NON-ISOTHERMAL EFFECTIVENESS FACTORS IN BIDISPERSED CATALYSTS UNDER CONDITIONS OF EXTERNAL TRANSPORT RESISTANCES

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ABSTRACT

Effectiveness factors in bidispersed catalysts under conditions of internal and external heat and mass transport resistances have been obtained numerically. The parametric dependence of the utilization of such catalysts has been formulated in the form of an equation which predicts the values to within a per cent accuracy.

INTRODUCTION

A large number of studies concerning reaction and diffusion in bidispersed catalysts are reported in chemical engineering and a recent survey by Dogu\(^1\) covers some of the important aspects and conclusions of these studies. Ors and Dogu\(^2\) define a parameter \(\alpha\) for bidispersed system to characterize the ratio of diffusion times in the macro and micro particle regions. The approach of Ors and Dogu\(^2\) is convenient and its application to analysis of several nonlinear rate forms such as power law\(^3,4\), Langmuir-Hinshelwood types of rate models\(^5\) and to analysis of bidispersed catalysts under non-isothermal conditions\(^6\) has been subsequently reported. In this note we present the results of reaction-diffusion behaviour of bidispersed catalysts for a first order reaction under non-isothermal conditions where external transport resistances become important.

This particular aspect does not appear to have been studied so far and is relevant when problems such as parametric sensitivity in pellets or in fixed-bed reactors become important. In fact, while a large number of studies dealing with parametric sensitivity have been reported\(^1\) for conventional monodispersed catalysts those in presence of bidispersed catalyst have yet to be studied. This apparently a simple problem is compounded by severe numerical problems and as will be seen in the present work, numerical integration of the governing equations requires substantial amount of central processing unit (CPU) time even when run on fast computers. As a first step towards this more general study, in the present work, we report the values of effectiveness factors for bidispersed catalysts under non-isothermal conditions in the presence of external transport resistances. The study covers substantial variations in the system parameters and final results are correlated in the form 2 of an equation which predicts the effectiveness factors as a function of system parameters with reasonable degree of accuracy.

Problem formulation

Referring to figure 1, which shows a schematic diagram of the spherical geometry of micro-macro particle catalyst system, the governing

![Figure 1. Diagram of macro-micro particles in a pellet.](image-url)
equations for mass and heat conservation can be written in non-dimensional form as:

**Microparticle**

\[
\nabla^2_x C_i = \phi^2 \exp \left[ \gamma \left(1 - \frac{1}{T_i} \right) \right] C_i
\]

with boundary conditions;

\[
C_i(x = 1) = C_a, \frac{dC_i}{dx} \bigg|_{x=0} = 0 \quad (1)
\]

\[
\nabla^2_x T_i = -\phi^2 \beta \exp \left[ \gamma \left(1 - \frac{1}{T_i} \right) \right] C_i
\]

with boundary conditions;

\[
T_i(x = 1) = T_a, \frac{dT_i}{dx} \bigg|_{x=0} = 0. \quad (2)
\]

**Macroparticle**

\[
\nabla^2_y C_a = \alpha \frac{dC_a}{dx} \bigg|_{x=1}
\]

with boundary conditions,

\[
\frac{dC_a}{dy} \bigg|_{y=0} = 0; \quad \frac{dC_a}{dy} \bigg|_{y=1} = \text{Sh}(1 - C_{as}) \quad (3)
\]

\[
\nabla^2_y T_a = \alpha_T \frac{dT_a}{dx} \bigg|_{x=1}
\]

with boundary conditions,

\[
\frac{dT_a}{dy} \bigg|_{y=0} = 0; \quad \frac{dT_a}{dy} \bigg|_{y=1} = \text{Nu}(1 - T_a^*) \quad (4)
\]

where \( C_a \) is defined as dimensionless concentration in the macroparticle normalized with respect to concentration in the bulk and the various parameters appearing in (1), (2), (3) and (4) are defined as follows:

\[
\phi = \sqrt{(k/D_i)r_i^2}; \quad \beta = (-\Delta H) \frac{D_i}{k_i} \frac{C_{as}}{T_i^*};
\]

\[
\alpha = 3(1-\varepsilon) \frac{D_i R_i^2}{D_a r_i^2}; \quad \alpha_T = 3(1-\varepsilon) \frac{k_i R_i^2}{k_p r_i^2}
\]

\[
\gamma = (E/R_g T_i^*); \quad \text{Sh} = \frac{k_g R}{D_a}; \quad \text{Nu} = \frac{h_f R}{k_p}.
\]

\( T_i \) and \( T_a \) from (1), (2), (3) and (4) can be eliminated appropriately to give governing microparticle mass conservation equation as:

\[
\nabla^2_x C_i = \phi^2 \exp \left[ \gamma \left(1 - \frac{1}{C + \beta C_a \left(1 + \frac{\alpha_T}{\alpha} \right) - \beta C_i} \right) \right] C_i
\]

with boundary conditions;

\[
C_i(x = 1) = C_a, \frac{dC_i}{dx} \bigg|_{x=0} = 0. \quad (5)
\]

