  $u(x, t)$  be the solution of (5.1)–(5.4); then

$$u(x, t) = -[\log \phi_0(x)]_x + O[\exp [(\lambda_1 - \lambda_0)t]]$$

uniformly in  $x \in [0, 1]$ .

*Proof.* From (5.5) we have,

$$\begin{aligned} u(x, t) &= - \frac{\left[ \phi_0^1(x) + \sum_1^\infty \exp [(\lambda_n - \lambda_0)t] \left( \frac{a_n}{a_0} \right) \phi_n^1(x) \right]}{\phi_0(x)} \\ &\quad \times \left[ 1 + \sum_1^\infty \exp [(\lambda_n - \lambda_0)t] \left( \frac{a_n}{a_0} \right) \left( \frac{\phi_n}{\phi_0} \right) \right]^{-1} \\ &= - \frac{\left[ \phi_0^1(x) + \sum_1^\infty (\exp [(\lambda_n - \lambda_0)t]) \right]}{\phi_0(x)} \\ &\quad \times [1 + O(\exp [(\lambda_1 - \lambda_0)t])]^{-1} \text{ uniformly in } x \\ &= - \frac{\phi_0^1(x)}{\phi_0(x)} [1 + O(\exp [(\lambda_1 - \lambda_0)t])] \end{aligned}$$

uniformly in  $x \in [0, 1]$ . The proof of theorem is complete.

q.e.d.

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