Equations (3) and (5) together describe the behaviour of the system subject to appropriate initial and boundary conditions. Numerical integration of these equations involves considerable trial and error besides the small step size requirements due to stiffness of the equations. It is desirable to avoid trial and error in the numerical computations and hence an alternate scheme, as will be explained shortly, yield the concentration and temperature profiles in the micro and macro particles and the information is used to compute the eventual quantities of interest which are defined as

\[
\eta = \left[ \frac{dC_a}{dy} \bigg|_{y=1} \right] / \alpha \phi^2 \quad (6)
\]

\[
\frac{dT_a}{dy} \bigg|_{y=1} = \beta \frac{\alpha_T}{\alpha} \left[ \frac{dC_a}{dy} \bigg|_{y=1} \right] \quad (7)
\]

\[
T_a^* = C + \beta \frac{\alpha_T}{\alpha} C_{as} \quad (8)
\]

\[
\text{Sh} = \left[ \frac{dC_a}{dy} \bigg|_{y=1} \right] / (1 - C_{as}) \quad (9)
\]

\[
\text{Nu} = \left[ \frac{dT_a}{dy} \bigg|_{y=1} \right] / (T_a^* - 1) \quad (10)
\]
The detailed steps involved in solving (3) and (5) are:

(i) Equation (5) represents a boundary value problem and solution of it to obtain the concentration profile in the micropore, for a defined set of parameter values $\alpha$, $\alpha_T$, $\beta$, $\gamma$ and $\phi$, requires considerable trial and error. It is possible however, to integrate this equation using Weisz and Hicks' method as an initial value problem. This however, requires to fix the calculation concentration at the center of the micropore and for the set of parameters $(\alpha, \alpha_T, \beta, \gamma)$ and other variables $(C$ and $C_0)$ defined, the integration gives the value of $\phi$.

(ii) For the set of parameter values $(\alpha, \alpha_T, \beta, \gamma)$ assumed and for the values of $C$ and $C_0$ fixed, step (i) gives $\phi$ for assumed values of $c_0$. These values of $C_{i0}$ and $\phi$ (for same values of other parameters) when used in (5) gives the value of surface flux at the micropore.

(iii) The value of surface flux $(dc_i/dx)|_{x=1}$ as obtained in step (ii) is then used in macropore (3) which is integrated in $y$ direction to obtain the new value of $C_a$ and $dC_a/dy$ at the next incremental step.

(iv) With this new value of $C_a$ obtained in step (iii) we shall return to step (i) to redefine the boundary condition $C_i|x=1 = C_a$. For the other parameters held constant, this change in the value of $C_a$—for the same value of $\phi$—suggests that $c_{i0}$ would now be different. There is no easy way, however, to calculate this changed value of $C_{i0}$. In fact, since step (i) envisages to regard (5) as an initial value problem we should specify this value of $C_{i0}$ a priori.

To circumvent this problem the procedure followed here assumes several values of $C_{i0}$ in the neighbourhood of the value assumed in the earlier stage of step (i) and corresponding to each of these $C_{i0}$ values obtains the values of $\phi$ for the value of $C_a$ from step (iii) and all other parameters same as before. The information is interpolated and used to obtain the precise value of $C_{i0}$ that corresponds to the value of $\phi$ as has been originally fixed in step (i). With this new set of $C_{i0}$, $C_a$ and the other parameters that remain unaltered, step (ii) suggesting integration of (5) gives the surface flux $dC_i/dx$. We then proceed to step (iii) and the new values of $C_a$ and $dC_a/dy$.

(v) The entire procedure as suggested in steps (i–iv) is repeated until we reach the macroparticle surface (i.e. $y = 1$) where the values of $C_a|s$ and $(dC_a/dy)|s$ are noted.

(vi) Knowing the macroparticle surface concentration and the flux now allows us to determine the extent of external film resistance ($Sh$) and the value of the effectiveness factors using (6), (7), (8), (9) and (10).

(vii) The entire procedure is repeated for different sets of parameter values $(\alpha, \alpha_T, \beta, \gamma)$ and different values of initial conditions such that $C_{i0}$, $C_a$, $C$ etc. in step (i). The general results obtained in the form of $\eta$ vs $Sh$ or $Nu$ number as a function of $\phi$ and other parameters can then be plotted using the calculated values.

A substantial range of parameters was investigated although the equations were becoming stiff in certain regions. The step size employed was of the order of $10^{-4}$ to $10^{-5}$ and in such instances the CPU time consumed by the NEC 1000 supercomputer was in the range of 4–5 h.

RESULTS AND DISCUSSION

The governing equations were solved by the procedure outlined earlier to obtain the values of $\eta$, $Sh$ and $Nu$. The corresponding $\phi$ values are also listed in the plots.

Figures 2a and 3a show $Sh$ vs $\eta$ and $Nu$ vs $\eta$ obtained by varying $\phi$. Here $\eta$ shows a monotonic increase as the value of $\phi$ is increased. The appropriate value of $C_{i0}$ as calculated using the procedure suggested is also included along with the value of $\phi$ to provide checks on the numerical calculations. $\eta$ monotonically increases as the center micropore concentration decreases. The subsequent results will be compared with this case to seek the influence of various parameters.
Figure 2a-i. Sherwood number vs effectiveness factor plot.

Figures 2b and 3b indicate the influence of parameter \( C \). It can be easily seen that the nature of these plots does not differ significantly from the previous case except for the actual magnitudes of the derived quantities. From the definition of parameter \( C \), it is clear that \( C \) can increase either due to the increase in surface temperature or decrease in surface concentration and results in enhancing \( \eta \) and marginal decrease in \( \phi \) values.
The initial macropore concentration $C_0$ was varied (0.7 and 0.5) and the results for each, for different values of $C_{\text{so}}$ are depicted in figures 2e, f and 3c, f. It is observed that lowering of the $C_a$ value results in the substantial reduction in the $\phi$ value.

The results obtained by changing $\beta$ value are shown in figures 2g and 3g. The well known feature of reduction in $\eta$ values with decrease in $\beta$ is clearly noticed. The Sherwood number rises sharply for the lower $C_{\text{so}}$ values tried.

In figures 2h and 3h the effect of higher $\alpha_T$ is shown. The high $\eta$'s obtained are obviously due to the increase in the thermal conductivity of macropore effecting higher reaction rates.

Increase in the magnitude of parameter $\gamma$ signifying dimensionless activation energy results in very high values of effectiveness factors (figures 2i and 3i).

In view of the large computer time required for the procedure adopted here, it is more appropriate to summarize these results and present them in the form of an equation. The several data points generated have been fitted in the form of a correlation by using nonlinear optimization routine based on Marquardt's algorithm. As could be noticed the general form of the equation has been chosen in analogy with the form for limiting situation and appropriately modified by inserting arbitrary exponents and coefficients. The form of the equation is thus given by:

$$\eta = \frac{b_1 \left( \frac{\text{Sh.Nu}}{\alpha_T} \right)^{b_2} \left( \frac{\text{Sh.Nu}}{\alpha} \right)^{b_3} \left( \frac{\text{Sh.Nu}}{\alpha_T} \right)^{0.2-\beta}}{c \phi^2 (\text{Sh} + \text{Nu})},$$

where the values of the coefficients and the exponents are: $b_1 = 0.642$, $b_2 = 0.8413$, $b_3 = 0.00935$ and $b_4 = 0.00172$. It should be noted that the functional form of (11) is purely empirical and should not mislead one to conclude that the heat and mass transfer resistances are in series.

The equation has been tested for all the results discussed above and for the parameter values lying in between. In each one of these
cases the value of \( \eta \) given by (11) conforms within \(< 1\%\) error in comparison with the value obtained numerically.

The expression thus obtained should therefore prove useful in predicting the effectiveness factors and also in related studies on the sensitivity behaviour of pellets or that of fixed bed reactors. While for all other parameter variations reported here (11) yields almost accurate results, the prediction becomes sensitive to the variations in the parameter \( \beta \) beyond the range investigated.

**Notation**

- \( b_1, b_2 \): coefficients and exponents appearing in equation (11)
- \( C \): parameter defined as \( T_s - C_{as} \beta (\alpha T / \alpha) \)
- \( C_a \): dimensionless concentration in macroparticle
- \( C_i \): dimensionless concentration in microparticle
- \( C_{as} \): concentration of species at the surface of the catalyst pellet
- \( D_a \): effective diffusivity in macroparticle
- \( D_i \): effective diffusivity in microparticle
- \( h_f \): heat transfer coefficient
- \( k_i \): thermal conductivity in microparticle
- \( k_p \): thermal conductivity in macroparticle
- \( k_y \): mass transfer coefficient
- \( N_u \): Nusselt number
- \( r_i \): radius of the microparticle
- \( R \): radius of macroparticle
- \( Sh \): Sherwood number
- \( T_w \): dimensionless temperature in macroparticle
- \( T_i \): dimensionless temperature in microparticle
- \( T_s \): temperature at the surface of the catalyst pellet
- \( x \): dimensionless distance variable used in microparticle
- \( y \): dimensionless distance variable used in macroparticle
- \( \alpha \): parameter defined as \( 3(1 - \varepsilon)(D_i / D_a) \) \((R^2 / r_i^2)\)
- \( \alpha_T \): parameter defined as \( 3(1 - \varepsilon)(k_i / k_p) \) \((R^2 / r_i^2)\)
- \( \beta \): exothermicity factor defined as \((- \Delta H)D, C_{as}/k, T_s\)
- \( \gamma \): defined as \( E/R_{k}, T_s\)
- \( \eta \): effectiveness factor
- \( \phi \): Thiele modulus
- \( \varepsilon \): void fraction of the pellet.

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